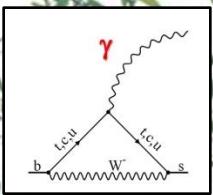




# Radiative decays of beauty hadrons at LHCb

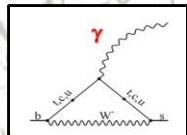
Vanya Belyaev

NIKHEF/Amsterdam & ITEP/Moscow



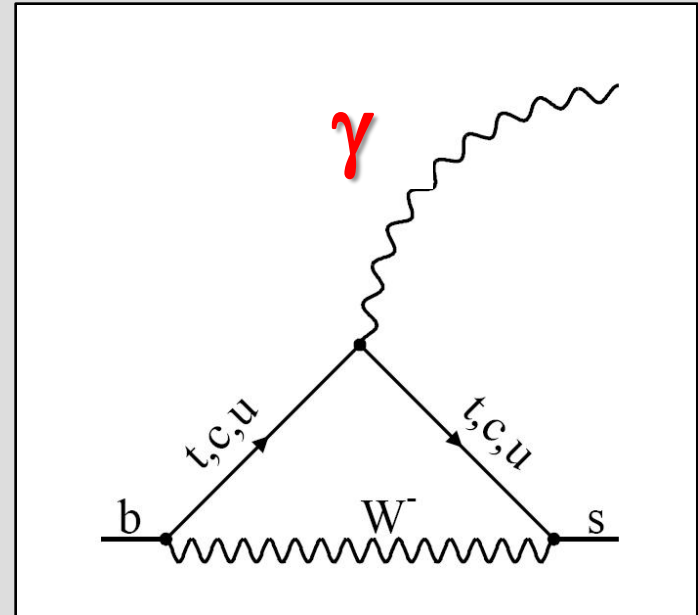


- (Exclusive) radiative penguin decays of beauty hadrons
  - Why they are interesting?
    - "The standard candle..."
- LHC collider and LHCb experiment
- Selected measurements with the radiative decays at LHCb:
  - Measurement of photon polarization in  $b \rightarrow s\gamma$  transitions:
    - Time dependent  $CP$ -asymmetry in  $B_s \rightarrow \phi\gamma$
    - Angular correlations in decays of polarized baryons  $\Lambda_b \rightarrow \Lambda\gamma$
    - Dalitz-analysis of  $B^\pm \rightarrow \phi K^\pm\gamma$
  - Direct  $CP$ -asymmetry in  $B^0 \rightarrow K^{*0}\gamma$
  - Extra rare decays:  $b \rightarrow d\gamma$  transitions
  - ...
- Summary

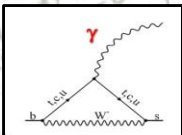


- Radiative decays in Standard Model occurs through "penguin" diagram:

- The simple example of *Flavor Changing Neutral Current FCNC*

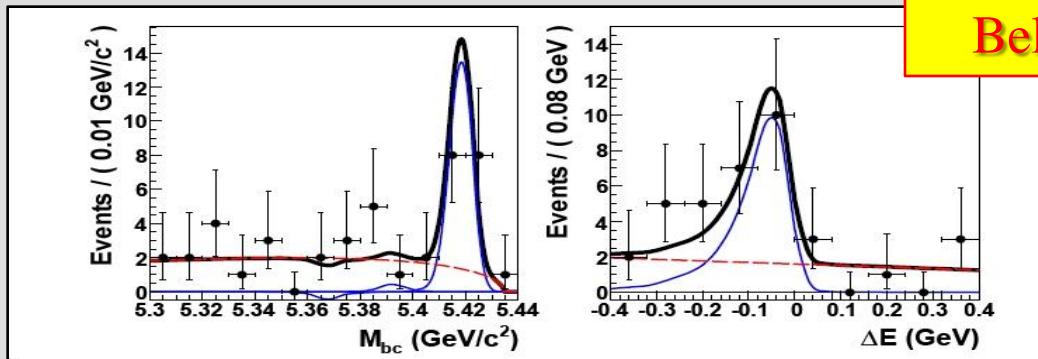


- Absent on the tree level
- Appears in the loop : "rare decay"

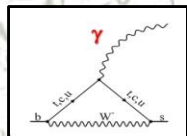




- Actually they are not so rare:
  - $Br(B^0 \rightarrow K^{*0} \gamma) = (40.1 \pm 2.0) 10^{-6}$  PDG'08
- First measured by CLEO'93, and later studies in great details by BaBar and Belle
- Recently also the radiative decays of  $B_s$ 
  - $Br(B_s \rightarrow \phi \gamma) = (57^{+18}_{-12} \text{ } ^{+12}_{-11}) 10^{-6}$  Belle'08



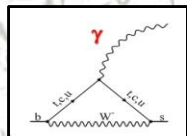
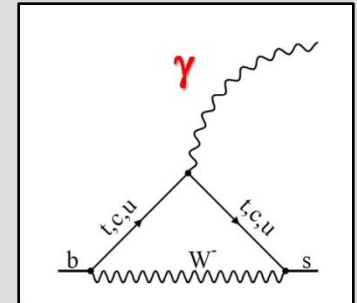
Belle:  $O(1 B_s \rightarrow \phi \gamma)$ /day at Y(5S)





# Why penguins are attractive?

- Very clear picture in standard model
  - 1-diagram dominance
  - Reliable theoretical description at (N)NLO allows the numerically precise predictions
    - the prerequisite for addressing NP
- "Loop":
  - New Physics contribution can be comparable and even dominating to (small) SM amplitudes
  - NP appears not only in modifications of Br, but also in asymmetries and the angular effects
    - "Sensitive also to spin structure of NP"

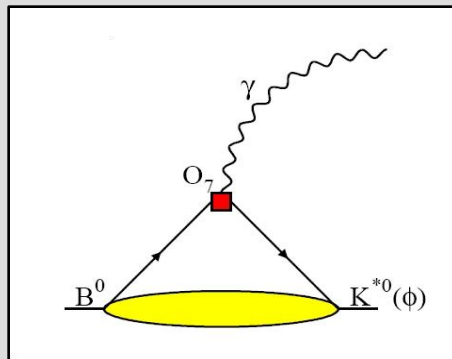


$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \sum_{p=u,c} \lambda_p^{(q)} \left( C_1(\mu) Q_1^p(\mu) + C_2(\mu) Q_2^p(\mu) + \sum_{i=3}^8 C_i(\mu) Q_i(\mu) \right),$$

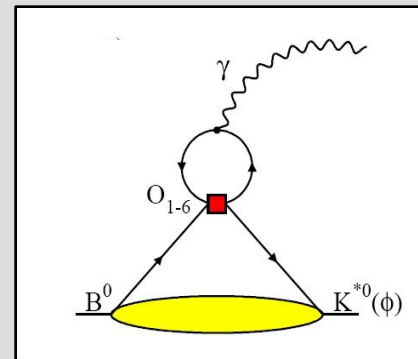
• **At LO.**

$$C_7^{\text{eff}} = C_7 + \dots$$

$$Q_7 = -\frac{e^2}{8\pi^2} \bar{m}_b(\mu) \bar{q} \sigma^{\mu\nu} (1 + \gamma_5) b F_{\mu\nu}$$

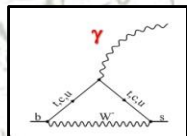


+

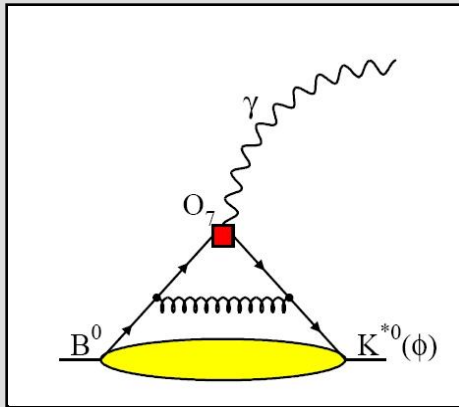


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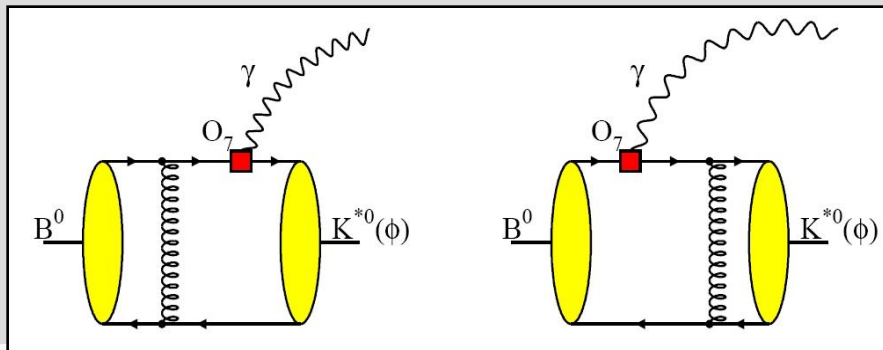
...



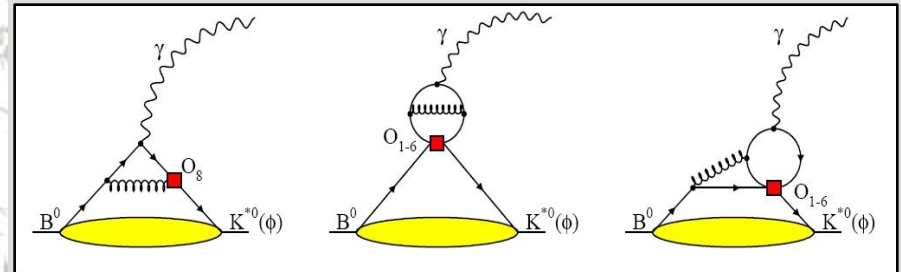
- + factorizable vertex graph



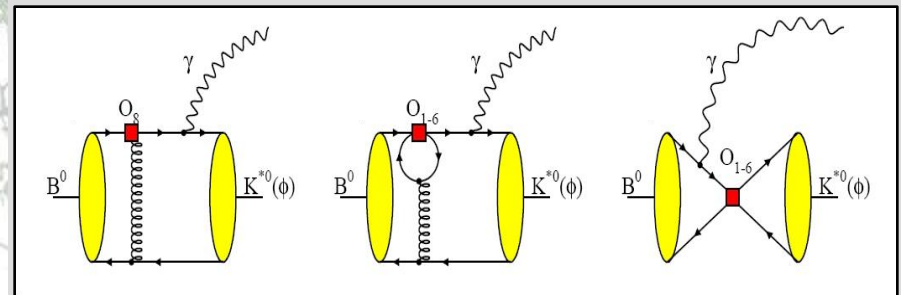
- + non-spectator



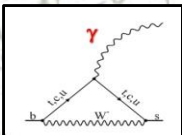
- Non-factorizable vertex:



- Non-spectator +...



... + annihilation



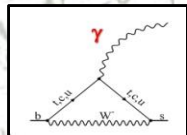


- 1-amplitude dominance
- Strong phase appears only in NLO
- Direct  $CP$ -asymmetries are small
  - ( $<1\%$ ) for  $b \rightarrow s\gamma$  &  $\mathcal{O}(10\%)$  for  $b \rightarrow d\gamma$
- *Isotopic asymmetries* due for non-spectator & annihilation  $\sim (C_3 + C_5/N_C)$  and suppressed by  $1/m_b$  and CKM-factors

$\mathcal{O}(5-12\%)$

$$A_{B \rightarrow K^* \gamma}^{\parallel} = \frac{\Gamma_{B_d^0 \rightarrow K^{*0} \gamma} - \Gamma_{B^- \rightarrow K^{*-} \gamma}}{\Gamma_{B_d^0 \rightarrow K^{*0} \gamma} + \Gamma_{B^- \rightarrow K^{*-} \gamma}}$$

- Q7 produces left-polarized photons
  - $A_R/A_L \sim m_s/m_b$
  - Mixing  $CP$ -asymmetries are vanished
    - $B \rightarrow f^{CP} \gamma$  is not  $CP$  eigenstate!



