

# Heavy Flavour in a Nutshell

(for a 27-km annular nut at 1.8K)

Robert W. Lambert, CERN



# Flavour physics timeline



## EXPERIMENT

Rutherford: Proton

Chadwick: Neutron

Gell-Mann: Strange

AGS : CPV in Kaons

SLAC: Up+Down

SLAC: Charm

E288: Bottom

Argus: B-mixing

$D^0$  + CDF: Top

1815 1919 1920 1932 1953 1963 1964 1968 1970 1973 1974 1977 1987 1995

Prout: Proton

Rutherford: Neutron

Cabbibo: angle

Gell-Mann: 8-fold way

Anderson: EWSB

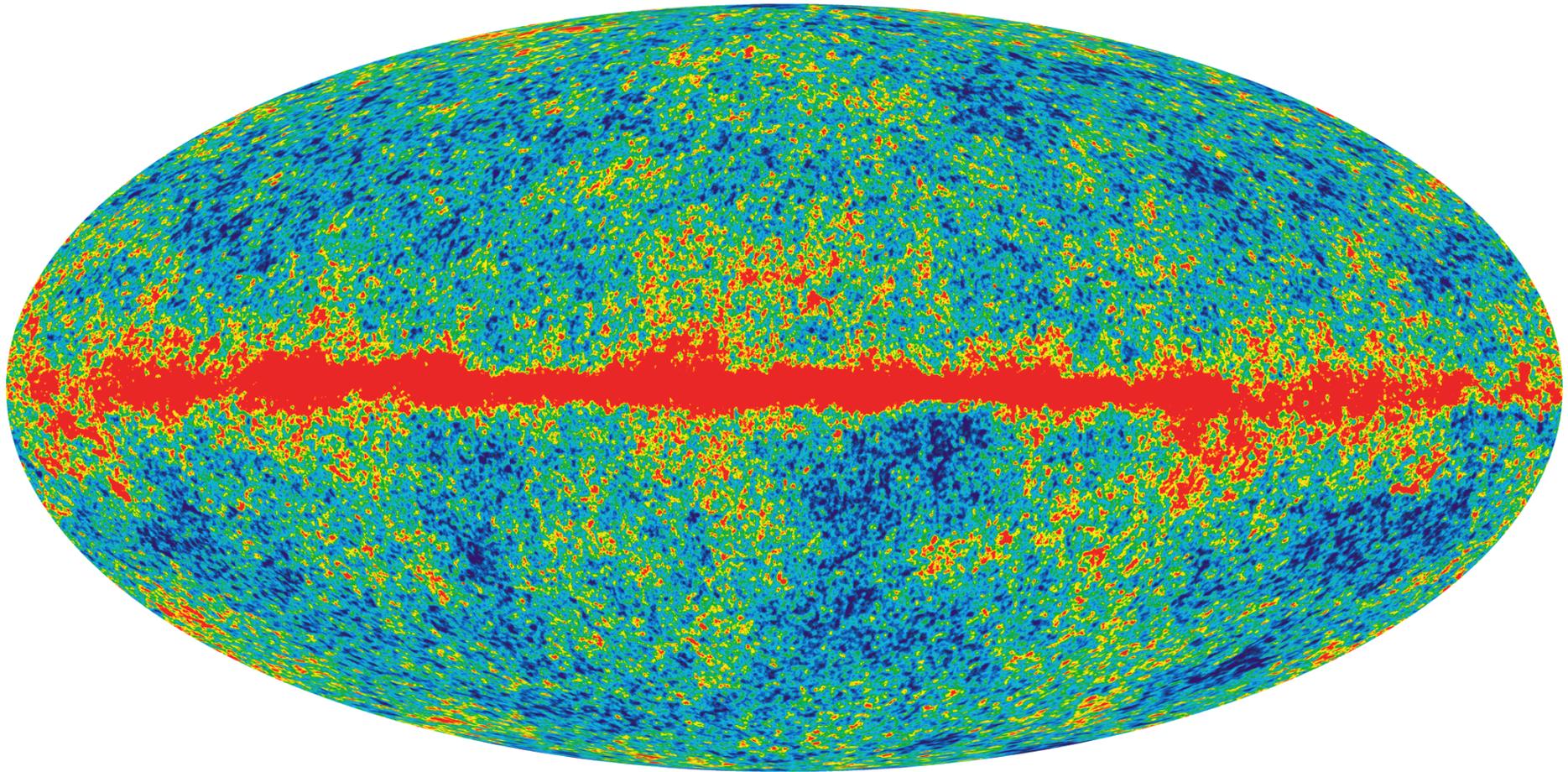
Brout-Englert-Higgs-  
Guralnik-Hagen-Kibble

