



New results on heavy flavour production at LHCb

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On behalf of the LHCb collaboration 2019.01.29

CERN LHC seminar

Outline

> Introduction

New results on heavy flavour production

• b-hadron production fractions at 13 TeV

I HCh-PAPER-2

LHCb-PAPER-2018-050, in preparation

• The mass and production rate of \mathcal{Z}_b^- baryons

LHCb-PAPER-2018-047, arXiv:1901.07075 submitted to Phys.Rev.D

• $\psi(2S)$ production at 13 TeV and 7 TeV

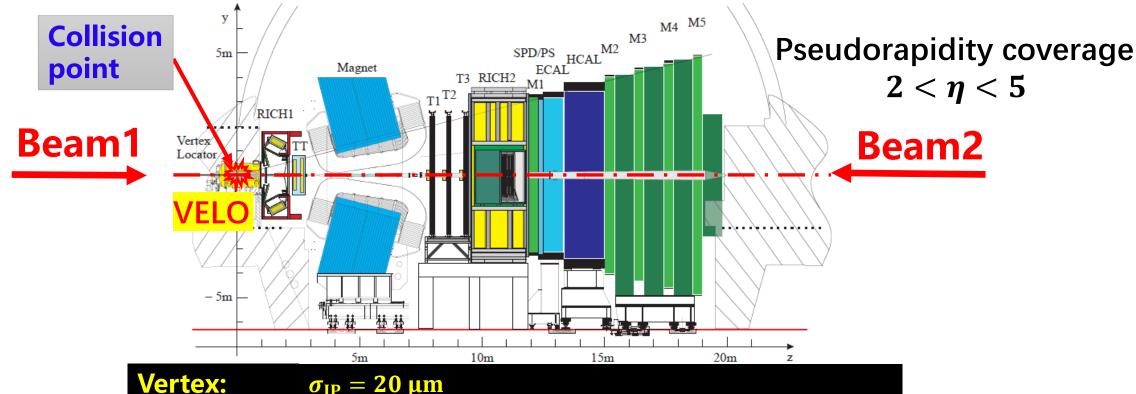
LHCb-PAPER-2018-049, in preparation

Prospects and summary

The LHCb experiment

The LHCb detector

JINST 3 (2008) S08005 Int. J. Mod. Phys. A 30 (2015) 1530022



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Vertex: \sigma_{\text{IP}} = 20 \ \mu\text{m}

Time: \sigma_{\tau} = 45 \ \text{fs} for B_s^0 \to J/\psi \phi or D_s^+ \pi^-

Momentum: \Delta p/p = 0.4 \sim 0.6\% \ (5-100 \ \text{GeV}/c)

Mass: \sigma_m = 8 \ \text{MeV}/c^2 for B \to J/\psi X (constraint m_{J/\psi})

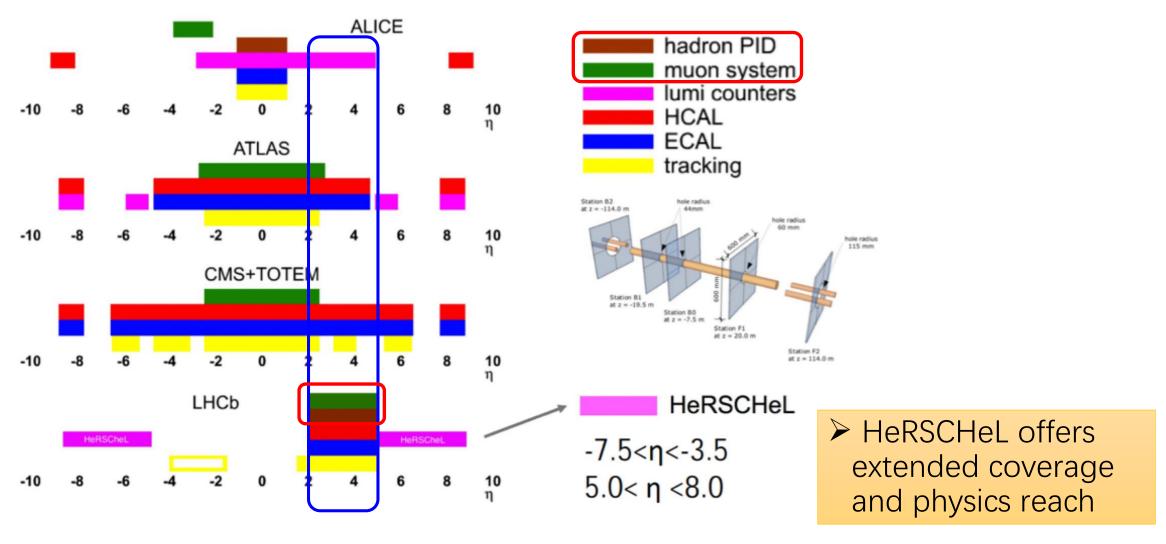
Hadron ID: \varepsilon(K \to K) \sim 95\% mis-ID \varepsilon(\pi \to K) \sim 5\%

Muon ID: \varepsilon(\mu \to \mu) \sim 97\% mis-ID \varepsilon(\pi \to \mu) \sim 1-3\%
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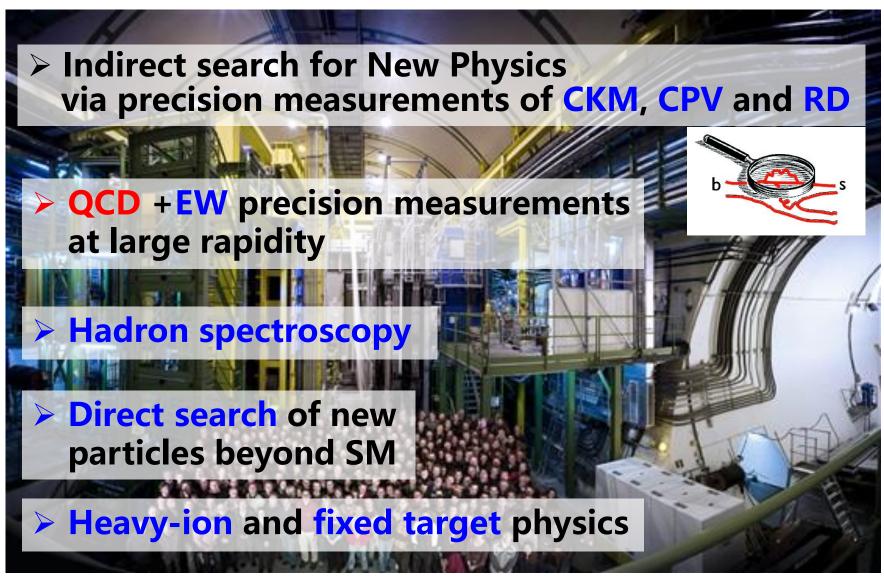
ECAL:

 $\Delta E/E = 1\% \oplus 10\%/\sqrt{E \text{ (GeV)}}$

Fully instrumented at forward coverage

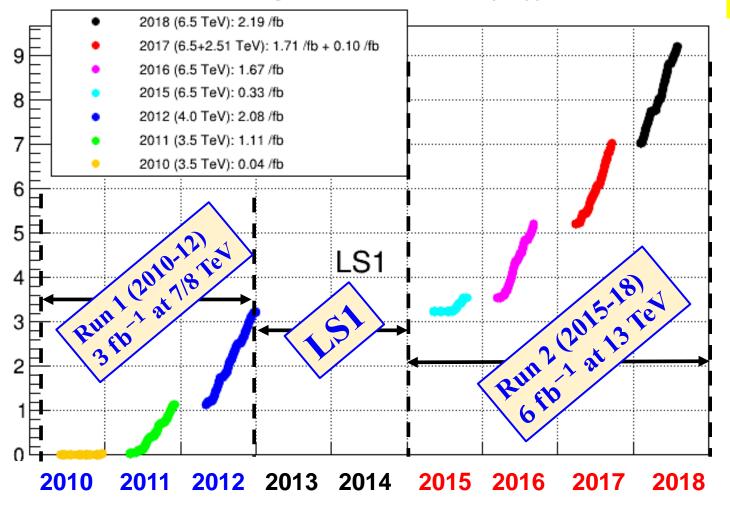


The Physics of LHCb



Data taking (run1+run2)

LHCb Cumulative Integrated Recorded Luminosity in pp, 2010-2018

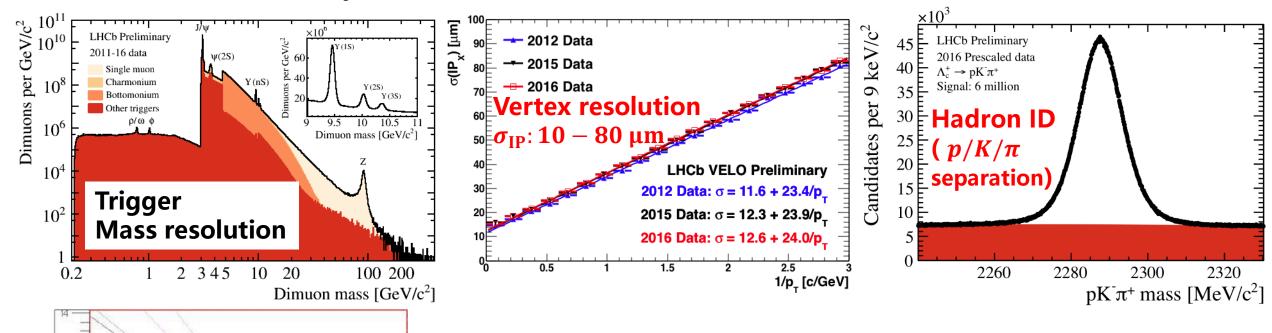


Great thanks to the LHC!