*"Mixing asymmetries are vanished, but..."*

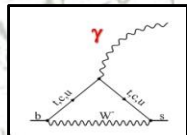
- $B \rightarrow f^{CP} \gamma$  is not  $CP$  eigenstate!  $\gamma_R/\gamma_L \approx m_s/m_b$
- Take it into account:

$$\Gamma(B_q(\bar{B}_q) \rightarrow f^{CP} \gamma) \propto e^{-\Gamma_q t} \left( \cosh \frac{\Delta\Gamma_q t}{2} - \mathcal{A}^\Delta \sinh \frac{\Delta\Gamma_q t}{2} \pm \right. \\ \left. \pm \mathcal{C} \cos \Delta m_q t \mp \mathcal{S} \sin \Delta m_q t \right)$$

- **SM:**

- $\mathcal{C} = 0$  direct  $CP$ -violation
- $\mathcal{S} = \sin 2\psi \sin \phi$
- $\mathcal{A}^\Delta = \sin 2\psi \cos \phi$

$$\tan \psi \equiv \left| \frac{A(\bar{B} \rightarrow f^{CP} \gamma_R)}{A(\bar{B} \rightarrow f^{CP} \gamma_L)} \right|$$



# $B^0$ vs $B_s$ : $\Delta\Gamma = 0$ vs $\Delta\Gamma_s/\Gamma_s \sim O(10\%)$

- $B^0$  :  $\Delta\Gamma \approx 0$

- $B_s$ :  $\Delta\Gamma_s/\Gamma_s \sim O(10\%)$

$$C = 0$$

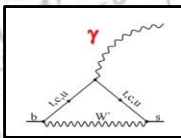
$$S = \sin 2\psi \sin \phi$$

$$A^\Delta = \sin 2\psi \cos \phi$$

- No sensitivity to  $A^\Delta$
- $\phi = 2\beta - \phi^{\text{peng}} \approx 2\beta$
- $\sin 2\psi = S / \sin 2\beta$

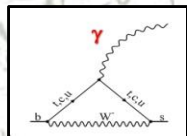
- Sensitive to  $A^\Delta$
- $\phi = 2\beta_s - \phi^{\text{peng}} \approx 0$ 
  - $S = 0$ , double smallness
- $\sin 2\psi = A^\Delta$

F.Muheim, Y.Xie & R.Zwicky, *Phys.Lett.B*664:174-179,2008





- In SM, the admixture of “right”-photon is “very small”  $\sim m_s/m_b$
- However the gluon emission could affect the pattern:
  - $\sim 10\%$ 
    - B. Grinstein, Y. Grossman, Z. Ligeti and D. Pirjol, Phys. Rev. D **71**,011504 (2005)
    - B. Grinstein and D. Pirjol, Phys. Rev. D **73**, 014013 (2006)
  - Hard gluon emission, pQCD  $\leq 3-4\%$ 
    - M. Matsumori and A.I. Sanda, Phys. Rev. D **73**, 114022 (2006)
  - Non-perturbative, slow gluon emission,  $c$ -loop +  $O_2$ , annihilation, etc... lightcone sum rules  $\sim \mathcal{O}(1\%)$ 
    - P. Ball and R. Zwicky, Phys. Lett. B **642**, 478 (2006)



- Sensitive to the spin structure & couplings of new particles
- Left-Right supersymmetry
  - minimal uSUSI
  - Enhancement:  $m_g/m_b$

$$\lambda_\gamma = \frac{|\mathcal{A}_R|^2 - |\mathcal{A}_L|^2}{|\mathcal{A}_R|^2 + |\mathcal{A}_L|^2}$$

$$\lambda_\gamma = \cos 2\psi$$

$$\lambda_\gamma(\text{SM}) = -1$$

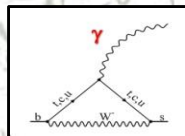
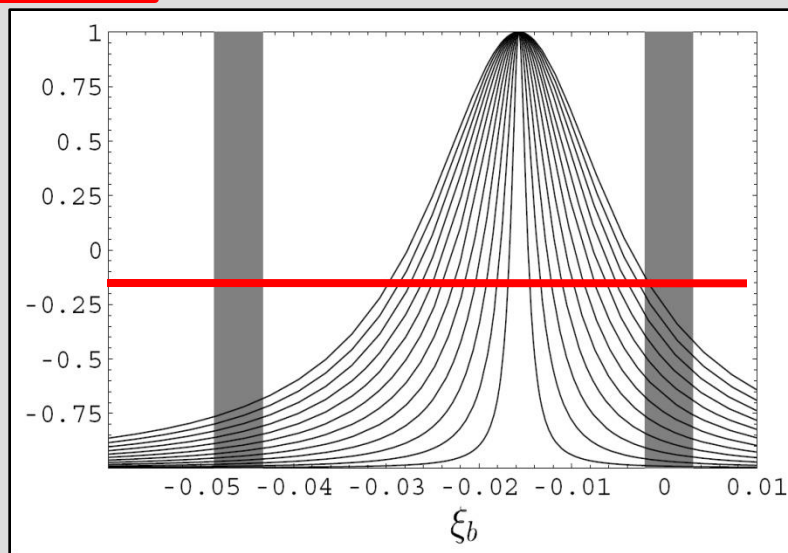
$$-1 \leq \lambda_\gamma \leq 1$$

L. Everett, G.L. Kane, S. Rigolin, L.-T. Wang and T.T. Wang, *JHEP* 0201, 022 (2002)

- Anomalous top-coupling:

$$-1 \leq \lambda_\gamma \leq -0.12$$

J.-P. Lee, *Phys. Rev. D* 69, 014017 (2004).

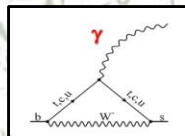


# How to measure $\sin 2\psi/\lambda_\gamma$ ?

- Time-dependent CP-asymmetry for  $B^0/B_s$  mesons
- Angular correlations (Dalitz-analysis)

$$\frac{d\Gamma(B \rightarrow \Phi\gamma)}{d\Omega} \propto (|\mathcal{A}_R|^2 + |\mathcal{A}_L|^2) + \lambda_\gamma (|\mathcal{A}_R|^2 - |\mathcal{A}_L|^2)$$

- $B \rightarrow K\pi\pi\gamma$  : many amplitudes, difficult analysis
  - M. Gronau, Y. Grossman, D. Pirjol and A. Ryd, *Phys. Rev. Lett.* **88**, 051802 (2002)
  - M. Gronau and D. Pirjol, *Phys. Rev. D* **66**, 054008 (2002)
- $B^\pm \rightarrow \phi K^\pm\gamma$ 
  - V.D. Orlovsky and V.I. Shevchenko, *Phys. Rev. D* **77**, 093003 (2008)
- Polarized b-baryon decays
  - Mannel, Recksiegel, *JPG: NPP* **24**, 979 (1998)
  - Hiller, Kagan, *PRD* **65**, 074038 (2002)
- Study for converted photons and  $B \rightarrow V l^+ l^-$  in the limit of  $m(l^+ l^-) \rightarrow 0$ 
  - Melikov, Nikitin, Simula, *PLB* **442**, 381 (1998)
  - Grossman, Pirjol, *JHEP* **06**, 029 (2000)







- Baryons are polarized
  - Polarization can be measured via  $\Lambda_b \rightarrow J/\psi \Lambda^0$
- exploit the angular correlations between initial and final state
- one can apply to higher resonances  $\Lambda(X) \rightarrow pK$ 
  - spin 5/2 is useless (lack of observables)
  - spin  $\frac{1}{2}$  need to be separated from 3/2
    - The measurement both photon and proton distributions are required



$$\frac{d\Gamma}{d \cos \theta_\gamma} \propto 1 - \alpha_\gamma P_B \cos \theta_\gamma$$

$$\frac{d\Gamma}{d \cos \theta_p} \propto 1 - \alpha_\gamma \alpha_p \cos \theta_p$$

Hiller, Kagan, PRD 65, 074038 (2002)

- $\Lambda_b \rightarrow \Lambda(\frac{1}{2}) \gamma$   
 $\cos \theta_\gamma$  the same  
 $\cos \theta_p$  flat



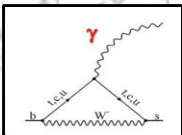
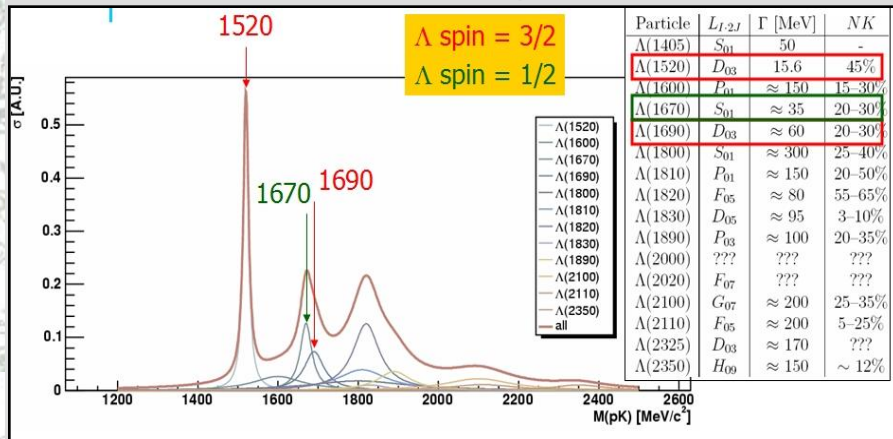
$$\frac{d\Gamma}{d \cos \theta_\gamma} \propto 1 - \alpha_{\gamma,3/2} P_B \cos \theta_\gamma$$

$$\frac{d\Gamma}{d \cos \theta_p} \propto 1 - \alpha_{p,3/2} \cos^2 \theta_p$$

$$\alpha_{\gamma,3/2} = \frac{1 - \eta}{1 + \eta} \alpha_\gamma$$

$$\alpha_{p,3/2} = \frac{\eta - 1}{\eta + \frac{1}{3}}$$

F. Legger and T. Schietinger, *Phys.Lett.B* 645, 204 (2007)

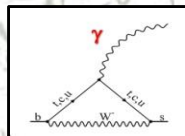
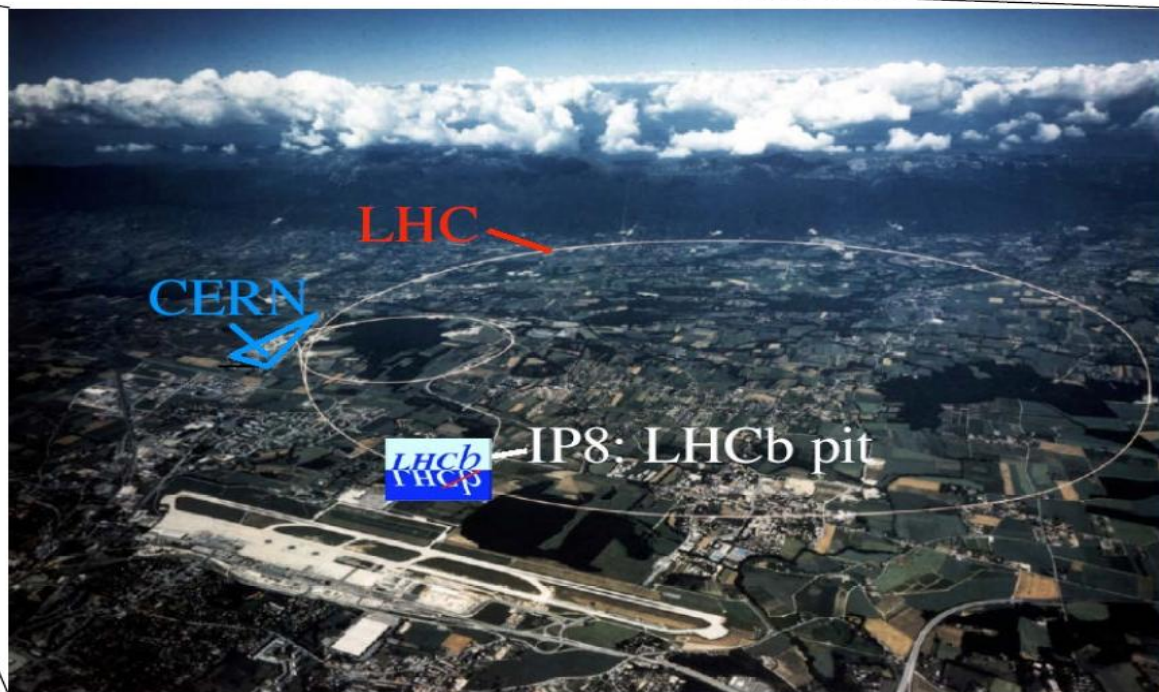