GIM: charm

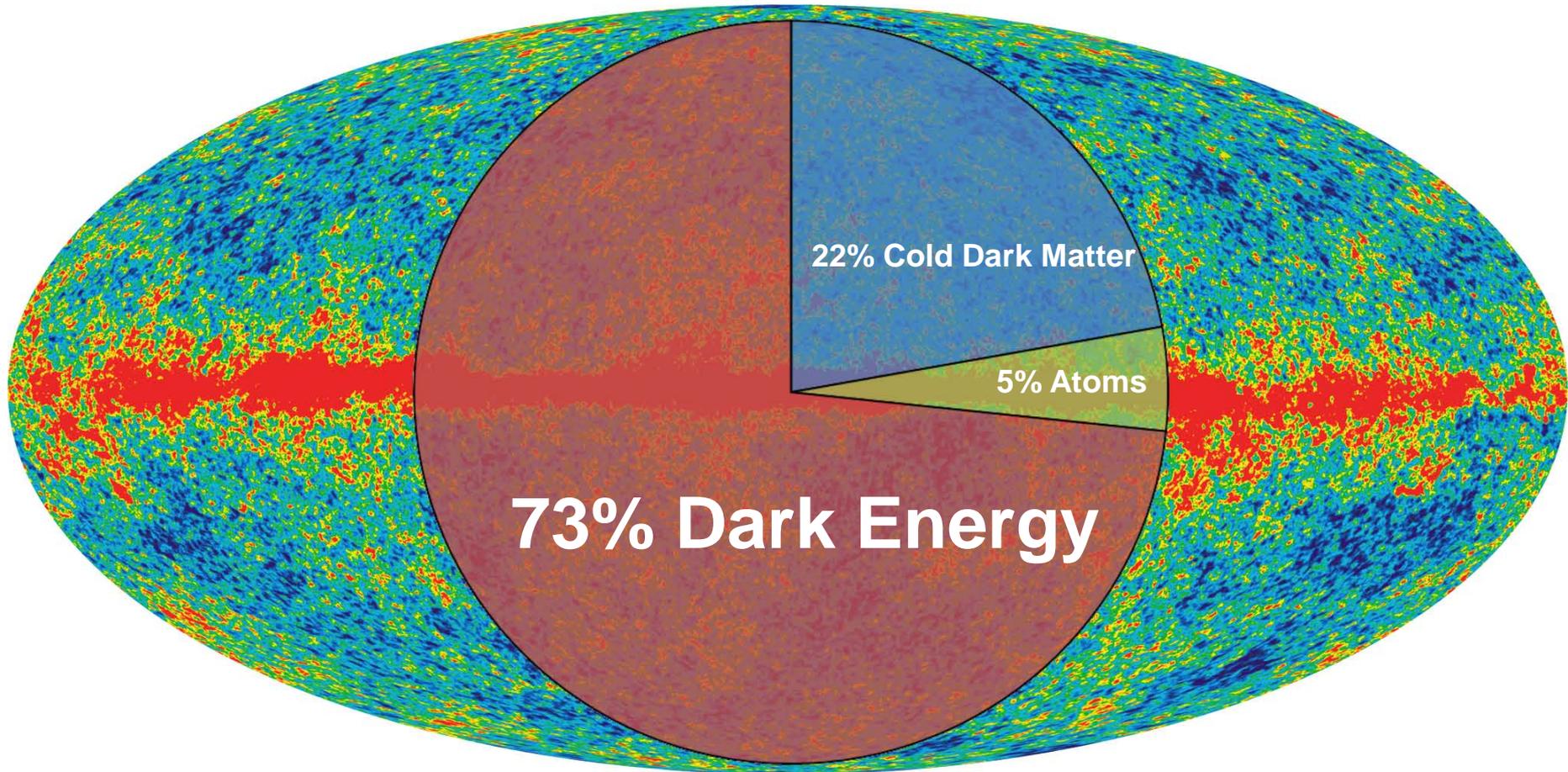
CKM: 3<sup>rd</sup> Generation

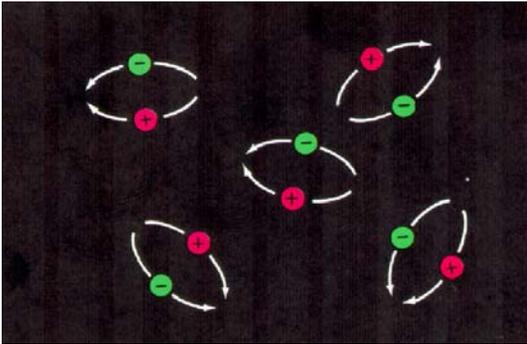
## THEORY

1. Welcome to our universe
  2. Introduction to flavour physics
  3. Hottest new physics searches
  4. Flavour-specific asymmetry
- Recent papers:
- $D\bar{D}$  measurement of  $A^b$ ,  $3.2\sigma$  deviation from the SM (May 2010)  
Evidence for an anomalous like-sign dimuon charge asymmetry  
[PRL. 105, 081801 \(2010\)](#)
  - Nierste and Lenz B-mixing update (Feb 2011)  
Numerical updates of lifetimes and mixing parameters of B mesons  
[hep-ph arxiv:1102.4274](#)
  - WMAP 7-year sky maps (Feb 2011)  
Seven-Year Wilkinson Microwave Anisotropy Probe (WMAP) Observations: Sky Maps, Systematic Errors, and Basic Results  
[Jarosik, N., et.al., 2011, ApJS, 192, 14](#)