- ightharpoonup A huge amount of $b\bar{b}$ and $c\bar{c}$ have been produced
 - $\sim 10^{12} \ b\bar{b}$
 - $\sim 10^{13} \ c\bar{c}$
- Many impressive results have been achieved

More than 9 fb⁻¹ accumulated in Run1+Run2

Pros of heavy flavour measurement at LHCb



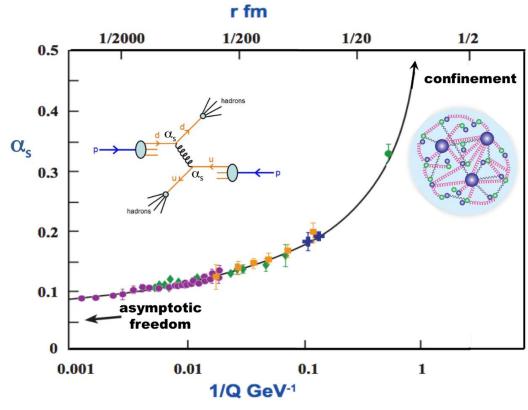
- Large production cross-section
- Efficient trigger
- Vertex locator with high precision
- High precision tracking system
- Powerful hadron identification
- Efficient muon system

LHCb Preliminary EVT: 49700980 RUN: 70684

2019/01/29

Physics motivation

- ➤ The theory of strong interactions, quantum chromodynamics (QCD), is the least understood part of the Standard Model
 - Well tested in the perturbative region
 - Nonperturbative behaviours remain mysterious
- ➤ Measurements of heavy flavour production and other properties can provide essential information to deeply understand QCD
 - Understanding of QCD is also important to improve sensitivity of New Physics searches



[S. Olsen et al, Rev. Mod. Phys. 90 (2018) 015003]

b-hadron fractions in pp collisions at 13 TeV: $f_s/(f_u+f_d)$ and $f_{\Lambda_b^0}/(f_u+f_d)$

LHCb-PAPER-2018-050 in preparation

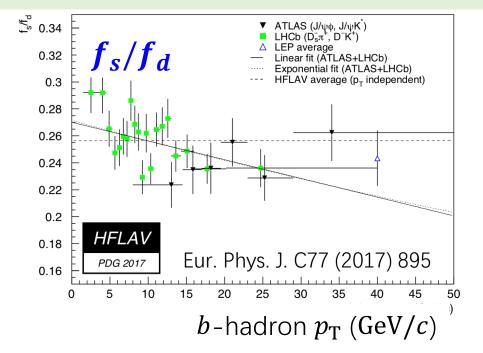
b-hadron fragmentation fractions

- ➤ Knowledge of *b*-hadron fragmentation fractions is essential in many aspects
 - To allow for relating the $b\bar{b}$ production cross-section from pQCD to the observed b hadrons
 - To convert the observed \bar{B}^0_s or Λ^0_b production ratios at the LHC into absolute branching fractions
 - For example, Measurement of $\mathcal{B}(B_s^0 \to \mu^+\mu^-)$
 - To characterise the signal (background) composition in inclusive (exclusive) b-hadron analyses
- > These fractions must be determined experimentally
 - They cannot be reliably predicted because they are dominated by longdistance strong interactions

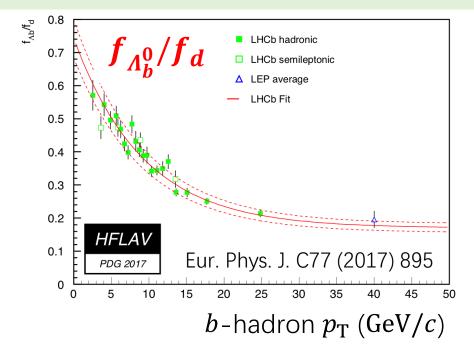
Previous measurements

Phys. Rev. D 85 (2012) 032008 JHEP 04 (2013) 001 JHEP 08 (2014) 143

- \succ LHCb measured kinematic dependences of f_s/f_d and $f_{\Lambda_b^0}/f_d$ at 7 TeV using both semileptonic and hadronic decays
 - The dependence of f_s/f_d on the b-hadron $p_{\rm T}$ is not conclusive



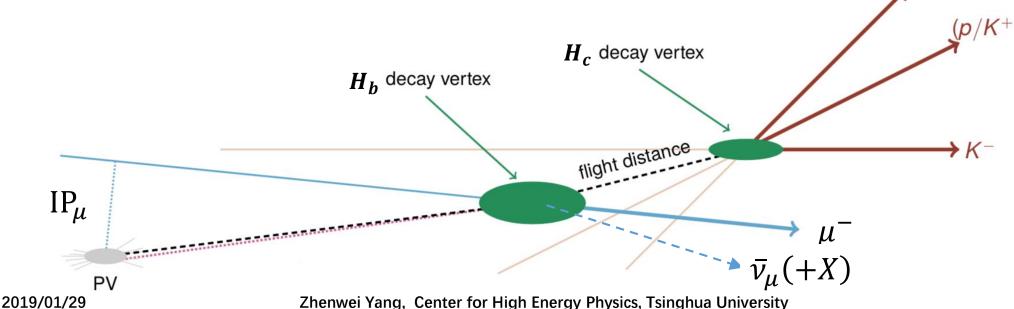
• $f_{A_b^0}/f_d$ is observed to be strongly dependent on the b-hadron $p_{\rm T}$



Analysis strategy at 13 TeV

- ➤ Data sample: 1.67 fb⁻¹ collected in 2016
- > Inclusive semileptonic decays $H_b \to H_c \mu^- \bar{\nu}_{\mu} X$
- > Theoretical basis: Semileptonic widths for all b-hadrons are almost equal $(\Gamma_{SL}(H_h) = \Gamma_{SL})$ [I. Bigi et al, JHEP 09 (2011) 012]

• Differences predicted to be around 1% (heavy quark expansion)



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Analysis strategy (cont.)

- \succ Semileptonic branching fractions $\mathcal{B}_{\mathrm{SL}}$ for \bar{B}_{s}^{0} and Λ_{b}^{0} calculated with well measured lifetimes and $\mathcal{B}_{\mathrm{SL}}$ for \bar{B}^{0} and B^{-}
 - $\mathcal{B}_{SL}(H_b) = \Gamma_{SL}/\Gamma(H_b) = \Gamma_{SL} \cdot \tau_{H_b}$
- \triangleright With known $\mathcal{B}_{SL}(H_b)$, only the ratios of yields need to be measured

Particle	$\tau \text{ (ps)}$	$\mathcal{B}_{\mathrm{SL}}$ (%)	Correction (%)	$\mathcal{B}_{\mathrm{SL}}$ $(\%)$
	measured	measured	[5]	used
$\overline{B}{}^0$	1.520 ± 0.004	10.30 ± 0.19		10.30 ± 0.19
B^-	1.638 ± 0.004	11.08 ± 0.20		11.08 ± 0.20
$\langle \overline{B}{}^0 + B^- \rangle$		10.70 ± 0.19		10.70 ± 0.19
$ar{B}^0_s$	1.526 ± 0.015		-1.0 ± 0.5	10.24 ± 0.21
$ar{B}^0_s \ A^0_b$	1.470 ± 0.010		3.0 ± 1.5	10.24 ± 0.25

Ref. [5]: I. Bigi et al, JHEP 09 (2011) 012

Formula of the fragmentation ratio: $f_{s(\Lambda_h^0)}/(f_u + f_d)$

$$\frac{N_{\rm SL}^{\rm obs}(\bar{B}_{S}^{0})}{N_{\rm SL}^{\rm obs}(B)} = \frac{\sigma_{b\bar{b}} f_{s}}{\sigma_{b\bar{b}}(f_{u} + f_{d})} \frac{\mathbf{B}_{\rm SL}(\bar{B}_{S}^{0})}{\mathbf{B}_{\rm SL}(B)} \frac{\varepsilon(\bar{B}_{S}^{0})}{\varepsilon(B)}$$

$$= \frac{f_{s}}{f_{u} + f_{d}} \frac{\Gamma_{\rm SL}(\bar{B}_{S}^{0}) \tau_{\bar{B}_{S}^{0}}}{\Gamma_{\rm SL}(B)(\tau_{B^{-}} + \tau_{\bar{B}^{0}})/2} \frac{\varepsilon(\bar{B}_{S}^{0})}{\varepsilon(B)}$$

Similar for Λ_h^0

Consider the corrections and use $\mathcal{B}_{SL}(H_b) = \Gamma_{SL} \cdot \tau_{H_b}$

$$\frac{f_{s}}{f_{u}+f_{d}} = \frac{N_{\rm SL}^{\rm corr}(\bar{B}_{s}^{0})}{N_{\rm SL}^{\rm corr}(B)} + \frac{2\tau_{\bar{B}_{s}^{0}}}{\tau_{B^{-}}+\tau_{\bar{B}^{0}}} (1-\xi_{s}) + \text{corr. term}$$
Measurable