# LHCb Experiment

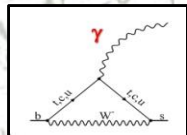
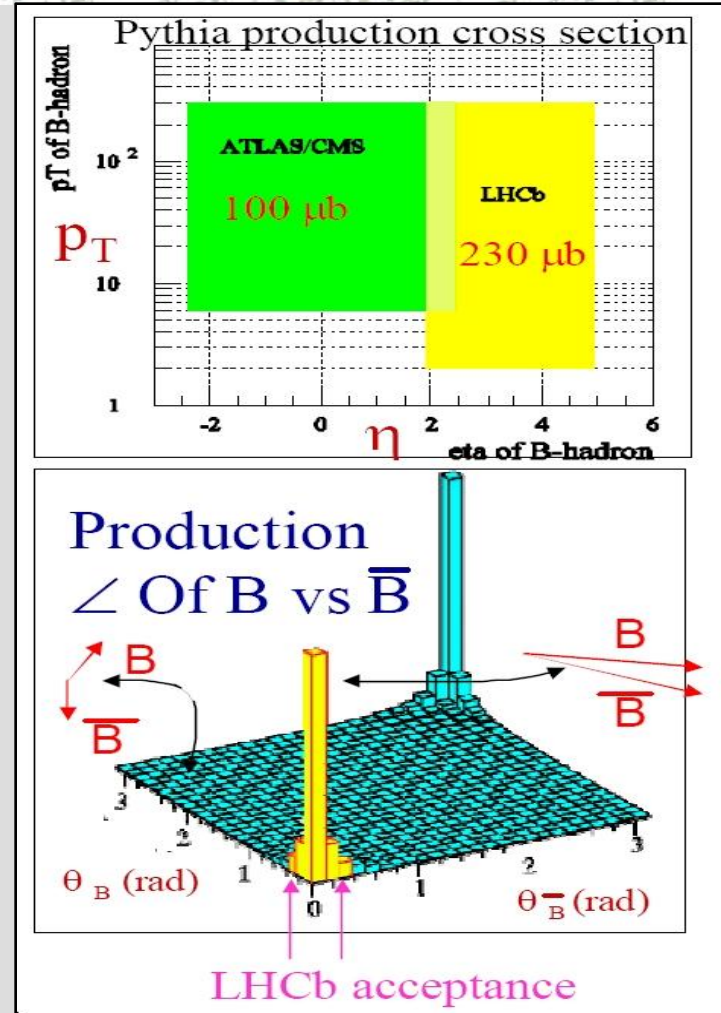


14 countries  
48 institutions  
~600 people



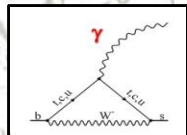
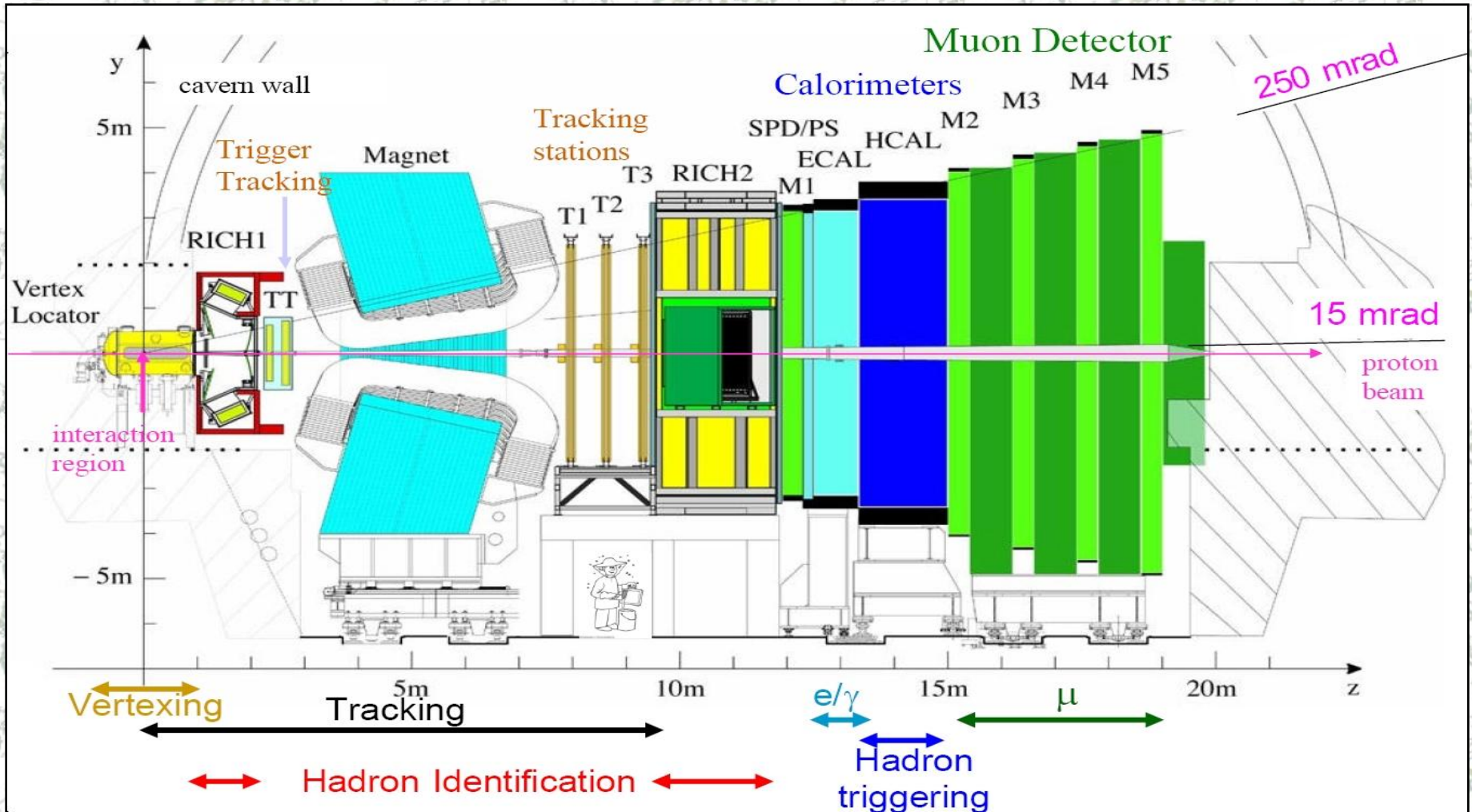


- In the forward region  $\sigma_b$  is large
  - $O(10^{12})$  B-hadrons per  $10^7\text{s}$  at  $L=2 \cdot 10^{32}$
  - Limited solid angle  $\rightarrow$  Limited cost
- Both B-hadrons in events are likely to be in acceptance
  - b-flavor tagging!
- B-hadrons are moving with large momentum  $\beta\gamma=O(20)$ ,  $l_B \sim 1\text{cm}$ 
  - Minimize the multiply scattering
  - Background rejection via detached vertex
  - Improved decay time resolution



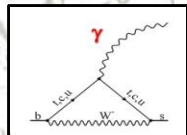


# LHCb Detector (sketch)





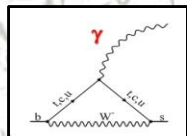
# LHCb Detector in Cavern





- Beauty particles:
  - $m_b \sim 5 \text{ GeV}/c^2$
  - $\beta\gamma c\tau \sim \mathcal{O}(1\text{cm})$
- Particles from B-decays:
  - Large  $p_T$
  - Large impact parameters
  - *Exploited in a smart way by L0/HLT triggers!*
- Background:
  - $b\bar{b}$ -production with at least one B within 400mrad cone

- We do have L0 trigger for energetic photons!
- Analyses are based on full detailed detector simulation with the realistic reconstruction chain
  - Common features:
- B-decay products do not point to reconstructed primary vertices
- Exclusively reconstructed B-candidate does point to primary vertex
- B-candidate is associated with primary vertex with minimal impact parameter (significance)





# Selection of $B_d \rightarrow K^{*0} \gamma$ and $B_s \rightarrow \phi \gamma$

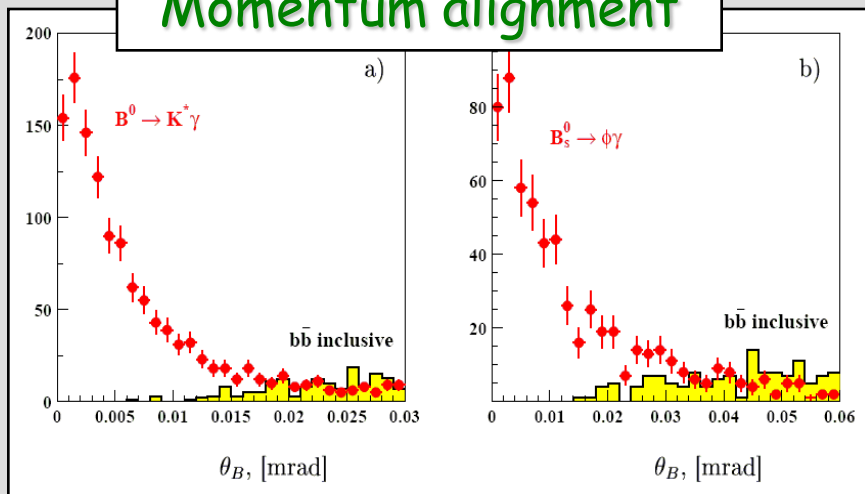
## Important factors

- Good hadron PID :
  - $K/\pi/p$  separation
- Good quality of 2-prong vertex fit

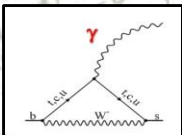
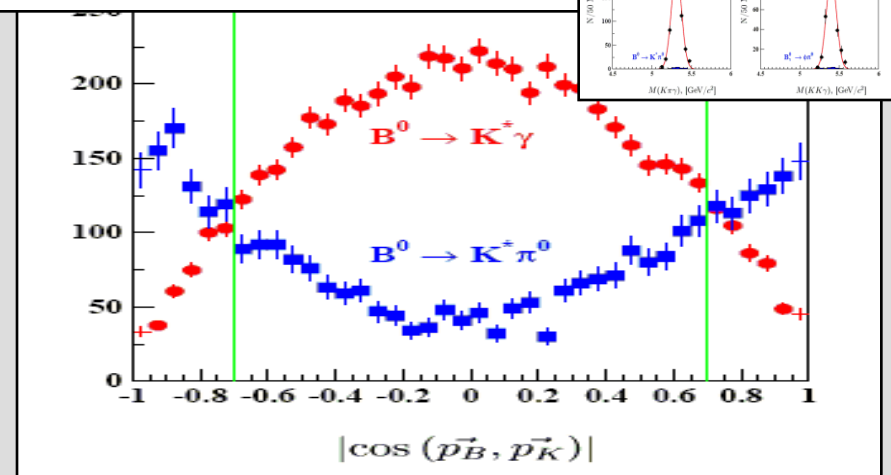
## Correlated feeddown from energetic $\pi^0 \rightarrow \gamma \gamma$