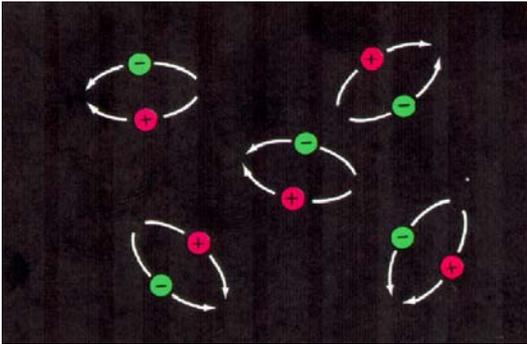


$(13.75 \pm 0.13)$  Gyr



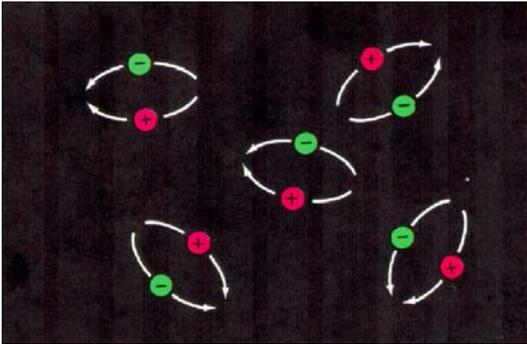


**Matter** + **Antimatter** = photons



**Matter + Antimatter = photons**

**CP-violation, CPV**  
observable difference between  
matter and antimatter