Known

Formula of the fragmentation ratio: $f_{s(\Lambda_b^0)}/(f_u + f_d)$

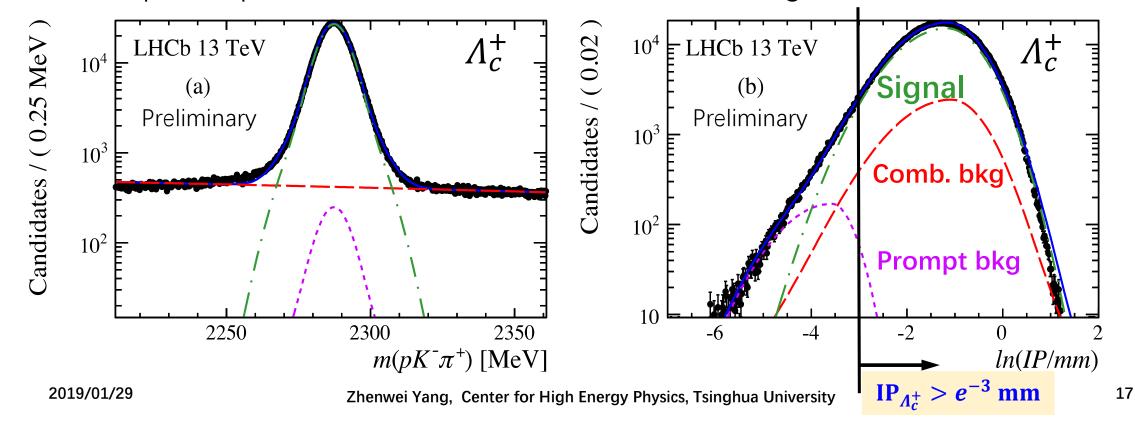
$$\frac{f_s}{f_u+f_d} = \frac{n_{\mathrm{corr}}(\overline{B}_s^0 \to D_s \mu)}{n_{\mathrm{corr}}(B \to D^0 \mu) + n_{\mathrm{corr}}(B \to D^+ \mu)} \underbrace{\tau_{B^-} + \tau_{\overline{B}^0}}_{2\tau_{\overline{B}_s^0}} (1-\xi_s) \qquad \text{SU(3) breaking correction} \\ \xi_s = (-1 \pm 0.5)\%$$

$$\frac{\mathcal{B}(B \to D_s K \mu)}{\langle \mathcal{B}_{SL} \rangle} \underbrace{\epsilon(\overline{B}_s^0 \to D_s^+)}_{\epsilon(\overline{B}_s^0 \to D_s^+)} \qquad \text{Subtraction of } \overline{B}^0(B^-) \to D_s^+ K \mu^- \overline{\nu}_\mu X \text{ contributions in } \overline{B}_s^0 \text{ signals}$$

$$\frac{f_{\Lambda_b^0}}{f_u + f_d} = \underbrace{\frac{n_{\mathrm{corr}}(\Lambda_b^0 \to D\mu)}{n_{\mathrm{corr}}(B \to D^0\mu) + n_{\mathrm{corr}}(B \to D^+\mu)}}_{\begin{array}{c} T_{B^-} + \tau_{\bar{B}^0} \\ 2\tau_{\Lambda_b^0} \end{array} \underbrace{(1 - \xi_{\Lambda_b^0})}_{\begin{array}{c} T_{B^-} + \tau_{\bar{B}^0} \\ 2\tau_{\Lambda_b^0} \end{array}}_{\begin{array}{c} T_{B^-} + \tau_{\bar{B}^0} \\ 2\tau_{\Lambda_b^0} \end{array}}_{\begin{array}{c} T_{B^-} + \tau_{\bar{B}^0} \\ 2\tau_{\Lambda_b^0} \end{array}} \underbrace{(1 - \xi_{\Lambda_b^0})}_{\begin{array}{c} T_{B^-} + \tau_{\bar{B}^0} \\ 2\tau_{\Lambda_b^0} \end{array}}_{\begin{array}{c} T_{B^$$

Removal of prompt charmed hadrons

- > Greatly suppressed by lifetime related requirements
 - χ^2 of charmed hadron flight distance and $\chi^2_{\rm IP}$ of final tracks (μ, p, K, π)
- \triangleright Remaining prompt H_c removed by requiring $\ln(IP_{H_c}/mm) > -3$
 - Prompt component reduced to below 0.1%, while signal loss is around 3%



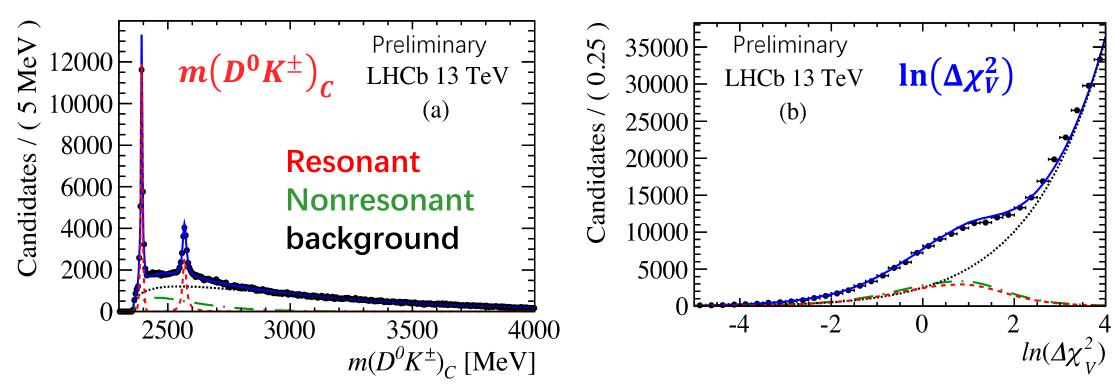
Nonresonant contribution: $\bar{B}_s^0 \to DK\mu^-\bar{\nu}_{\mu}X$

LHCb-PAPER-2018-050 in preparation

- \triangleright Signals for \bar{B}_s^0 and background for $\bar{B}^0(B^-)$
- \triangleright Extracted by 2D fits: $m(D^0K^{\pm})_C$ v.s. $\ln(\Delta\chi_V^2)$

$$m(D^0K^{\pm})-m(D^0)+m(D^0)_{\mathrm{PDG}}$$

logarithm of the vertex χ^2 difference between $D\mu K$ and $D\mu$



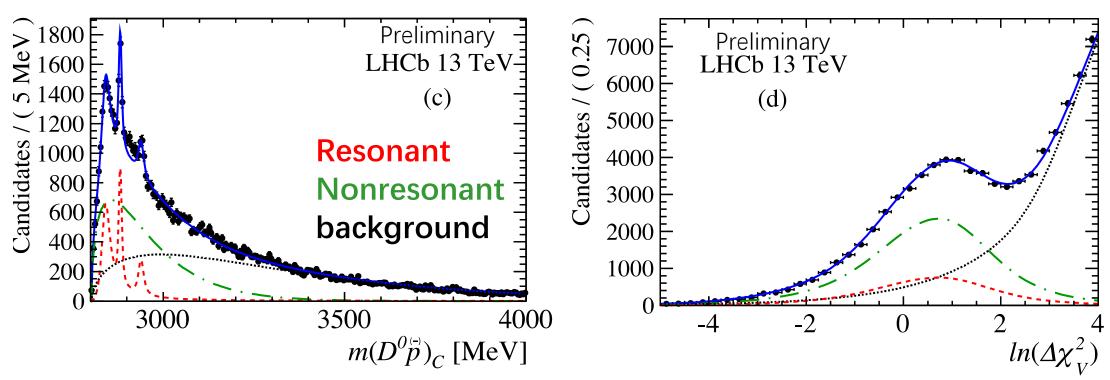
Nonresonant contribution: $\Lambda_b^0 \to D^0 p \mu^- \bar{\nu}_{\mu} X$

LHCb-PAPER-2018-050 in preparation

- \triangleright Signals for Λ_b^0 and background for $\bar{B}^0(B^-)$
- \triangleright Extracted by 2D fits: $m(D^0p)_C$ v.s. $\ln(\Delta \chi_V^2)$

$$m(D^0p)-m(D^0)+m(D^0)_{PDG}$$

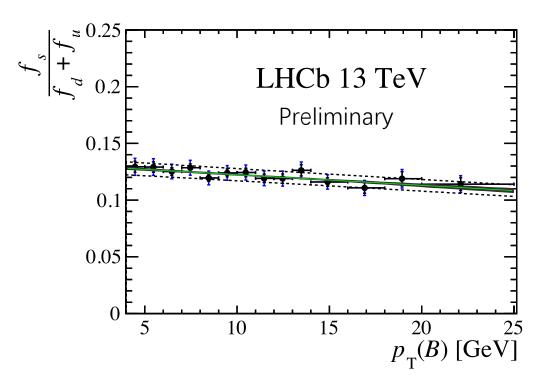
logarithm of the vertex χ^2 difference between $D\mu p$ and $D\mu$

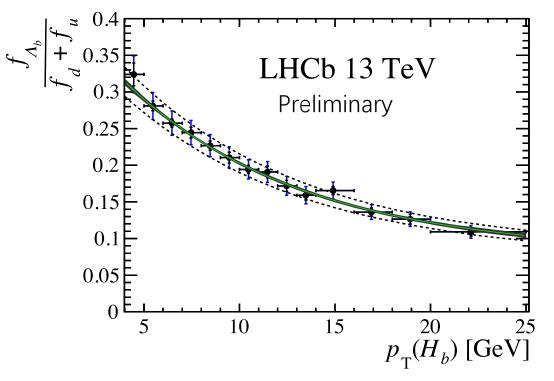


$p_{\rm T}$ dependence

LHCb-PAPER-2018-050 in preparation

 $> p_{\rm T}(H_b) = k p_{\rm T}(H_c\mu)$: correction factor $k = \langle p_{\rm T}^{\rm rec} \rangle / p_{\rm T}^{\rm true}$ from simulation



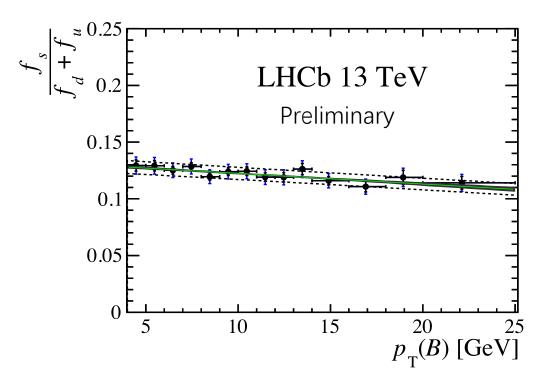


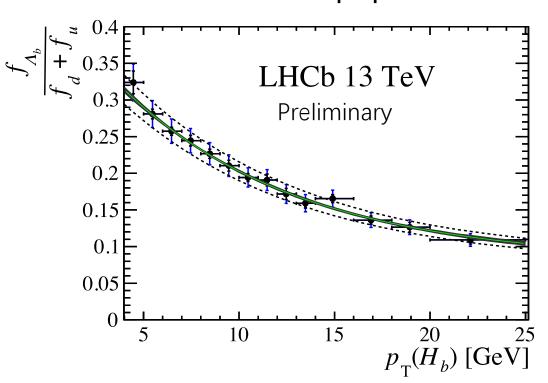
• $f_s/(f_u+f_d)$ slightly depends on $p_T(B)$