- $K^{*0} \pi^0 \rightarrow K^{*0} \gamma \gamma$
  - $\phi \pi^0 \rightarrow \phi \gamma \gamma$
- Different vector polarisation

## Momentum alignment



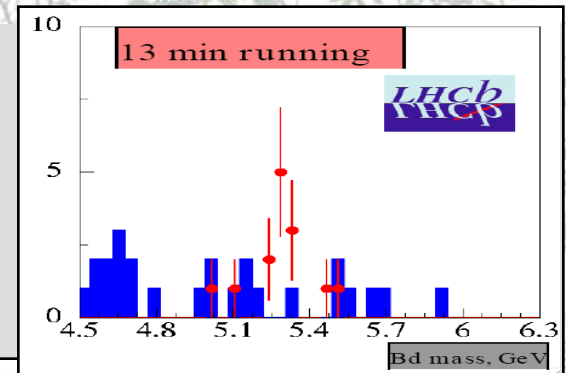
## Correlated feeddown



- Selection of other channels

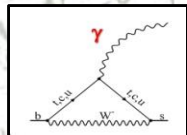
$$\Lambda_b \rightarrow \Lambda^0 \gamma, \Lambda_b \rightarrow (\Lambda^0(1670) \rightarrow pK^-) \gamma,$$

$$B^\pm \rightarrow (\phi \rightarrow K^+K^-) K^\pm \gamma, \dots \sim \text{the same}$$



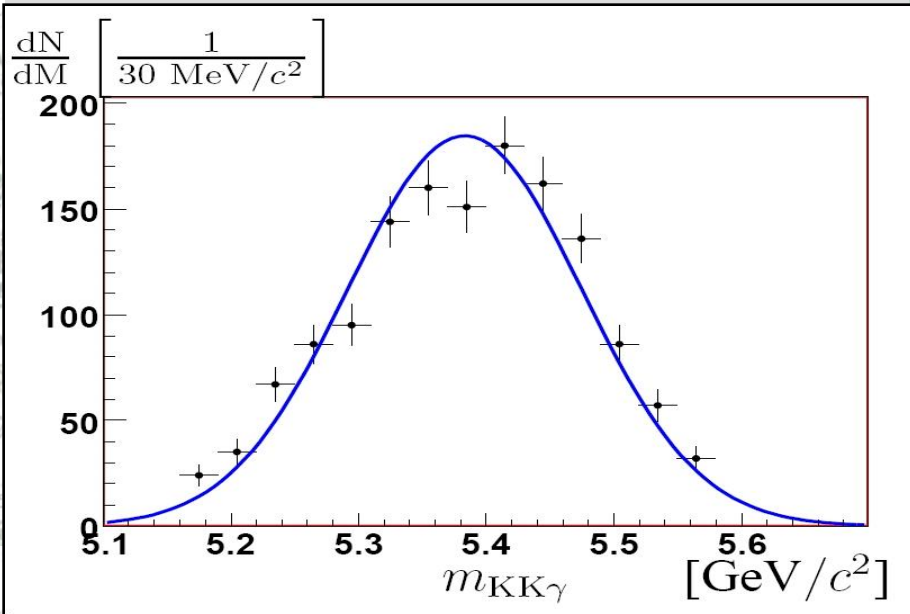
Decay Mode	#/ $2\text{fb}^{-1}$	$\mathcal{B}_{b\bar{b}}/S$
$B^0 \rightarrow K^{*0} \gamma$	$7 \cdot 10^4$	0.6-0.7
$B_s^0 \rightarrow \phi \gamma$	$1.1 \cdot 10^4$ <span style="border: 1px solid red; padding: 2px;"><math>\propto (1 B_s \rightarrow \phi \gamma) / \text{hour}</math></span>	$< 0.6 - 0.9$
$\Lambda_b \rightarrow \Lambda^0 \gamma$	750	$< 42$
$\Lambda_b \rightarrow \Lambda^0(1670) \gamma$	$2.5 \cdot 10^3$ <span style="font-size: small;"><math>\sim \Lambda^0(1690), \sim \frac{1}{2} \Lambda^0(1520)</math></span>	$< 18$
$B^+ \rightarrow \phi K^+ \gamma$	$\sim 7 \cdot 10^3$	$< 2 - 4$

- L. Shchutska, A. Golutvin and I. Belyaev, CERN-LHCb-2007-030
- F. Legger, CERN-LHCb-2006-012

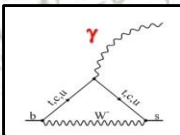
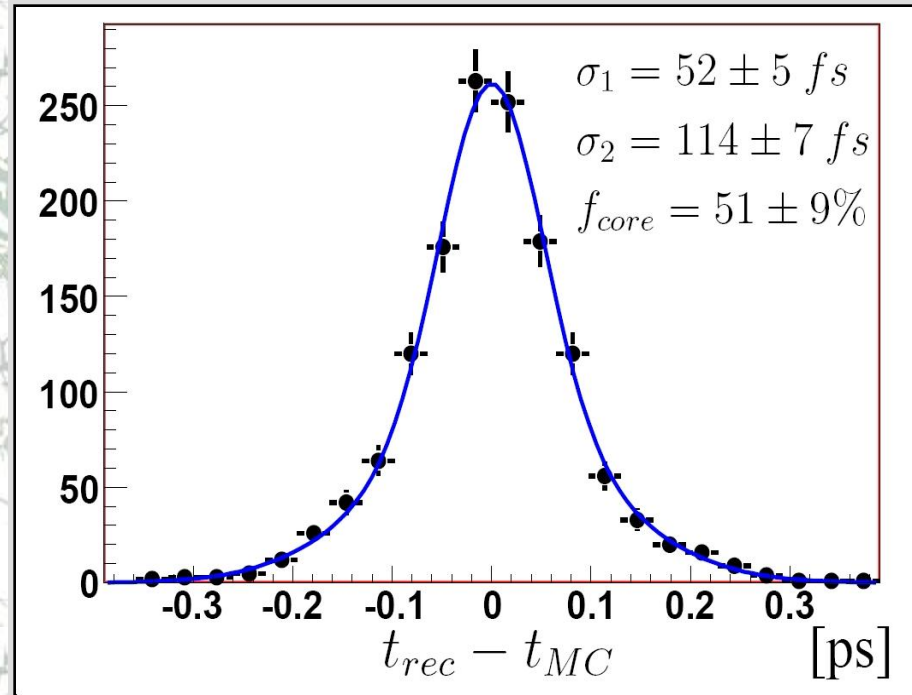


# What do we "know" about signal?

- Mass resolution
  - $\sim 90 \text{ MeV}/c^2$
  - Totally dominated by Ecal energy resolution



- Proper Life-time resolution
  - $\sigma \sim 78 \text{ fs}$



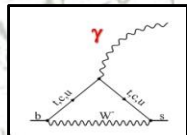
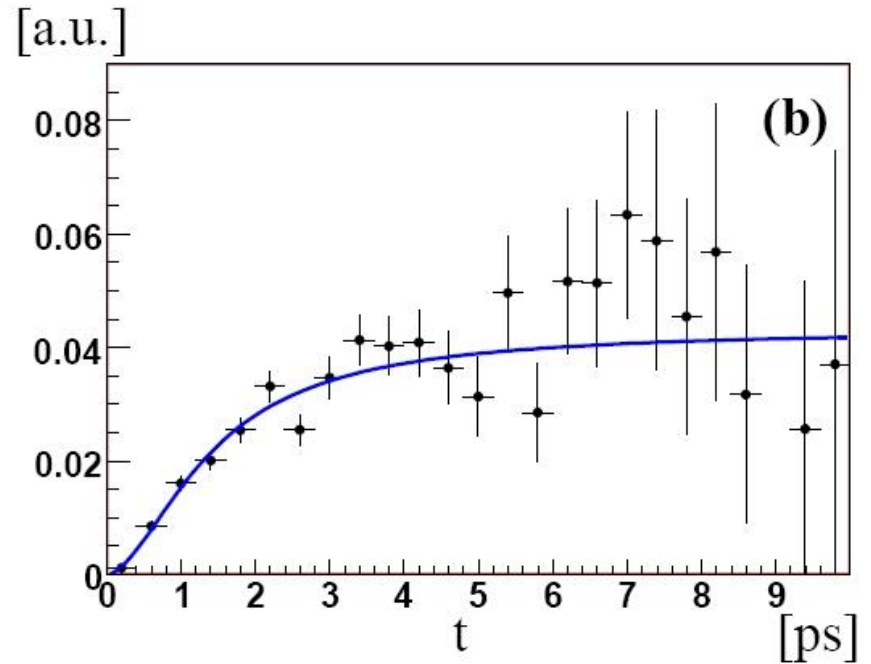
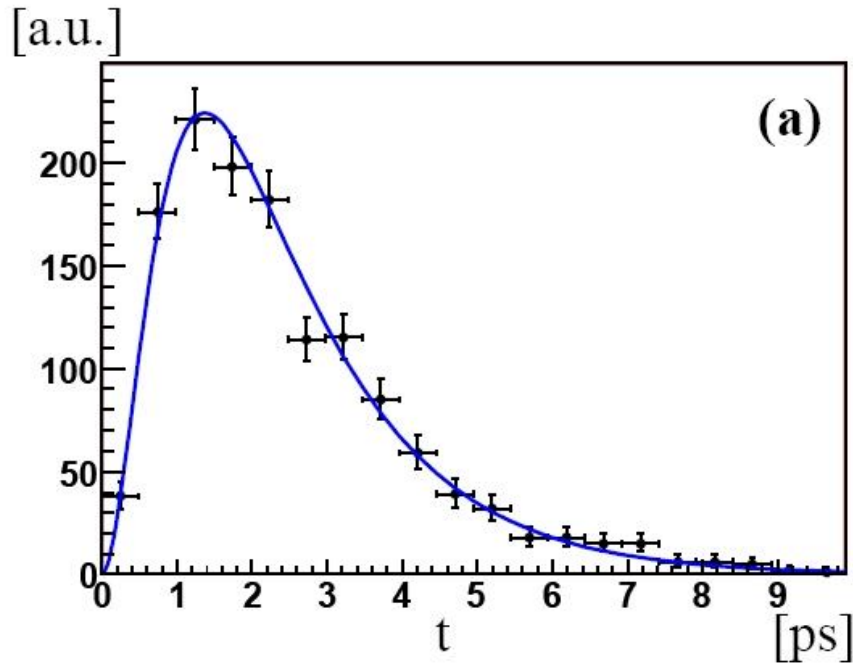