**Matter + Antimatter = photons**

**CP-violation, CPV**  
observable difference between  
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## REALITY

$$\frac{n_{baryon}}{n_\gamma} = (5.5 \pm 0.5) \times 10^{-10}$$

You Are Here

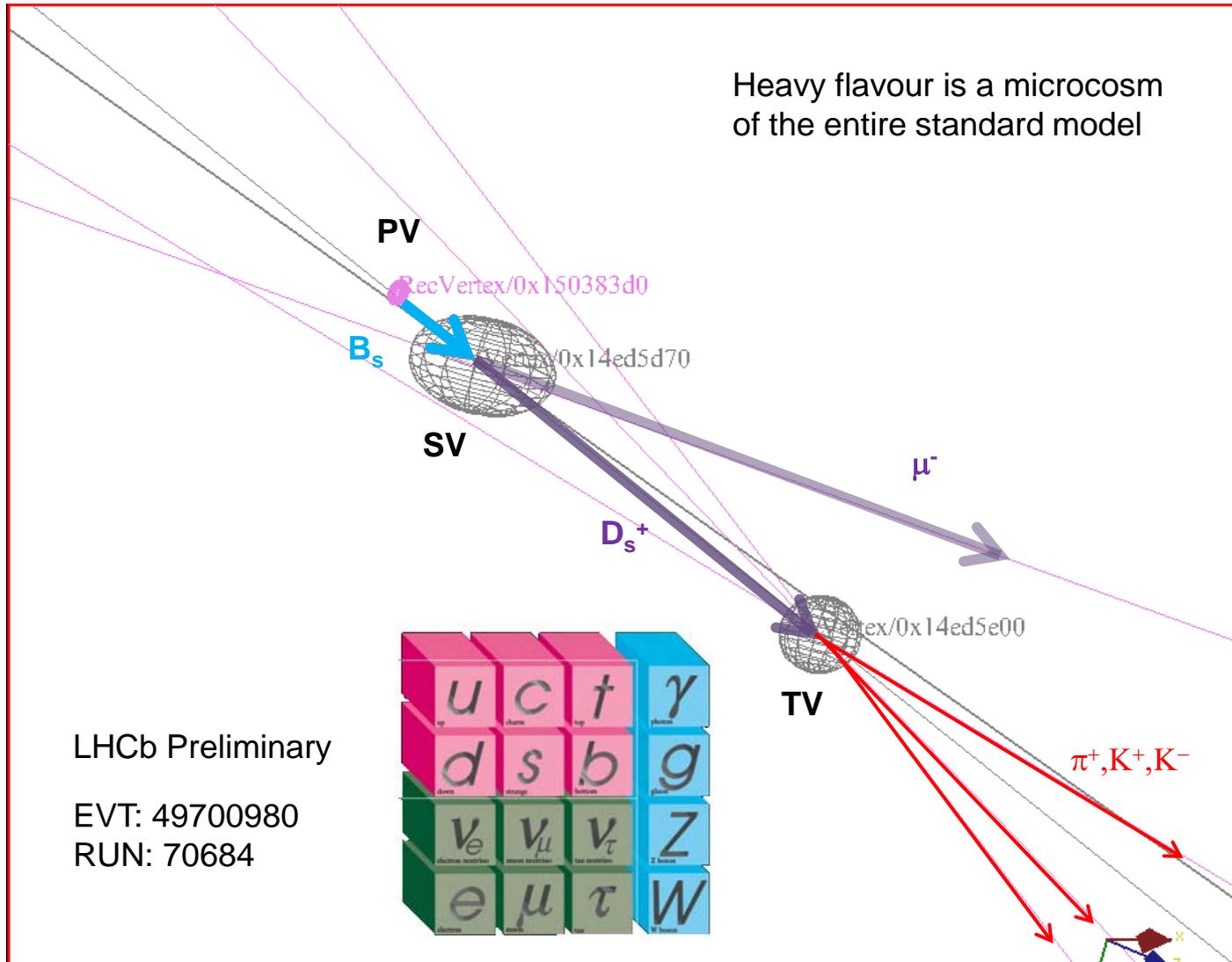
## SM (maximal CPV)