• $f_{\Lambda_b^0}/(f_u+f_d)$ strongly depends on $p_{\rm T}(\Lambda_b)$

$p_{\rm T}$ dependence

LHCb-PAPER-2018-050 in preparation





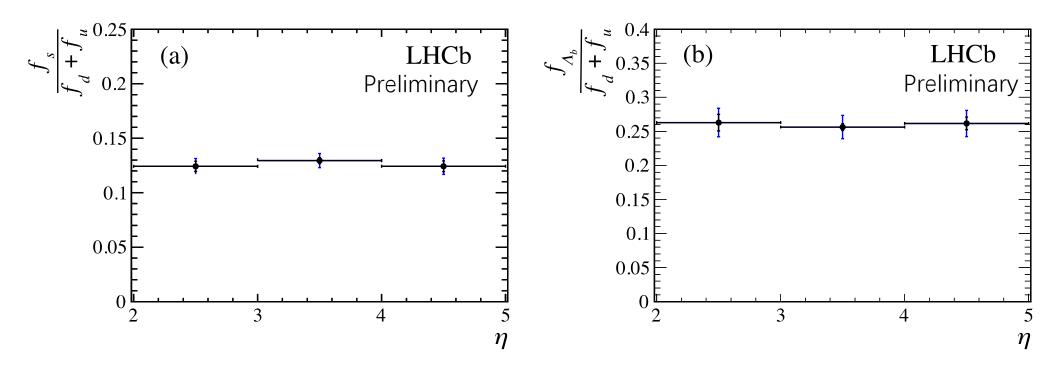
Linear

$$\geq \frac{f_{\Lambda_b^0}}{f_u + f_d}(p_{\rm T}) = A \left[(7.93 \pm 1.41) \cdot 10^{-2} + e^{(-1.022 \pm 0.047) + (-0.107 \pm 0.002)} p_{\rm T} / \text{GeV} \right]$$

Exponential

η dependence

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\triangleright No η dependence of the ratios is visible

Average fragmentation fractions

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Preliminary results

$$\frac{f_s}{f_u + f_d} = 0.122 \pm 0.006$$

$$\frac{f_{A_b^0}}{f_u + f_d} = 0.259 \pm 0.018$$

Kinematic region:

$$4 < p_T(H_b) < 25 \text{ GeV}/c$$

 $2 < \eta < 5$

- Statistical and systematic uncertainties combined
- > Systematic uncertainty dominates

LHCb 7 TeV result:

$$\frac{f_S}{f_u + f_d} = 0.128 \pm 0.010$$

[LHCb, JHEP 04 (2013) 001]

\mathcal{Z}_b^- production ratio

LHCb-PAPER-2018-047

arXiv:1901.07075, submitted to PRD

Introduction

- The *b*-hadron fragmentation fractions (f_u , f_d , f_s and $f_{\rm baryon}$) are available at the *Z* resonance and at $p\bar{p}$ collisions
- ➤ Complete measurements of *b*-hadron production fractions at the LHC do not exist yet
- > To achieve this, measurements of other b baryons are needed

			PDG 2018
$f_u + f_d + f_s + f_{\text{baryon}} = 1$	b hadron	Fraction at Z[%]	Fraction at $\overline{p}p$ [%]
$f_{\text{baryon}} = f_{\Lambda_h^0} + f_{\Xi_h^0} + f_{\Xi_h^-} + f_{\Omega_h^-}$	B^{+}, B^{0}	41.2 ± 0.8	34.0 ± 2.1
	B_s^0	8.8 ± 1.3	10.1 ± 1.5
$= f_{\Lambda_b^0} \left(1 + 2 \frac{f_{\Xi_b^-}}{f_{\Lambda_b^0}} + \frac{f_{\Omega_b^-}}{f_{\Lambda_b^0}} \right)$	$\frac{b}{a}$ baryons	8.9 ± 1.2	21.8 ± 4.7

DDC 2010

Measurement of $f_{\mathcal{Z}_h^{0(-)}}/f_{\varLambda_b^0}$

- ightharpoonup The best way is to measure $f_{\Xi_h^{0(-)}}/f_d$ using $\Xi_b^{0(-)} o \Xi_c^{+(0)} \mu^- \bar{\nu}_\mu X$ decays
 - Limited knowledge of absolute BRs of \mathcal{Z}_c^0 decays [Belle, arXiv:1811.09738]
 - No absolute BRs of \mathcal{Z}_c^+ decays available
 - Precision measurements should be feasible at Belle II
- \triangleright Alternative way is to measure $f_{\Xi_b^-}/f_{\Lambda_h^0}$ using the SU(3) related decays

$$\Lambda_b^0 \to J/\psi \Lambda$$
 and $\Xi_b^- \to J/\psi \Xi^-$

- - Uncertainty around 30%

• SU(3) symmetry
$$\Rightarrow$$
 the partial width ratio:
$$\frac{\Gamma(\Xi_b^- \to J/\psi\Xi^-)}{\Gamma(\Lambda_b^0 \to J/\psi\Lambda)} = \frac{3}{2} \text{ [M. Savage et al, NPB326 (1989) 15]}$$
• Uncertainty around 30%

$$R \equiv \frac{f_{\Xi_b^-}}{f_{\Lambda_b^0}} \frac{\mathcal{B}(\Xi_b^- \to J/\psi \, \Xi^-)}{\mathcal{B}(\Lambda_b^0 \to J/\psi \, \Lambda)} = \frac{f_{\Xi_b^-}}{f_{\Lambda_b^0}} \frac{\Gamma(\Xi_b^- \to J/\psi \, \Xi^-)}{\Gamma(\Lambda_b^0 \to J/\psi \, \Lambda)} \frac{\tau_{\Xi_b^-}}{\tau_{\Lambda_b^0}} = \frac{N(\Xi_b^- \to J/\psi \, \Xi^-)}{N(\Lambda_b^0 \to J/\psi \, \Lambda)} \frac{\epsilon_{\Lambda_b^0}}{\epsilon_{\Xi_b^-}}$$

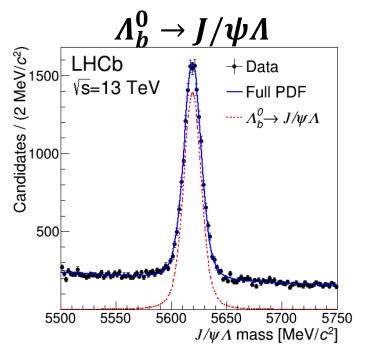
Known (theo. + exp.)

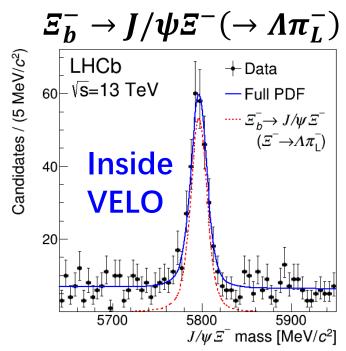
Measurable

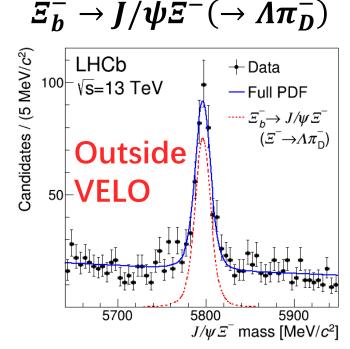
Data sample and mass fits

LHCb-PAPER-2018-047 arXiv:1901.07075

- Use Run1 (7&8 TeV) and 2016 (13 TeV) data
- \gt $E^- \rightarrow \Lambda \pi^-$ decays either inside VELO or outside VELO
- Signal: sum of two Crystal-Ball functions Background: exponential







Results: \mathcal{E}_b^- mass

LHCb-PAPER-2018-047 arXiv:1901.07075

- > Simultaneous fit to mass distributions of six subsamples:
 - $\Lambda_b^0 \rightarrow J/\psi \Lambda$: Run1 (7&8 TeV) and 2016 (13 TeV)
 - $\mathcal{E}_b^- \to J/\psi \mathcal{E}^-(\to \Lambda \pi_L^-)$: Run1 (7&8 TeV) and 2016 (13 TeV)
 - $\mathcal{E}_b^- \to J/\psi \mathcal{E}^-(\to \Lambda \pi_D^-)$: Run1 (7&8 TeV) and 2016 (13 TeV)
- > Mass difference $\delta m \equiv m(\Xi_b^-) m(\Lambda_b^0)$ and Ξ_b^- mass with systematic uncertainties taken into account

$$\delta m = 177.30 \pm 0.39 \pm 0.15 \,\text{MeV}/c^2,$$

 $m(\Xi_b^-) = 5796.70 \pm 0.39 \pm 0.15 \pm 0.17 \,\text{MeV}/c^2$

The most precise determination of \mathcal{Z}_b^- mass

Consistent with previous most precision result:

$$\delta m = 178.36 \pm 0.46 \pm 0.16 \, \mathrm{MeV}/c^2$$
 [LHCb, PRL113 (2014) 242002]

Results: \mathcal{E}_b^- production asymmetry

LHCb-PAPER-2018-047 arXiv:1901.07075

 \triangleright Repeat the fit by splitting into baryon (Ξ_b^-) and antibaryon (Ξ_b^+)

$$A_{\text{prod}}(\Xi_b^-) - A_{\text{prod}}(\Lambda_b^0) = \alpha(\Xi_b^-) - \alpha(\Lambda_b^0) - A_{\text{det}}(\pi^-)$$