# (Signal) proper time acceptance

$$a = 0.74 \pm 0.09 \text{ ps}^{-1}$$

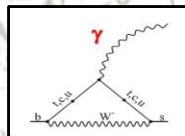
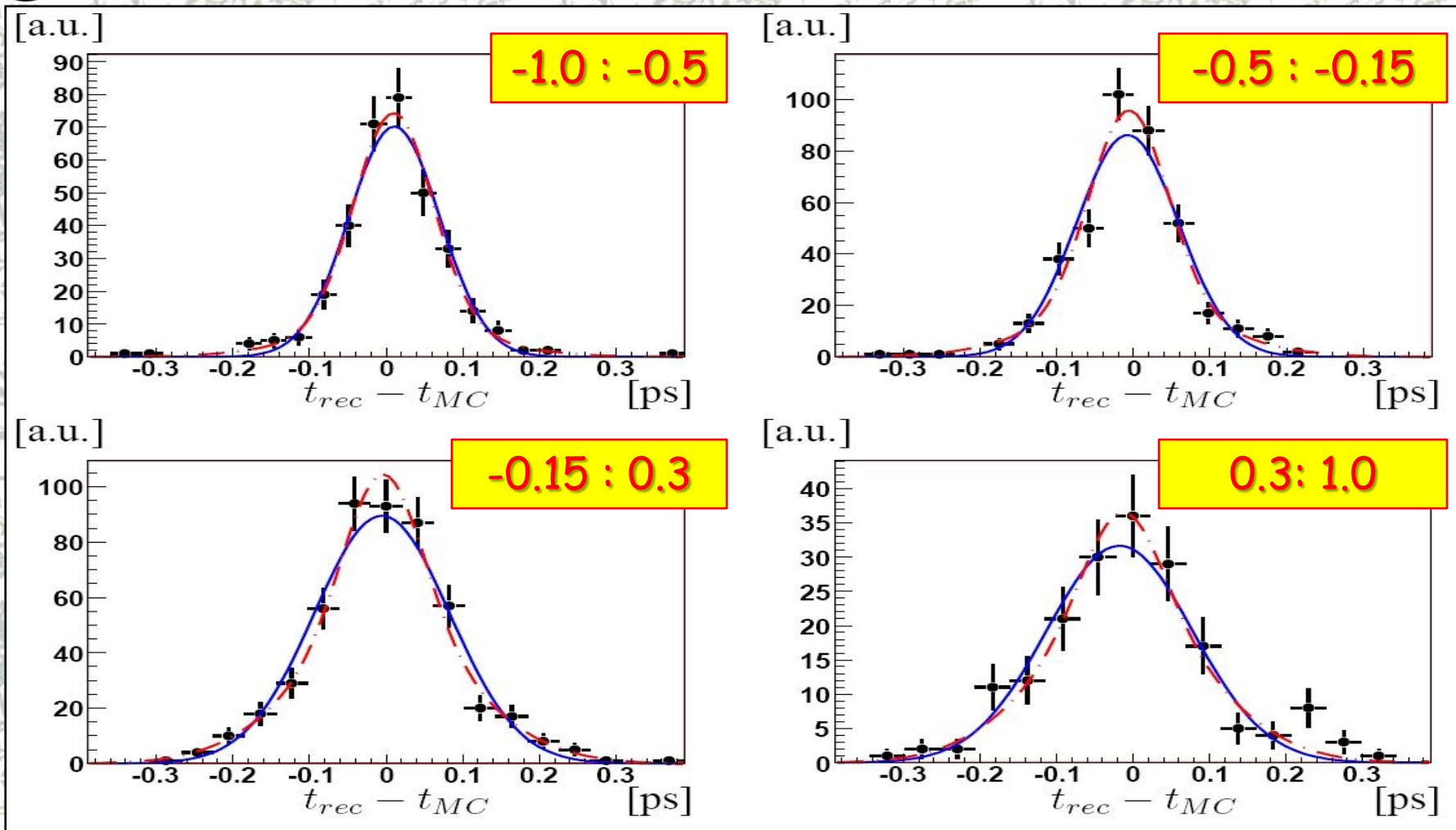
$$c = 1.86 \pm 0.15$$

$$\varepsilon_s(t) \propto \frac{(at)^c}{1 + (at)^c}$$





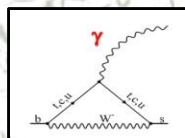
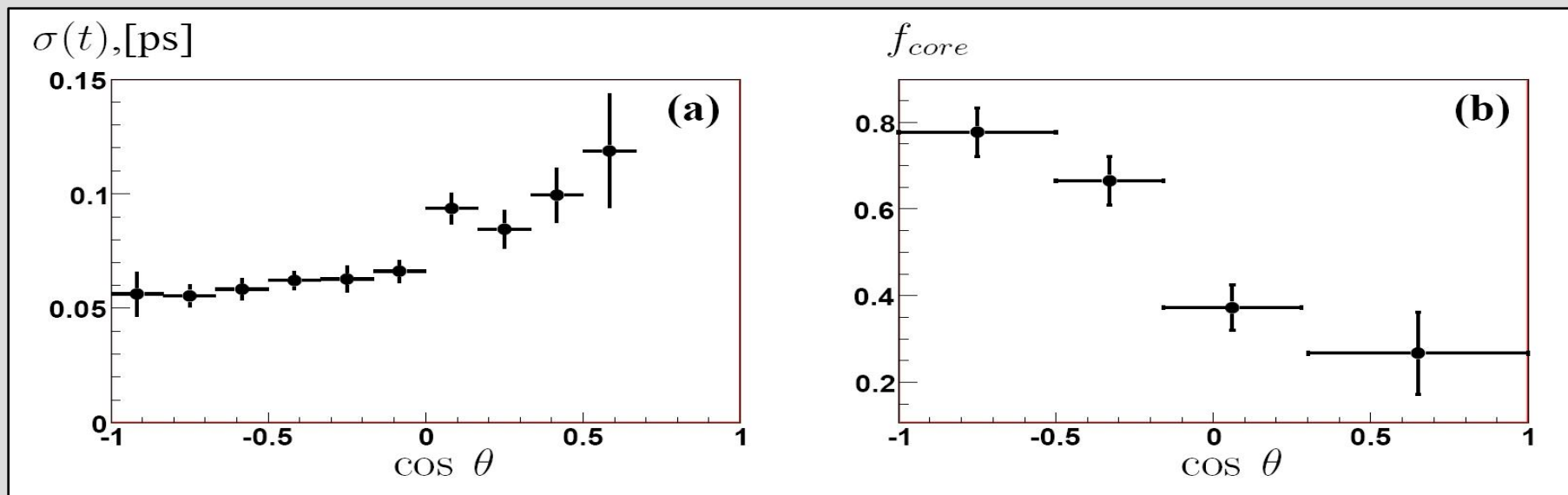
# Signal proper time resolution as function of $\cos\Theta$





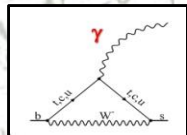
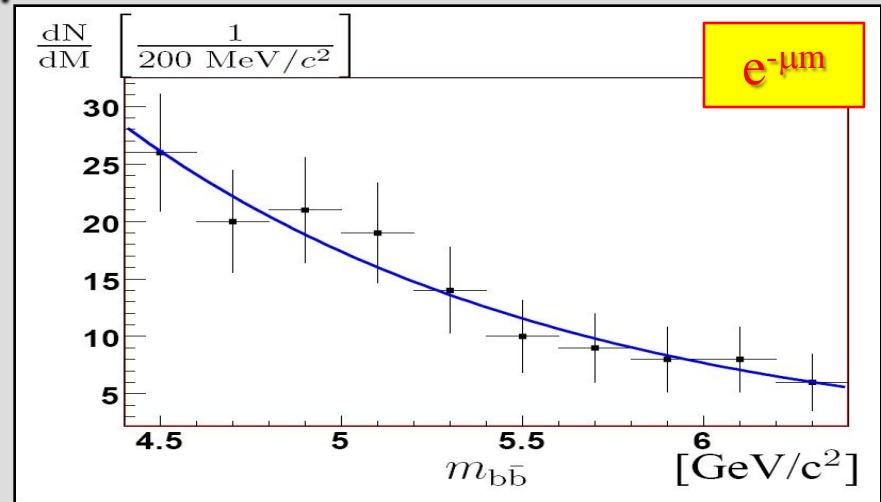
# Signal proper time resolution as function of $\cos\theta$

$\cos \theta$	mean <sub>1G</sub> , $f s$	$\sigma$ , $f s$	mean <sub>2G</sub> , $f s$	$f_{core}$
[-1,-0.5)	$10 \pm 3$	$59 \pm 3$	$10 \pm 3$	$0.78 \pm 0.06$
[-0.5,-0.15)	$-8 \pm 4$	$66 \pm 4$	$-6 \pm 4$	$0.66 \pm 0.06$
[-0.15,0.3)	$-5 \pm 4$	$88 \pm 4$	$-4 \pm 4$	$0.37 \pm 0.05$
[0.3,1]	$-18 \pm 8$	$96 \pm 7$	$-13 \pm 8$	$0.27 \pm 0.09$

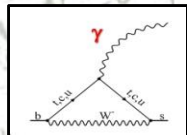
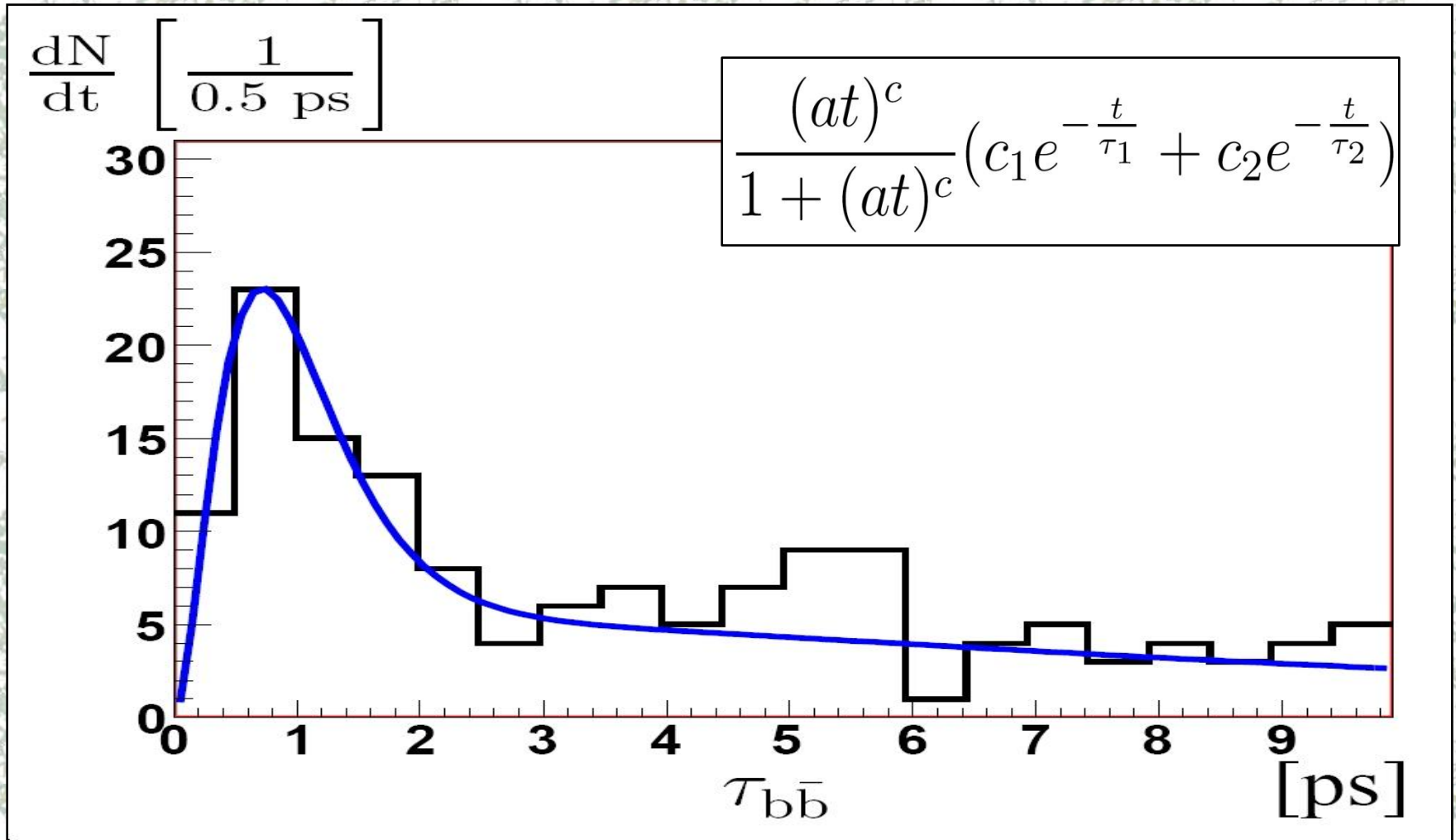




- Limited set of Monte Carlo statistics  $\sim (1-5)10^7$   
potentially dangerous "forward" bb-events
  - 15-20 minutes of LHC
- Many estimates are (MC) statistically limited and rely on some assumptions.
  - "loose cuts ..."