Previous measurements

$$A_{\text{prod}}(\Lambda_b^0) = (2.4 \pm 1.4 \pm 0.9)\%$$

LHCb, Phys.Lett. B774 (2017) 139 LHCb, Phys.ReV. D91 (2015) 054022 LHCb, Chin.Phys.C40 (2016) 011001 Raw yield asymmetry from fits

Previous measurements consistent with zero

LHCb, Phys.Rev.Lett. 117 (2016) 061803 LHCb, JHEP 08 (2018) 008

$$A_{\text{prod}}(\Xi_b^-) = (1.1 \pm 5.6 \pm 1.9)\% \quad [\sqrt{s} = 7, 8 \text{ TeV}],$$

 $A_{\text{prod}}(\Xi_b^-) = (-3.9 \pm 4.9 \pm 2.5)\% \quad [\sqrt{s} = 13 \text{ TeV}].$

 $\boldsymbol{\mathcal{Z}}_{b}^{-}$ production asymmetry consistent with zero

Results: Production ratio $f_{\Xi_b^-}/f_{\Lambda_b^0}$

LHCb-PAPER-2018-047 arXiv:1901.07075

$$R \equiv \frac{f_{\Xi_b^-}}{f_{\Lambda_b^0}} \frac{\mathcal{B}(\Xi_b^- \to J/\psi \, \Xi^-)}{\mathcal{B}(\Lambda_b^0 \to J/\psi \, \Lambda)} = \frac{f_{\Xi_b^-}}{f_{\Lambda_b^0}} \frac{\Gamma(\Xi_b^- \to J/\psi \, \Xi^-)}{\Gamma(\Lambda_b^0 \to J/\psi \, \Lambda)} \frac{\tau_{\Xi_b^-}}{\tau_{\Lambda_b^0}} = \frac{N(\Xi_b^- \to J/\psi \, \Xi^-)}{N(\Lambda_b^0 \to J/\psi \, \Lambda)} \frac{\epsilon_{\Lambda_b^0}}{\epsilon_{\Xi_b^-}}$$

Known

Measurable

Signal yields + efficiencies + systematic uncertainties

$$R = (10.8 \pm 0.9 \pm 0.8) \times 10^{-2} \quad [\sqrt{s} = 7, 8 \text{ TeV}],$$

 $R = (13.1 \pm 1.1 \pm 1.0) \times 10^{-2} \quad [\sqrt{s} = 13 \text{ TeV}],$

$$\frac{f_{\Xi_b^-}}{f_{A_b^0}} = (6.7 \pm 0.5 \pm 0.5 \pm 2.0) \times 10^{-2} \quad [\sqrt{s} = 7, 8 \, \text{TeV}],$$

$$\frac{f_{\Xi_b^-}}{f_{A_b^0}} = (8.2 \pm 0.7 \pm 0.6 \pm 2.4) \times 10^{-2} \quad [\sqrt{s} = 13 \, \text{TeV}].$$
 (stat.) (syst.) (SU3 breaking)

Theoretical calculations:

$$\frac{f_{\mathcal{Z}_b^-}}{f_{\Lambda_b^0}} = (5.4 \pm 2.0) \times 10^{-2}$$

[H.-Y. Jiang et al, EPJC78 (2018) 224]

$$\frac{f_{\mathcal{Z}_b^-}}{f_{\Lambda_b^0}} = (6.5 \pm 2.0) \times 10^{-2}$$

[D. Wang, arXiv:1901.01776]

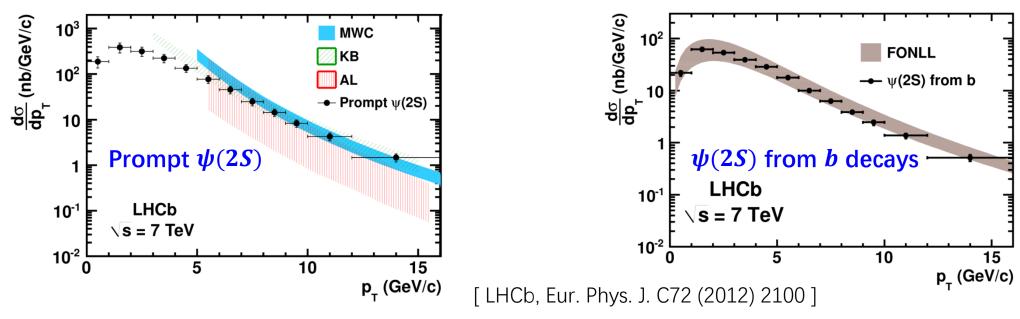
$\psi(2S)$ production cross-sections at 7 and 13 TeV

LHCb-PAPER-2018-049 in preparation



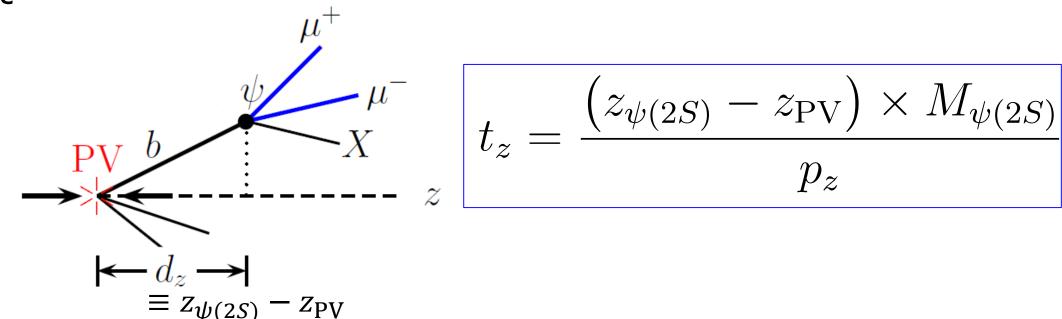
Introduction

- Study of heavy quarkonium production at hadron colliders provides important test to QCD models
 - Many models (NRQCD, CSM, COM, k_{T} factorization, FONLL, et al) available
 - Many measurements of heavy quarkonia performed at Tevatron and the LHC
- \triangleright Previous $\psi(2S)$ measurement in pp collisions at 7 TeV



Separation of prompt and non-prompt $\psi(2S)$

- \triangleright A fraction of $\psi(2S)$ comes from b-hadron decays
 - Prompt : direct (negligible feed-down contribution for $\psi(2S)$)
 - Non-prompt: from b-hadron decays (i.e. $\psi(2S)$ from b)
- Prompt and non-prompt separated by pseudo decay time in longitudinal or transverse direction



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Data sample and cross-section determination

- \geq 275 pb⁻¹ at 13 TeV (2015) and 614 pb⁻¹ at 7 TeV (2011)
 - Previous measurement at 7 TeV: 36 pb⁻¹ (2010)
- $\rightarrow \psi(2S) \rightarrow \mu^{+}\mu^{-}$ used owing to high efficiencies
- \triangleright Cross-section determined in each (p_T, y) bin

$\frac{\mathrm{d}^2\sigma}{\mathrm{d}y\mathrm{d}p_\mathrm{T}} = \underbrace{\frac{N(p_\mathrm{T},y)}{\Sigma_\mathrm{tot}(p_\mathrm{T},y)\times\mathcal{L}_\mathrm{int}\times\mathcal{B}\times\Delta y\times\Delta p_\mathrm{T}}}_{\mathcal{B}(\pmb{\psi}(\mathbf{2S})\to\pmb{\mu}^+\pmb{\mu}^-)}$ Efficiencies from simulation calibrated with data $\underbrace{R(\pmb{\psi}(\mathbf{2S})\to\pmb{\mu}^+\pmb{\mu}^-)}_{\text{Bin width luminosity}}$

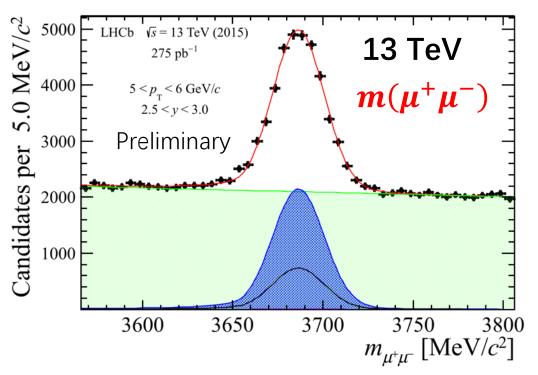
Signal yields

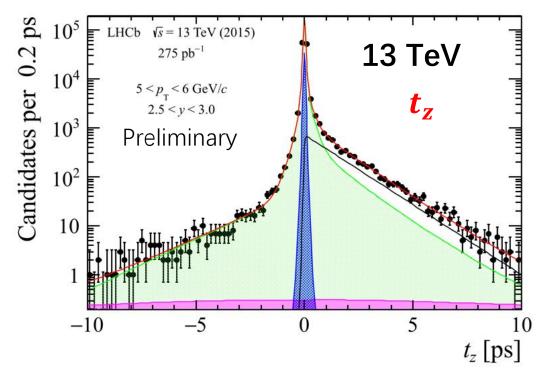
LHCb-PAPER-2018-049 in preparation

 \geq 2D fits to the $m(\mu^+\mu^-)$ and t_z distributions in each (p_T, y) bin

 $\rightarrow N_p(p_T, y)$: Signal yields of prompt $\psi(2S)$

 $N_b(p_T, y)$: Signal yields of $\psi(2S)$ from b





Results: Integrated cross-sections

LHCb-PAPER-2018-049 in preparation

Preliminary:

$$\sigma(\text{prompt } \psi(2S), 13 \text{ TeV}) = 1.430 \pm 0.005 \text{ (stat)} \pm 0.099 \text{ (syst) } \mu \text{b},$$

 $\sigma(\psi(2S)\text{-from-}b, 13 \text{ TeV}) = 0.426 \pm 0.002 \text{ (stat)} \pm 0.030 \text{ (syst) } \mu \text{b}.$

Kinematic region:

$$2 < p_{\rm T} < 20 ~{\rm GeV}/c$$
 and $2.0 < y < 4.5$

$$\sigma(\text{prompt } \psi(2S), 7 \text{ TeV}) = 0.471 \pm 0.001 \text{ (stat)} \pm 0.025 \text{ (syst) } \mu \text{b},$$

 $\sigma(\psi(2S)\text{-from-}b, 7 \text{ TeV}) = 0.126 \pm 0.001 \text{ (stat)} \pm 0.008 \text{ (syst) } \mu \text{b}.$

Kinematic region:

$$3.5 < p_{\rm T} < 14 ~{\rm GeV}/c$$
 and $2.0 < y < 4.5$

(due to tighter trigger selection)

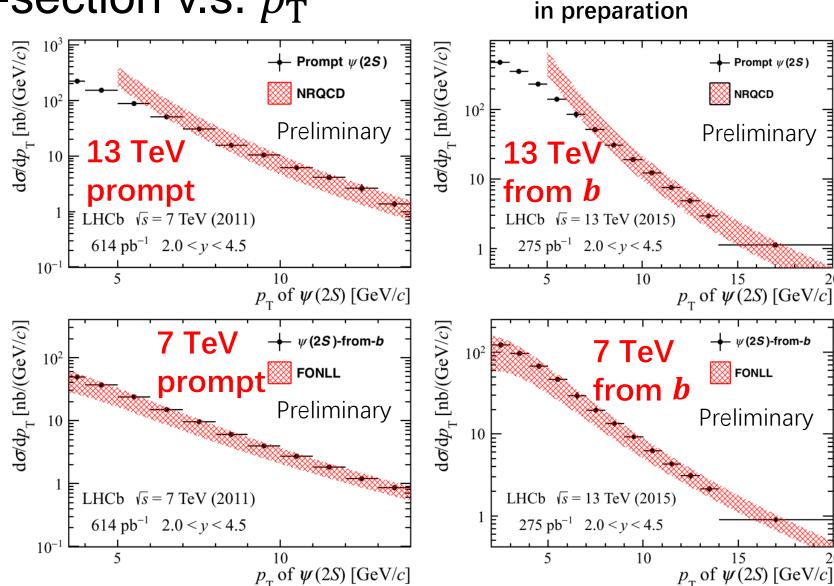
Results: cross-section v.s. $p_{\rm T}$

- Prompt results compared with NRQCD
- Non-prompt results compared with FONLL
- \succ Good agreement for high p_{T}

NRQCD:

[H.-S. Shao et al, JHEP 05 (2015) 103] FONLL:

- [M. Cacciari et al, JHEP 05 (1998) 007]
- [M. Cacciari et al, JHEP 10 (2012) 137]
- [M. Cacciari et al, EPJC75 (2015) 610]



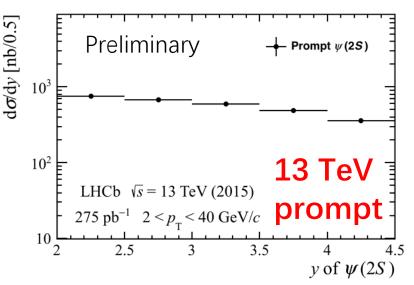
LHCb-PAPER-2018-049

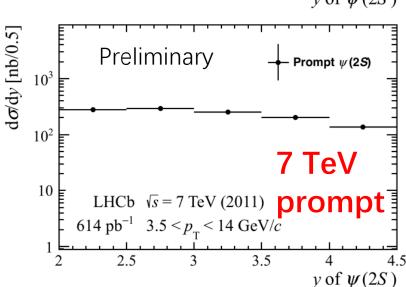
Results: cross-section v.s. y

- Non-prompt results compared with FONLL
- Good agreement

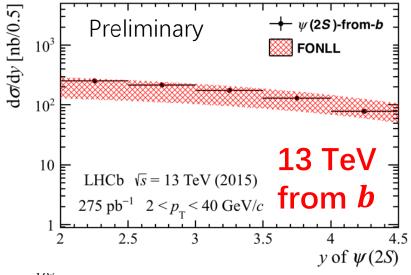
FONLL:

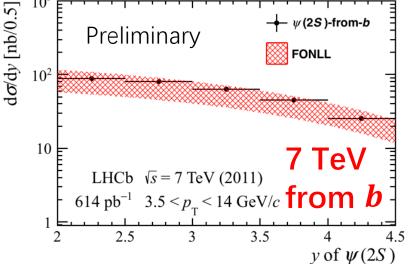
- [M. Cacciari et al, JHEP 05 (1998) 007]
- [M. Cacciari et al, JHEP 10 (2012) 137]
- [M. Cacciari et al, EPJC75 (2015) 610]





LHCb-PAPER-2018-049 in preparation





Results: 13 TeV and 7 TeV comparison

LHCb-PAPER-2018-049 in preparation

Most uncertainties cancel out in the ratios

$$R_{13/7} = \frac{\sigma(13 \text{ TeV})}{\sigma(7 \text{ TeV})}$$

- More precise test of theories
- Good agreement

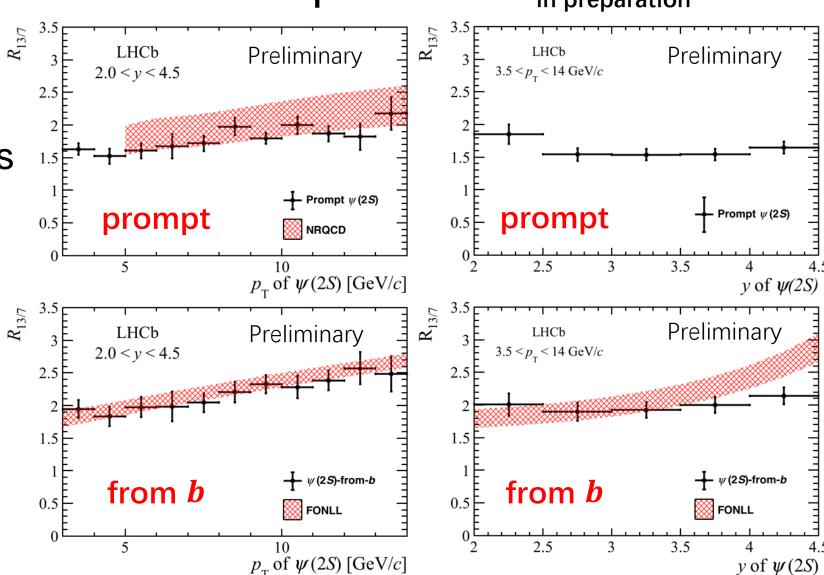
NRQCD:

[H.-S. Shao et al, JHEP 05 (2015) 103] FONLL:

[M. Cacciari et al, JHEP 05 (1998) 007]

[M. Cacciari et al, JHEP 10 (2012) 137]

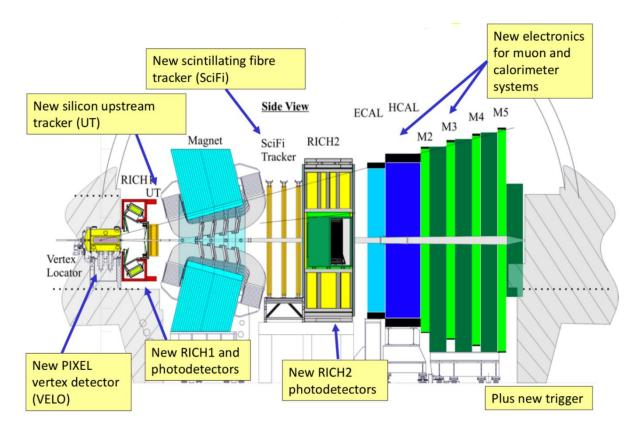
[M. Cacciari et al, EPJC75 (2015) 610]



Prospects

LHCb Upgrade (2019-2020)

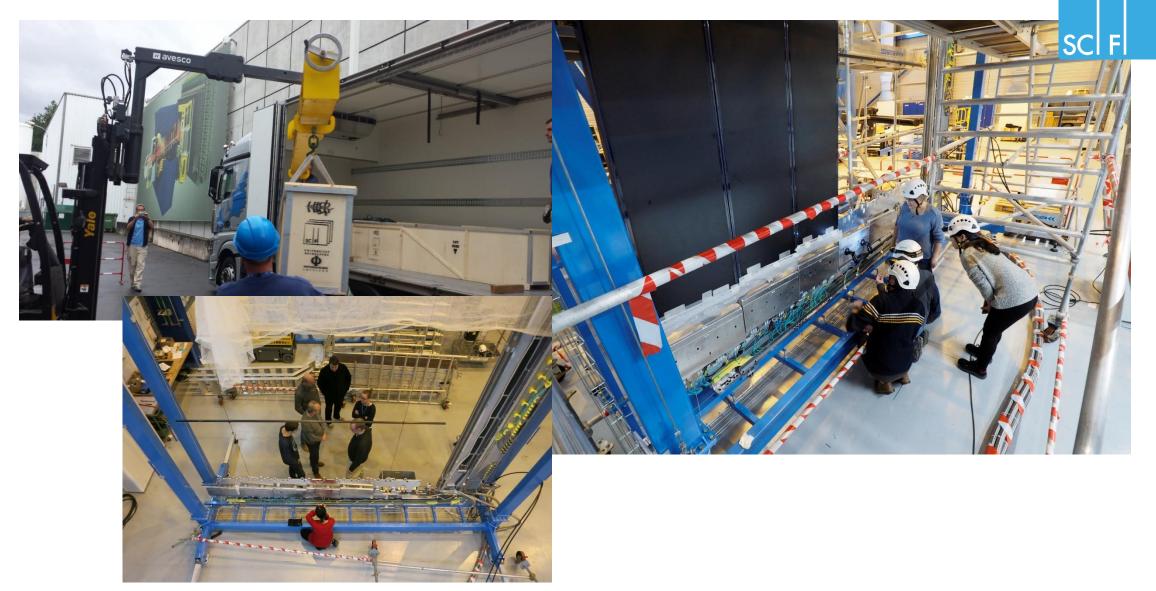
[LHCB-TDR-017]