# Background "proper time" distribution

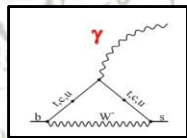
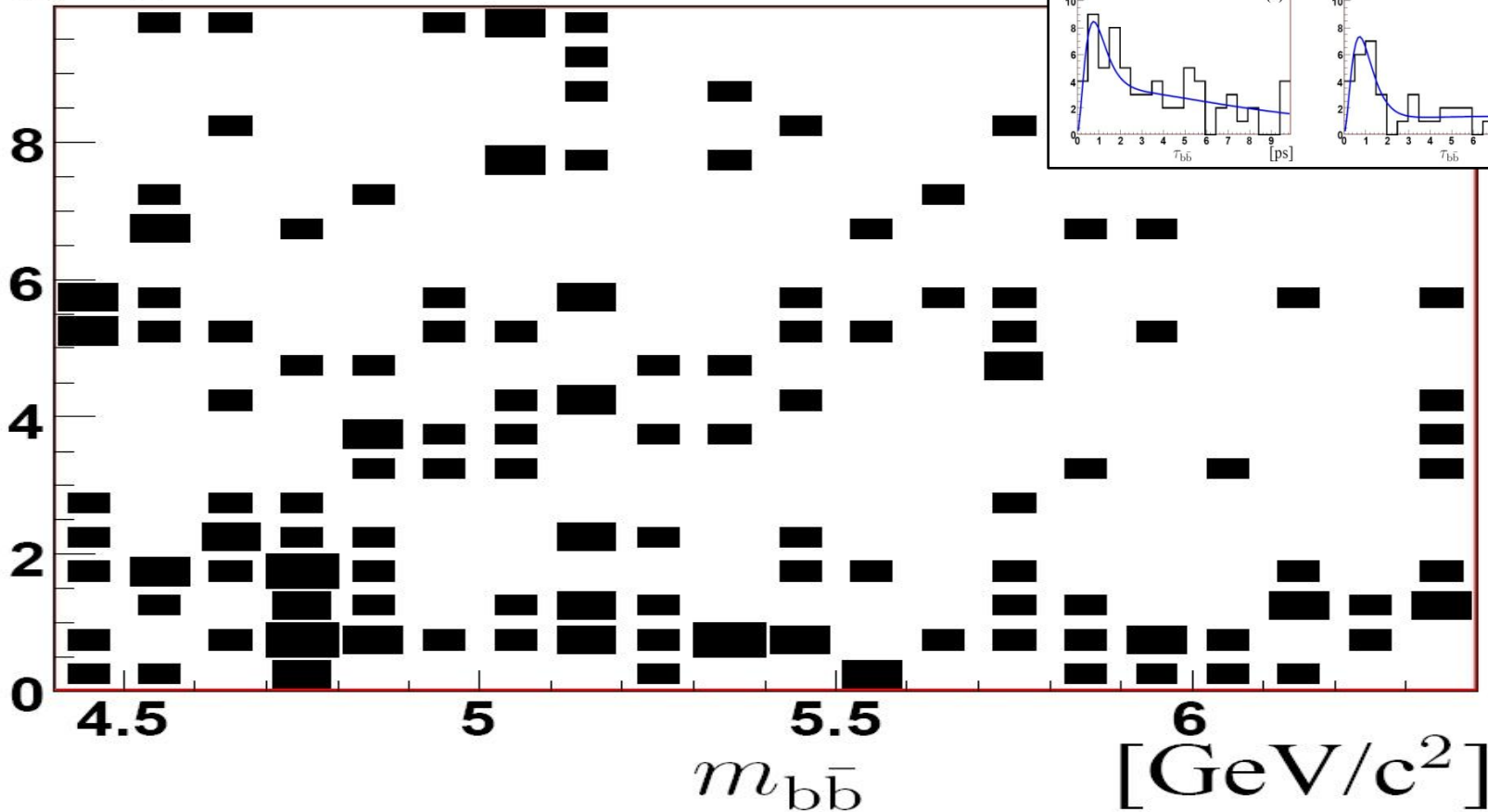




# B: "proper-time" vs mass

$$e^{-\mu m} \frac{(at)^c}{1 + (at)^c} \left( (\alpha_0 + \alpha_1 \Delta m) e^{-\frac{t}{\tau_1}} + (\beta_0 + \beta_1 \Delta m) e^{-\frac{t}{\tau_2}} \right)$$

$t, [\text{ps}]$



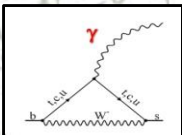


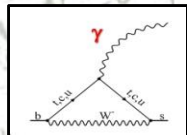
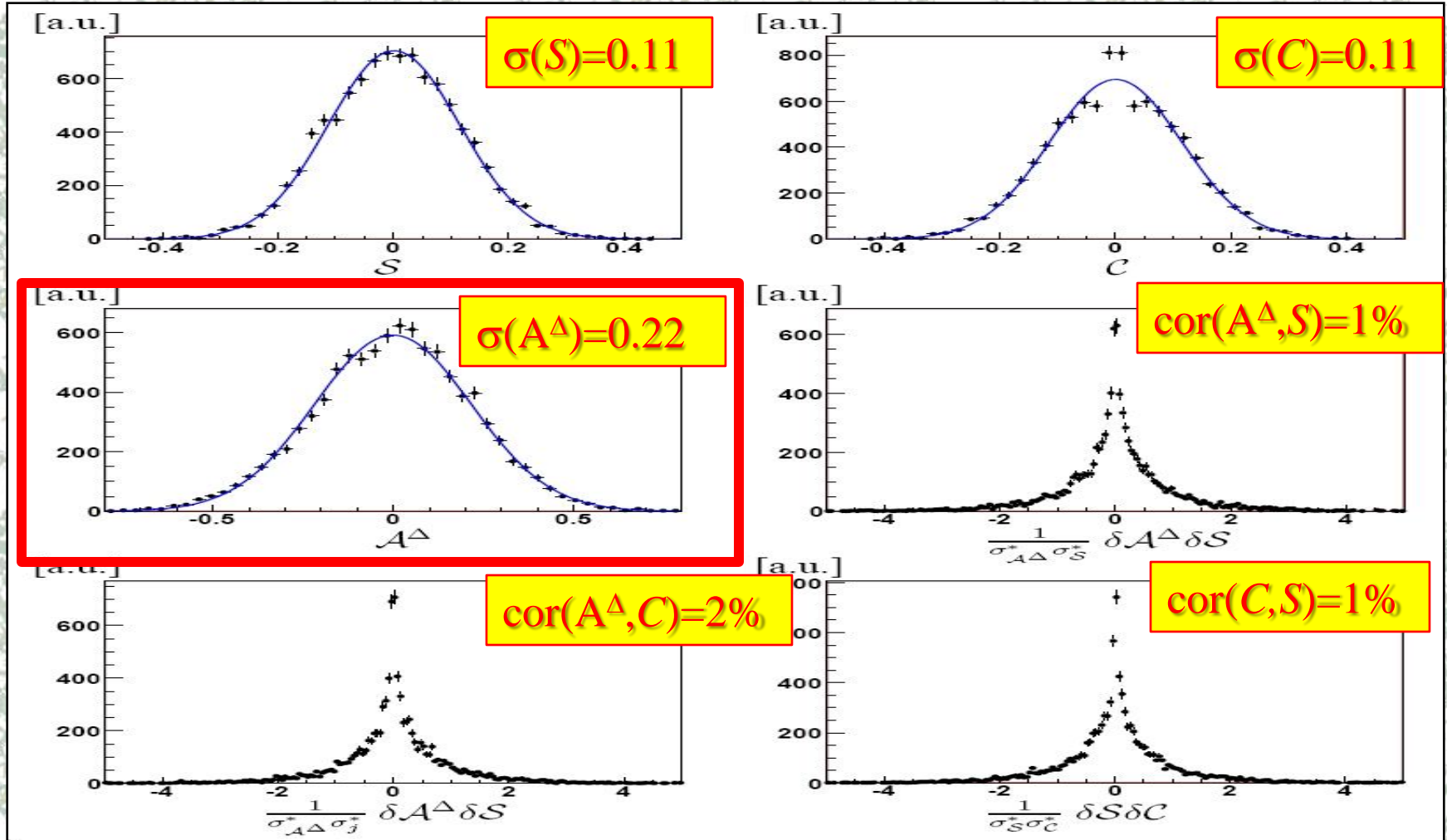
$$P_{\kappa}(t, m) = f_s \frac{\{e^{-\Gamma\tau}[I_+(\tau) + \kappa(1 - 2\omega)I_-(\tau)]\} \otimes G(t - \tau)\varepsilon(t)g_s(m)}{\int \{e^{-\Gamma\tau}[I_+(\tau) + \kappa(1 - 2\omega)I_-(\tau)]\} \otimes G(t' - \tau)\varepsilon(t')dt'} + (1 - f_s)\varepsilon_b(m, t),$$

$$I_+(\tau) = \cosh \frac{\Delta\Gamma\tau}{2} - \mathcal{A}^{\Delta} \sinh \frac{\Delta\Gamma\tau}{2}$$

$$I_-(\tau) = \mathcal{C} \cos \Delta m_s \tau - \mathcal{S} \sin \Delta m_s \tau$$

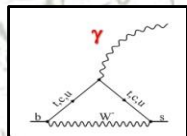
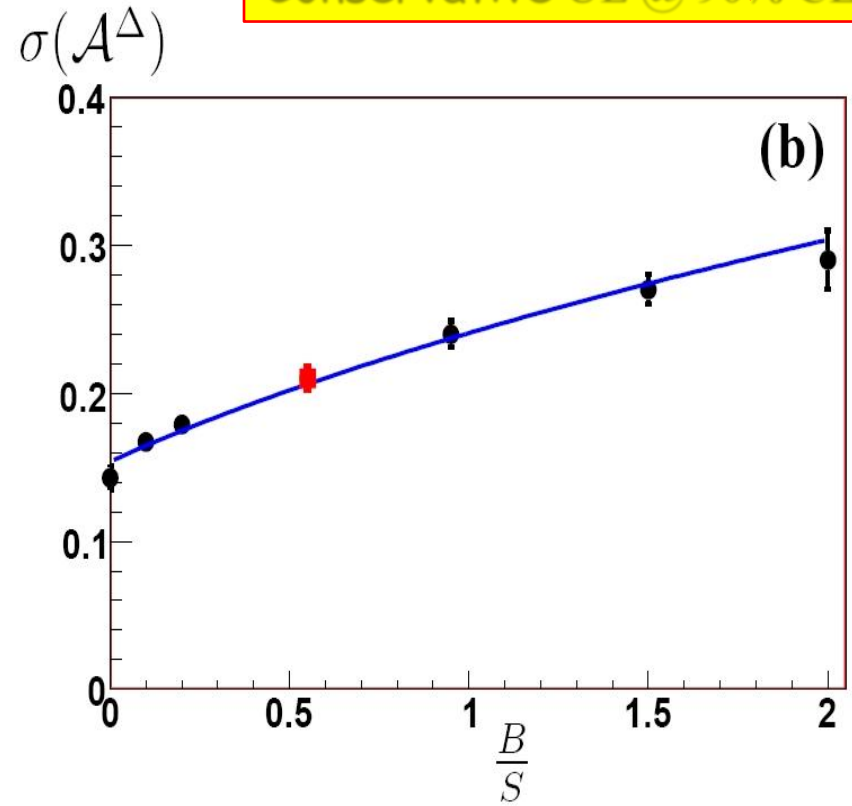
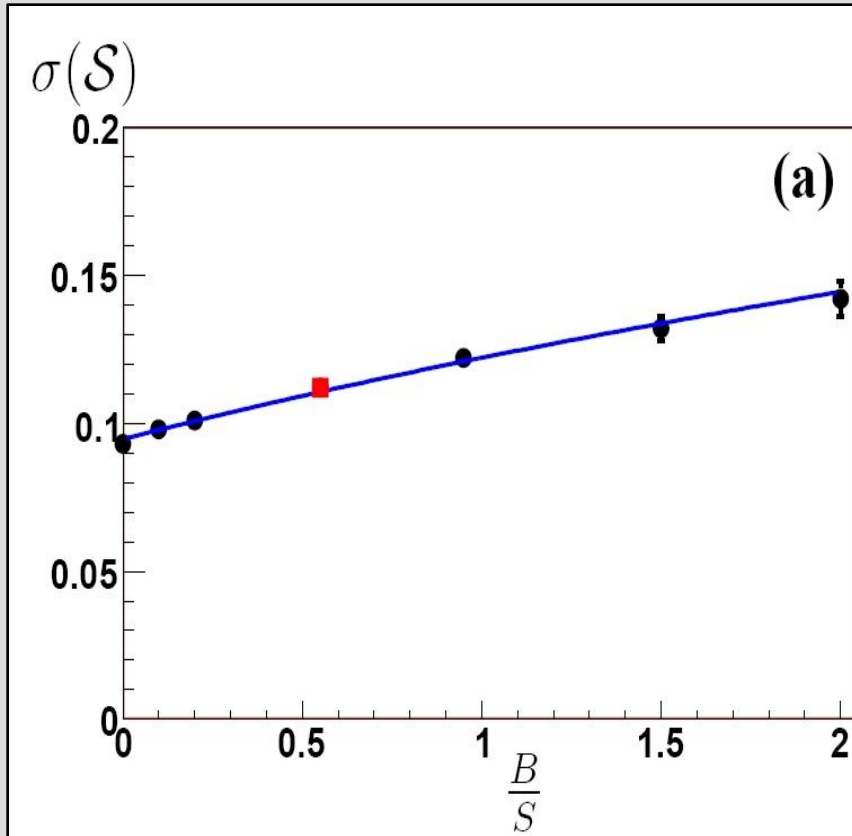
$$\mathcal{L}_0 = \prod_{i=1}^{N_{B_s}} P_{-1}(m_i, t_i, \sigma_{ti}) \prod_{i=1}^{N_{\bar{B}_s}} P_1(m_i, t_i, \sigma_{ti}) \prod_{i=1}^{N_{\text{untagged}}} P_0(m_i, t_i, \sigma_{ti}),$$