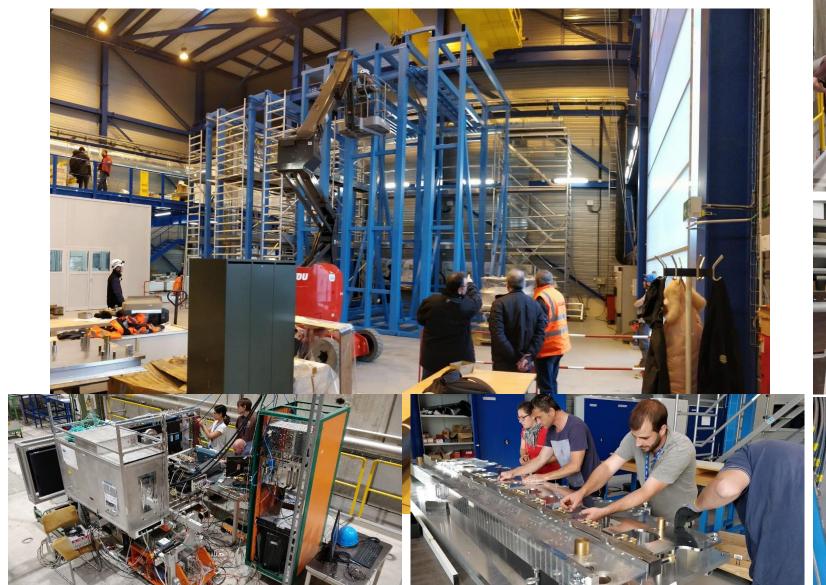
CERN-LHCC-2012-007

- > Increase luminosity to $2 \times 10^{33} \text{ cm}^{-2} \text{s}^{-1}$
 - 5 times larger than current maximum instantaneous luminosity
- All sub-detectors read out at 40 MHz for a full software trigger
 - Record with 10 GB/s
- All subdetector apart from muon and calorimeter systems will be fully replaced

Scintillating Fibre (SciFi) tracker installation



Scintillating Fibre (SciFi) tracker installation









Expected measurements

- ➤ A much larger sample of *b* and *c*-hadrons would be collected after LS2 with the Upgrade
- ➤ More precision measurements for SM tests and NP searches with heavy flavour, CKM, CPV, RD, spectroscopy, et al
- More heavy flavour production measurements could be performed or improved, e.g.,
 - Measurement of $f_{\Omega_b^-}/f_{\varLambda_b^0}$
 - Double heavy flavour production
 - $\Upsilon(nS) + \Upsilon(nS)$

$$f_u + f_d + f_s + f_{\text{baryon}} = 1$$

$$f_{\text{baryon}} = f_{\Lambda_b^0} + f_{\Xi_b^0} + f_{\Xi_b^-} + f_{\Omega_b^-}$$

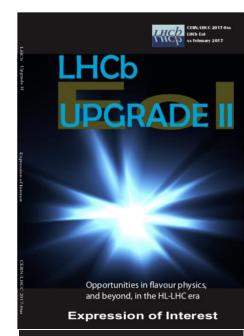
$$= f_{\Lambda_b^0} \left(1 + 2 \frac{f_{\Xi_b^-}}{f_{\Lambda_b^0}} + \frac{f_{\Omega_b^-}}{f_{\Lambda_b^0}} \right)$$

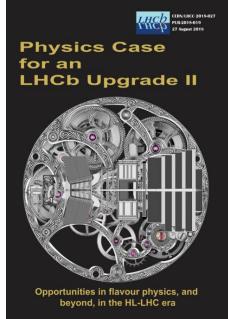
LHCb Upgrade 2

- Upgrade 2 proposed to take full profit of HL-LHC
 - $\mathcal{L} = 1 2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$, 10 times larger than Upgrade 1
 - Aiming at 300 fb⁻¹ after Run5



- Consolidate in LS3
- Major upgrade in LS4⁴
- ➤ EOI submitted in 2017 (CERN-LHCC-2017-003)
- > Physics document submitted in 2018 (arXiv:1808.08865)





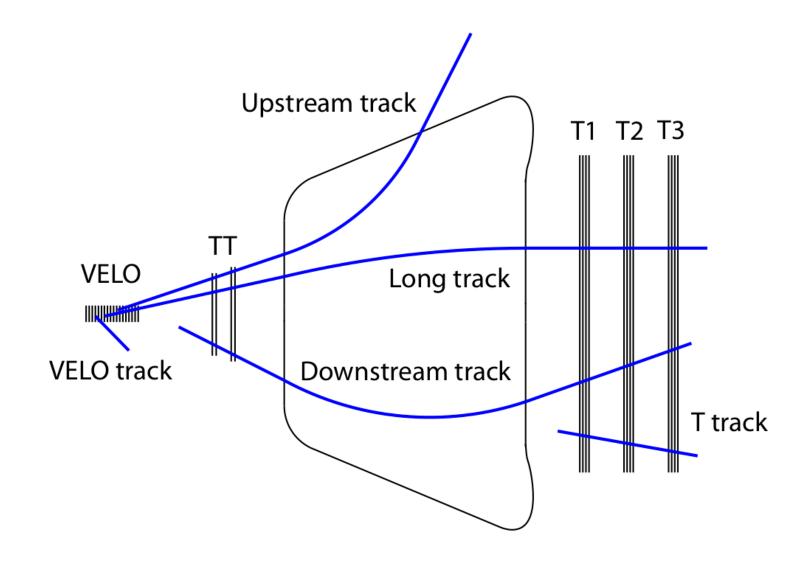
Summary

- LHCb has emphatically demonstrated its ability to perform important and unique measurements in various aspects
- New results of heavy flavour production are shown
 - *b*-hadron production fractions at 13 TeV
 - The mass and production rate of \mathcal{Z}_{h}^{-} baryons
 - $\psi(2S)$ production at 13 TeV and 7 TeV
- ➤ LHCb Upgrade I detector will be installed during LS2
 - Full software trigger at event rate ~30 MHz
 - Real time event reconstruction
 - Expect 23 fb^{-1} by 2025 and 50 fb^{-1} by 2029
- ➤ LHCb Upgrade II aiming at 300 fb⁻¹ with fully new detector to deepen our understanding of heavy flavour physics

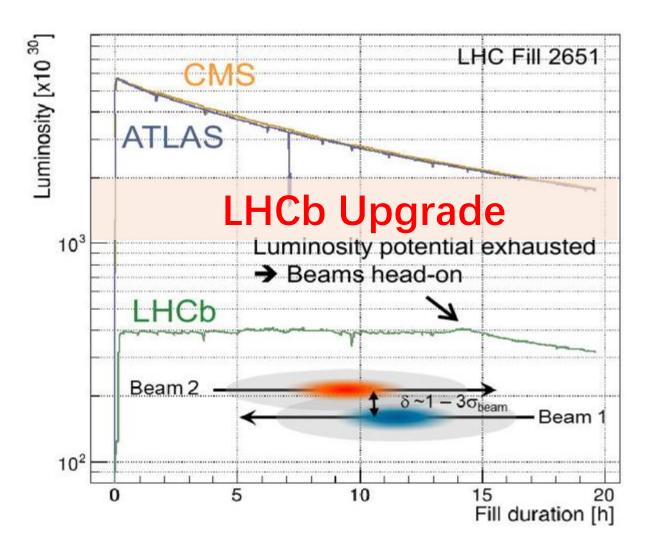


Backup slides

Track types for the LHCb Run I and II



How to increase the LHCb statistics significantly?



> LHCb up to LS2 (2018)

- Running at levelled luminosity of $\sim 4 \times 10^{32} \text{ cm}^{-2} \text{s}^{-1}$, pile-up~1
- First level hardware trigger running at event rate ~1 MHz
- Record ~12 kHz (0.6 GB/s)

> LHCb Upgrade I (2021-)

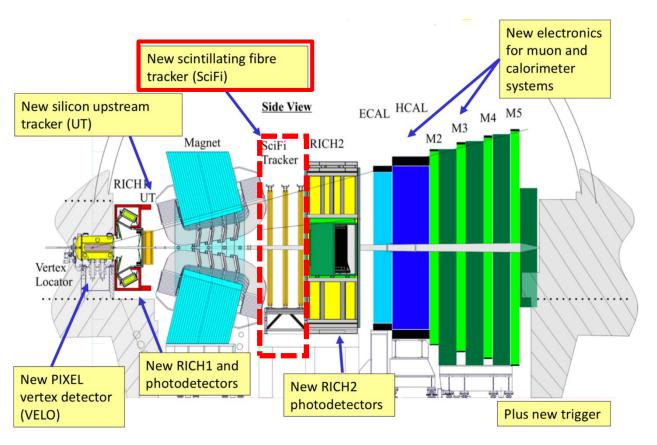
- Increase luminosity to a levelled 2×10^{33} cm⁻²s⁻¹, pile-up~5
- Run fully flexible and efficient software trigger up to 40 MHz
- Record with 10 GB/s

The most severe bottlenecks:

- Hardware trigger limited to ~ 1 MHz
- Tracking reconstruction

The LHCb Upgrade I detector

- > A complete new detector
 - All sub-detectors read out at 40 MHz for a fully software trigger