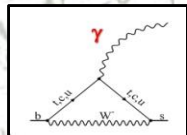
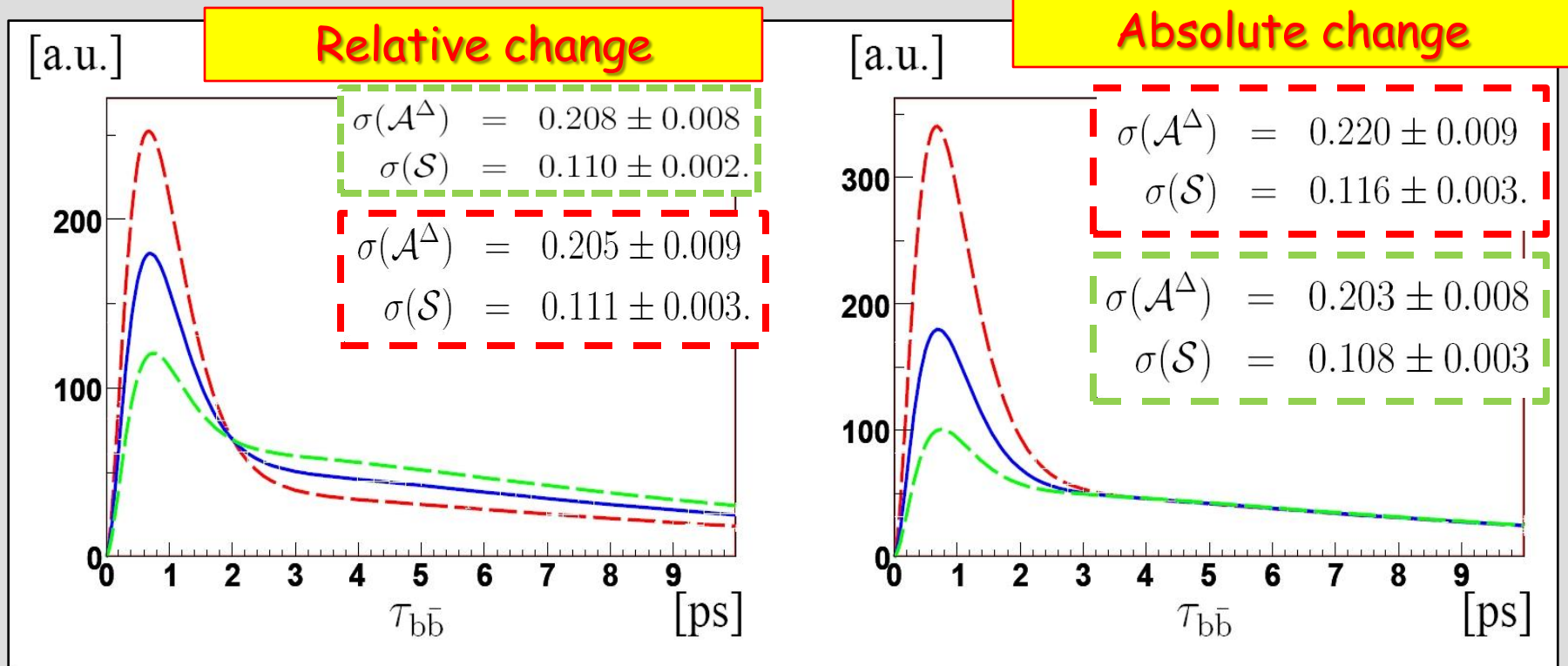
- There is some dependency on  $B/S$  level:

Conservative UL @ 90% CL

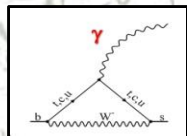
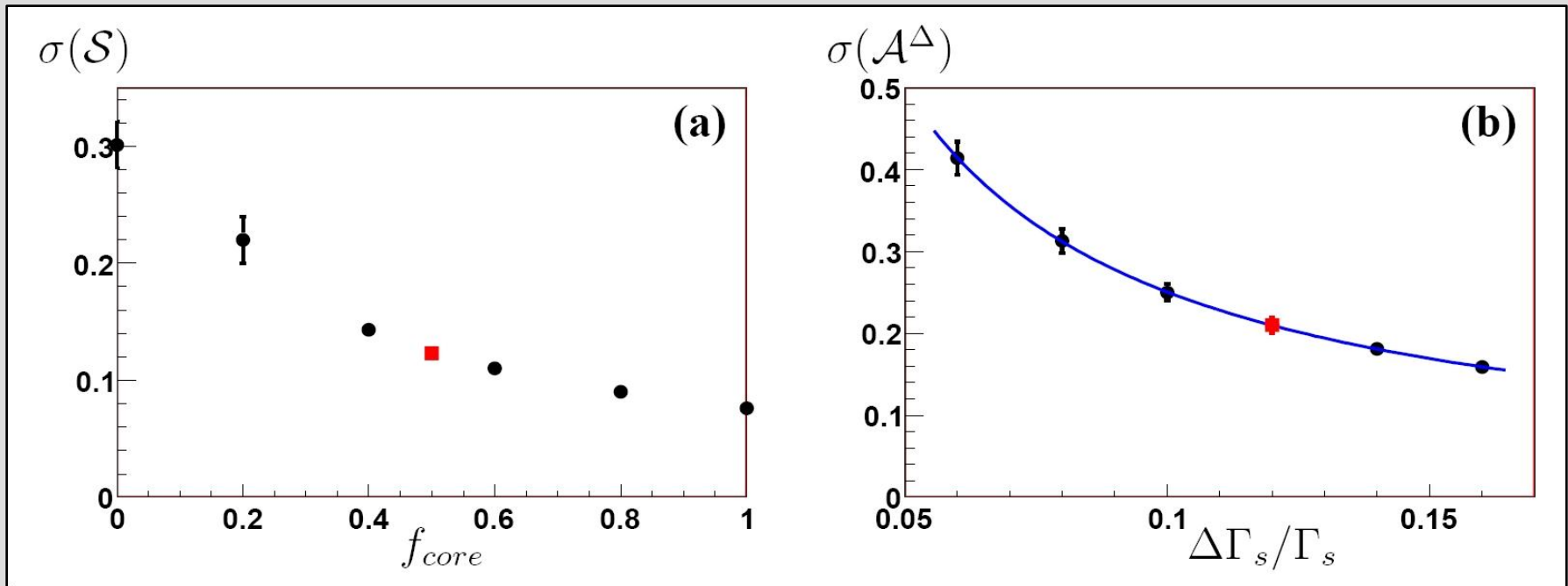




- Vary the "short/long"-lived components



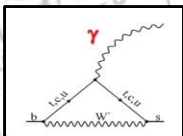
- Vary the proper time resolution
  - Use simple model with two Gaussians and vary the proportion



$$B_s \rightarrow \phi \gamma$$

- LHCb has good potential for measurement of photon polarization in  $B_s \rightarrow \phi \gamma$  decay
- For  $2 \text{ fb}^{-1}$ :  $\sigma(A^\Delta)=0.22$ ,  $\sigma(S)=\sigma(C)=0.11$ 
  - for  $500 \text{ pb}^{-1}$  ( $\int L dt$  for "the first year"):  $\sigma(A^\Delta) \sim 0.4$
- The result has *moderate* dependency on  $B/S$
- *The determination of proper time acceptance function from data is under the study now*

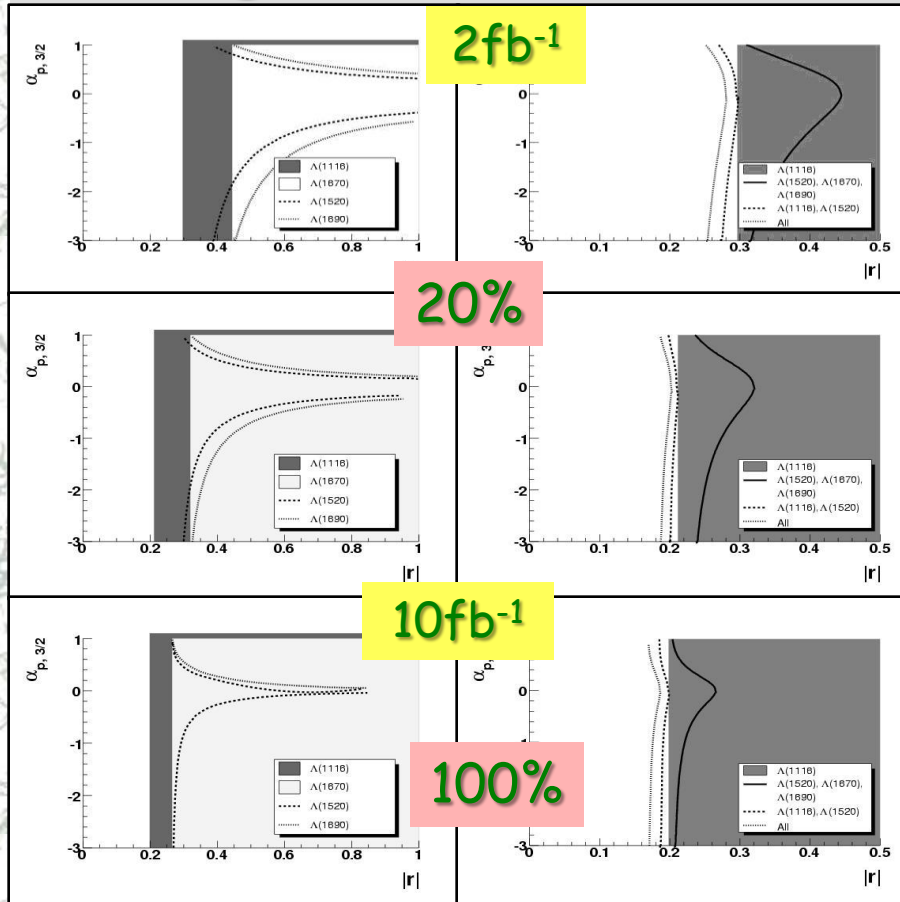
L.Shchutska, Y.Xie, A.Golutvin, V.Egorychev, V.Shevchenko, I.Belyaev LHCb-2007-147