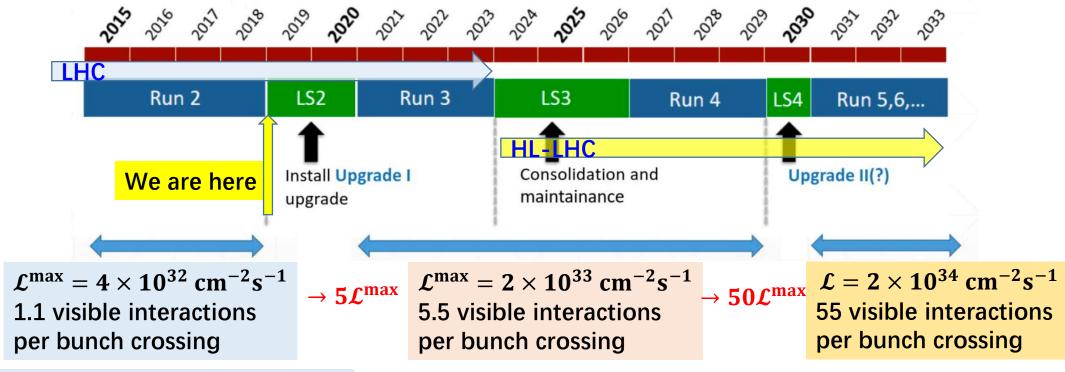
> Tracking system

- VELO: Silicon strip \rightarrow 55 × 55 μ m² PIXEL
- TT→ UT: Silicon strip → Silicon microstrip
- T1-T3→SciFi: Straw + silicon microstrip
 → Scintillating Fibre Tracker

> PID system

- RICH: HPD → MaPMT improved optics + mechanics
- ECAL/HCAL: remains the same ECAL inner modules replaced in LS3
- Muon: increased granularity

Plan of the LHC(b) upgrade



LHCb up to 2018 \rightarrow 9 fb⁻¹

- ✓ Demonstrated feasibility of high precision flavour physics at hadron colliders
- ➤ Find/rule out large sources of NP at the TeV scale

LHCb Upgrade I \rightarrow \geq 50 fb⁻¹

- ✓ Increase trigger efficiency
- ➤ Aim at experimental sensitivities comparable to theoretical uncertainties

LHCb Upgrade II \rightarrow \geq 300 fb⁻¹

- √ Take full profit of HL-LHC
- ➤ Physics document has been submitted to LHCC arXiv:1808.08865

Corrected yields of $B \rightarrow D\mu^-$

$$n_{\text{corr}}(B \to D^{0}\mu^{-}) = \frac{1}{\mathcal{B}(D^{0} \to K^{-}\pi^{+})\epsilon(B \to D^{0})} \times \left[n(D^{0}\mu^{-}) - n(D^{0}K^{+}\mu^{-}) \frac{\epsilon(\overline{B}_{s}^{0} \to D^{0})}{\epsilon(\overline{B}_{s}^{0} \to D^{0}K^{+})} - n(D^{0}p\mu^{-}) \frac{\epsilon(\Lambda_{b}^{0} \to D^{0})}{\epsilon(\Lambda_{b}^{0} \to D^{0}p)} \right]$$

$$n_{\text{corr}}(B \to D^{+}\mu^{-}) = \frac{1}{\epsilon(B \to D^{+})} \left[\frac{n(D^{+}\mu^{-})}{\mathcal{B}(D^{+} \to K^{-}\pi^{+}\pi^{+})} - \frac{n(D^{0}K^{+}\mu^{-})}{\mathcal{B}(D^{0} \to K^{-}\pi^{+})} \frac{\epsilon(\overline{B}_{s}^{0} \to D^{+})}{\epsilon(\overline{B}_{s}^{0} \to D^{0}K^{+})} - \frac{n(D^{0}p\mu^{-})}{\mathcal{B}(D^{0} \to K^{-}\pi^{+})} \frac{\epsilon(\Lambda_{b}^{0} \to D^{+})}{\epsilon(\Lambda_{b}^{0} \to D^{0}p)} \right].$$

Corrected yields of $\bar{B}_s^0 \to D\mu^-(K^+)$ and $\Lambda_b^0 \to D\mu^-$

$$n_{\text{corr}}(\overline{B}_{s}^{0} \to D_{s}^{+}\mu^{-}) = \frac{n(D_{s}^{+}\mu^{-})}{\mathcal{B}(D_{s}^{+} \to KK\pi)\epsilon(\overline{B}_{s}^{0} \to D_{s}^{+}\mu^{-})}$$
$$-N(\overline{B}^{0} + B^{-})\mathcal{B}(B \to D_{s}^{+}K)\frac{\epsilon(\overline{B} \to D_{s}^{+}K\mu^{-})}{\epsilon(\overline{B}_{s}^{0} \to D_{s}^{+}\mu^{-})}$$
$$n_{\text{corr}}(\overline{B}_{s}^{0} \to D^{0}K^{+}\mu^{-}) = 2\frac{n(D^{0}K\mu^{-})}{\mathcal{B}(D^{0} \to K\pi)\epsilon(\overline{B}_{s}^{0} \to D^{0}K\mu^{-})}$$

$$n_{\text{corr}}(\Lambda_b^0 \to D\mu^-) = \frac{n(\Lambda_c^+ \mu^-)}{\mathcal{B}(\Lambda_c^+ \to pK^-\pi^+)\epsilon(\Lambda_b^0 \to \Lambda_c^+)} + 2\frac{n(D^0 p\mu^-)}{\mathcal{B}(D^0 \to K^-\pi^+)\epsilon(\Lambda_b^0 \to D^0 p)}$$

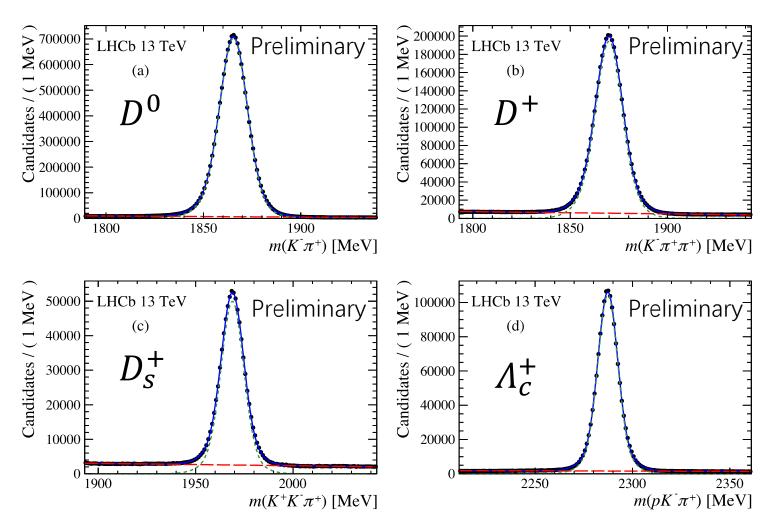
BRs of charmed hadron decays

Decay	\mathcal{B} $(\%)$	Source _
$D^0 \to K^-\pi^+$	3.93 ± 0.05	PDG average [3]
$D^+ \to K^- \pi^+ \pi^+$	9.22 ± 0.17	CLEO III [18]
$D_s^+ \to K^- K^+ \pi^+$	5.44 ± 0.18	CLEO III [18] PDG average [3]
$\Lambda_c^+ \to p K^- \pi^+$	6.23 ± 0.33	Weighted average of Belle
		and BES III results [19, 20]

Signal yields of charmed hadrons H_c

LHCb-PAPER-2018-050 in preparation

- ➤ Signal: two Gaussians
- ➤ Background: linear
- > Signal yields $N(D^{0}\mu^{-}\bar{\nu}_{\mu}X) = 13.8 \text{M}$ $N(D^{+}\mu^{-}\bar{\nu}_{\mu}X) = 4.28 \text{M}$ $N(D_{s}^{+}\mu^{-}\bar{\nu}_{\mu}X) = 0.85 \text{M}$ $N(\Lambda_{c}^{+}\mu^{-}\bar{\nu}_{\mu}X) = 1.75 \text{M}$
- > Yields obtained in bins of (p_T, η)



Cross-feed background with misidentified particles evaluated in (p_T, η) bins

A systematic check: f_+/f_0

LHCb-PAPER-2018-050 in preparation

- \blacktriangleright Measure the ratio of $D^0\mu^-\bar{\nu}_{\mu}X$ to $D^+\mu^-\bar{\nu}_{\mu}X$: f_+/f_0
- > Theoretical prediction

$$f_{+}/f_{0} = 0.387 \pm 0.012 \pm 0.026$$

[M. Rudolph, Int.J.Mod.Phys. A33 (2018) 1850176]

➤ Measured result: Preliminary

$$f_{+}/f_{0} = 0.359 \pm 0.006 \pm 0.009$$

No dependence of $p_{\rm T}$ and η is observed Consistent with theoretical prediction

Systematic of b-hadron production fraction

Preliminary

LHCb-PAPER-2018-050 in preparation

Source	Value (%)		
	$f_s/(f_u+f_d)$	$f_{A_b^0}/(f_u+f_d)$	f_+/f_0
Simulation	1.7	2.4	_
Backgrounds	0.9	0.3	
Cross-feeds	1.2	0.4	0.2
$\mathcal{B}(D^0 o K^-\pi^+)$	1.0	1.0	1.3
$\mathcal{B}(D^+ \to K^+\pi^-\pi^-)$	0.6	0.6	1.8
$\mathcal{B}(D_s^+ \to K^+ K^- \pi^+)$	3.3	_	
$\mathcal{B}(\Lambda_c^+ \to pK^+\pi^-)$	_	5.3	_
Measured lifetime ratio	1.2	0.7	
Γ_{sl} correction	0.5	1.5	_
Total	4.3	6.1	2.2

Systematic of \mathcal{Z}_b^- production ratio

LHCb-PAPER-2018-047 arXiv:1901.07075

Source	Value (%)	
Λ_b^0, Ξ_b^- polarization	3.0	
Signal and background shape	2.0	
Ξ_b^- production spectra	3.0	
π^- tracking efficiency	4.5	
Ξ^- mass resolution & non-resonant $\Lambda\pi^-$	3.0	
Ξ^- selections	1.4	
Ξ_b^- lifetime	0.5	
Simulated sample sizes	2.0	
Total	7.6	