# 3 $\sigma$ measurement of $\tan\psi$ for $\Lambda_b \rightarrow \Lambda^0 \gamma$ $\Lambda_b \rightarrow \Lambda(X) \gamma$

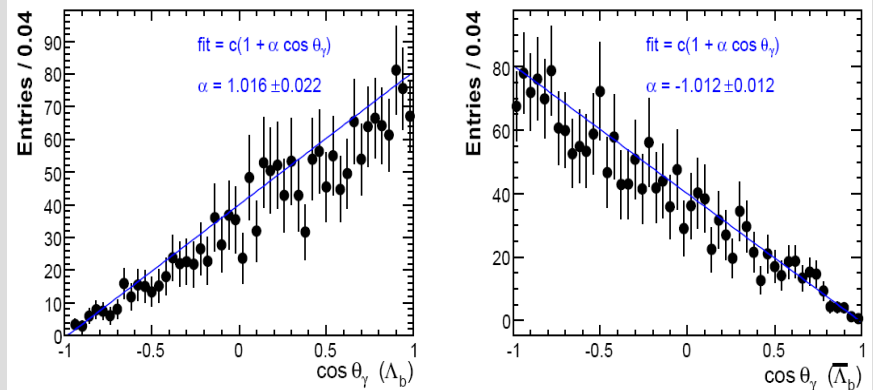
## Single Combined



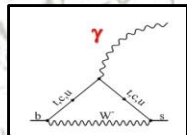
$$|A_R/A_L| = \tan\psi \equiv |r|$$

• Can be probed at 3 $\sigma$  for 5 years

- >20% with  $\Lambda_b \rightarrow \Lambda^0 \gamma$
- >25% with  $\Lambda_b \rightarrow \Lambda^* \gamma$



F.Legger & T.Schitinger, LPHE 2006-003, LHCb 2006-013



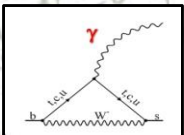
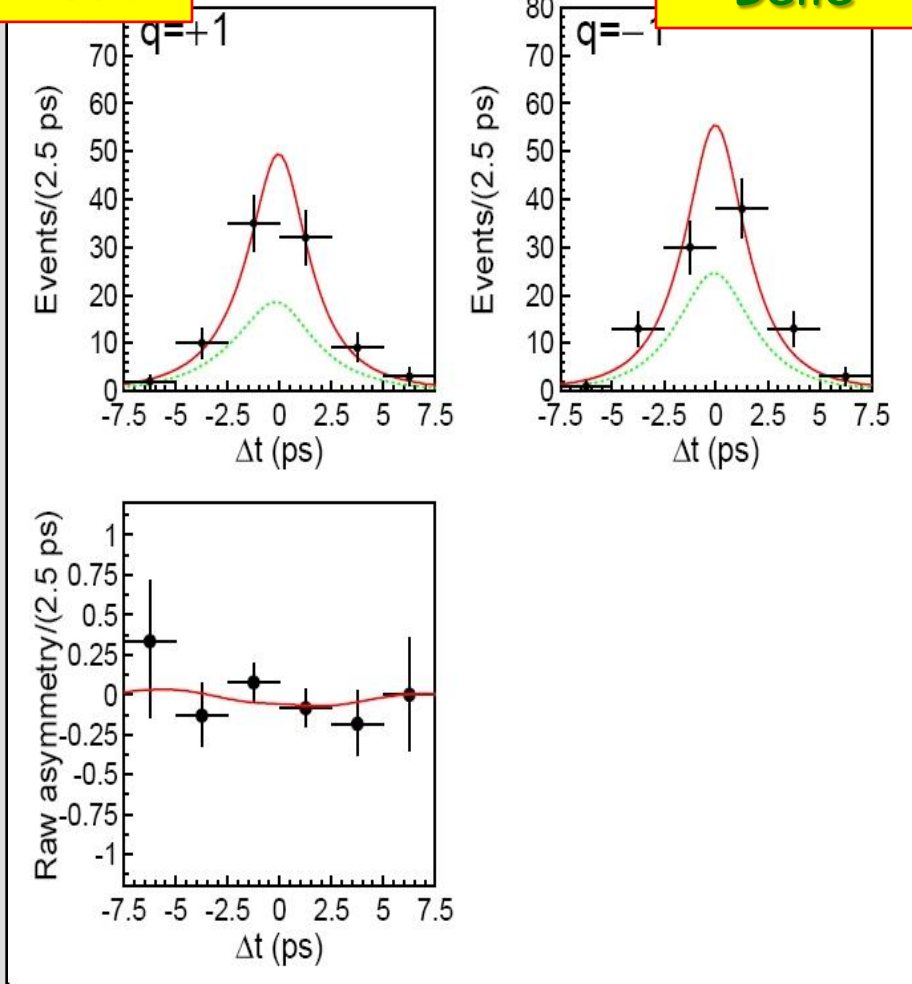
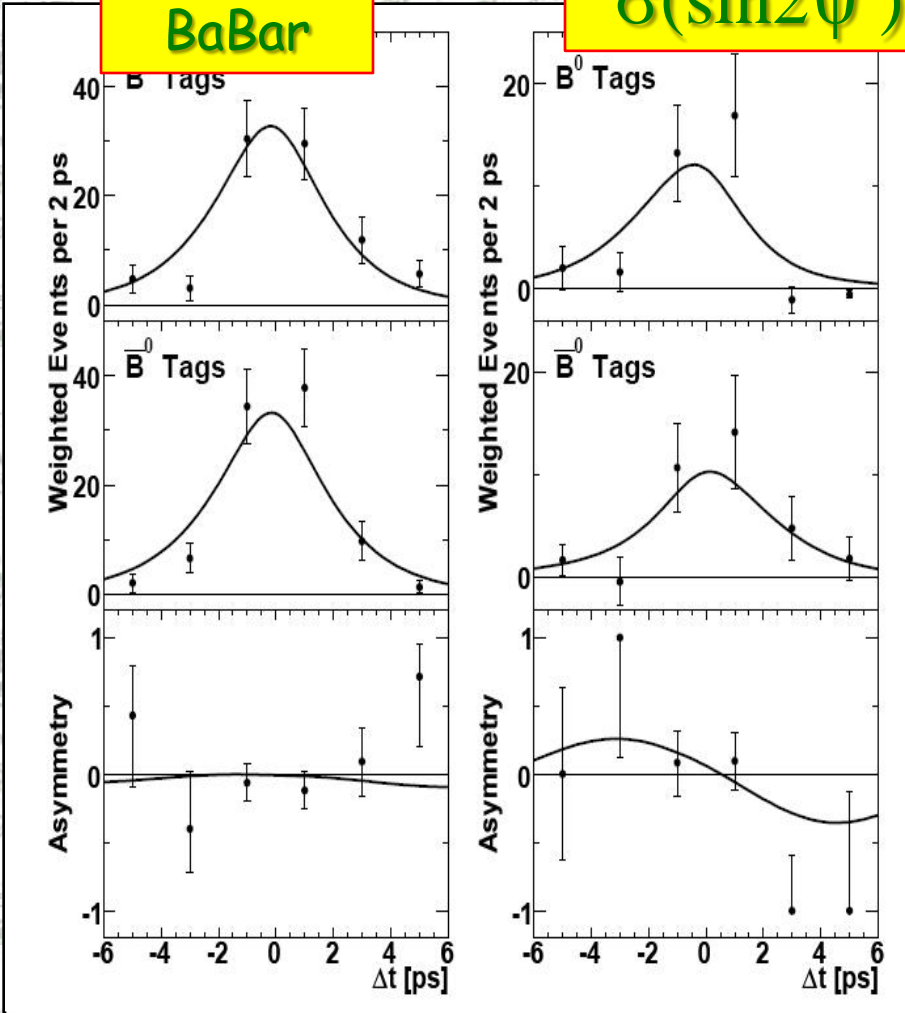


# sin2ψ now for B-factories

$\sigma(\sin 2\psi) \sim 0.4$

**BaBar**

**Belle**



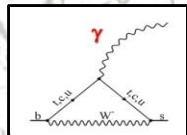
- Is expected to be small
- Statistical error of practically negligible
- Systematic error is an issue ☹
- But: use the normalization channel  $B^0 \rightarrow K^{*0} J/\psi$

$\mathcal{R} = \frac{\mathcal{N}_{B^0 \rightarrow K^{*0} \gamma}}{\mathcal{N}_{B^0 \rightarrow K^{*0} J/\psi}}$	$\bar{\mathcal{R}} = \frac{\mathcal{N}_{\bar{B}^0 \rightarrow \bar{K}^{*0} \gamma}}{\mathcal{N}_{\bar{B}^0 \rightarrow \bar{K}^{*0} J/\psi}}$	$\mathcal{A}_{\mathcal{R}} = \frac{\bar{\mathcal{R}} - \mathcal{R}}{\bar{\mathcal{R}} + \mathcal{R}}$	$\mathcal{A}_{\mathcal{R}} = \frac{\mathcal{A}_{K^{*0} \gamma}^{\text{dir}} - \mathcal{A}_{K^{*0} J/\psi}^{\text{dir}}}{2 - \frac{1}{2} \mathcal{A}_{K^{*0} \gamma}^{\text{dir}} \mathcal{A}_{K^{*0} J/\psi}^{\text{dir}}}$
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$$\mathcal{A}_{K^{*0} \gamma}^{\text{dir}} \equiv \frac{\mathcal{B}_{\bar{B}^0 \rightarrow \bar{K}^{*0} \gamma} - \mathcal{B}_{B^0 \rightarrow K^{*0} \gamma}}{\mathcal{B}_{\bar{B}^0 \rightarrow \bar{K}^{*0} \gamma} + \mathcal{B}_{B^0 \rightarrow K^{*0} \gamma}} = \frac{2\mathcal{A}_{\mathcal{R}} + \mathcal{A}_{K^{*0} J/\psi}^{\text{dir}}}{1 + \frac{1}{2} \mathcal{A}_{\mathcal{R}} \mathcal{A}_{K^{*0} J/\psi}^{\text{dir}}}$$

$$\frac{\int \varepsilon(\Phi) (p^s(\Phi) - p^c(\Phi)) d\Phi}{\int \varepsilon(\Phi) (p^s(\Phi) + p^c(\Phi)) d\Phi}$$

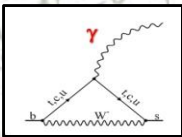
O(1%) is reachable





- Good prediction from theory
  - A.Ali, B.D. Pecjak and C.Greub, Eur.Phys.J. C55:577-595,2008
  - + test the formfactors
- The absolute measurements impossible?
- But *the double ratio* is easy to measure and almost free from the experimental uncertainties

$$\mathcal{R}_D^{B_s^0 \rightarrow \phi\gamma} = \frac{\mathcal{B}_{B_s^0 \rightarrow \phi\gamma} / \mathcal{B}_{B_s^0 \rightarrow \phi J/\psi}}{\mathcal{B}_{B^0 \rightarrow K^{*0}\gamma} / \mathcal{B}_{B^0 \rightarrow K^{*0} J/\psi}} = \frac{\mathcal{B}_{B_s^0 \rightarrow \phi\gamma} / \mathcal{B}_{B^0 \rightarrow K^{*0}\gamma}}{\mathcal{B}_{B_s^0 \rightarrow \phi J/\psi} / \mathcal{B}_{B^0 \rightarrow K^{*0} J/\psi}},$$



- LHCb experiment is well suited for detailed study of exclusive radiative decays of beauty hadrons
  - Experimental challenge
- Study for photon polarization in  $B_s \rightarrow \phi \gamma$ ,  $\Lambda_b \rightarrow \Lambda \gamma$  and  $B^\pm \rightarrow \phi K^\pm \gamma$  are of the primary interest
  - Combined sensitivity?
- Many other interesting measurements, like direct  $CP$ -asymmetries are possible

## Stay tuned!