$$\frac{n_{baryon}}{n_\gamma} < 10^{-20}$$

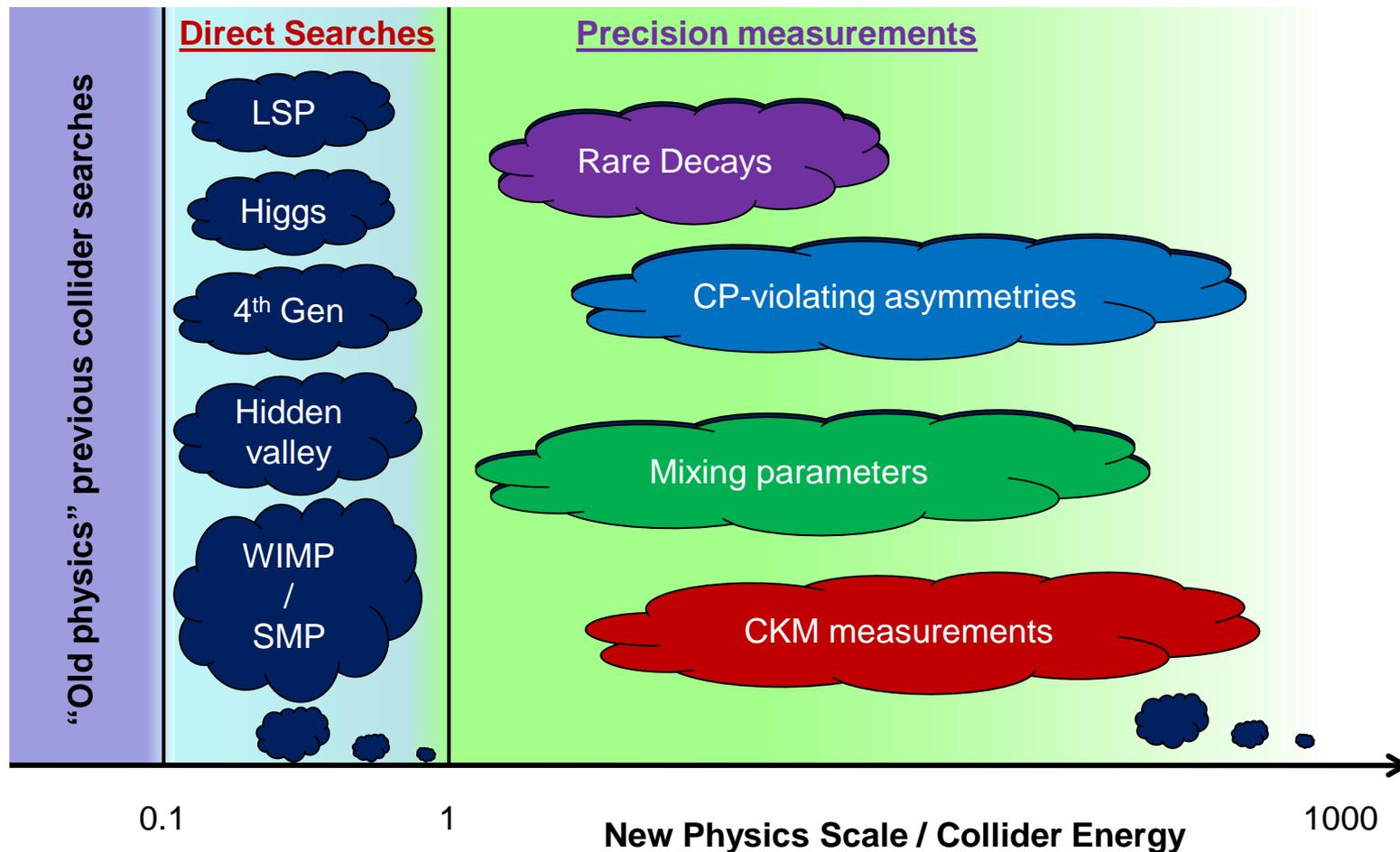
Where did you go?

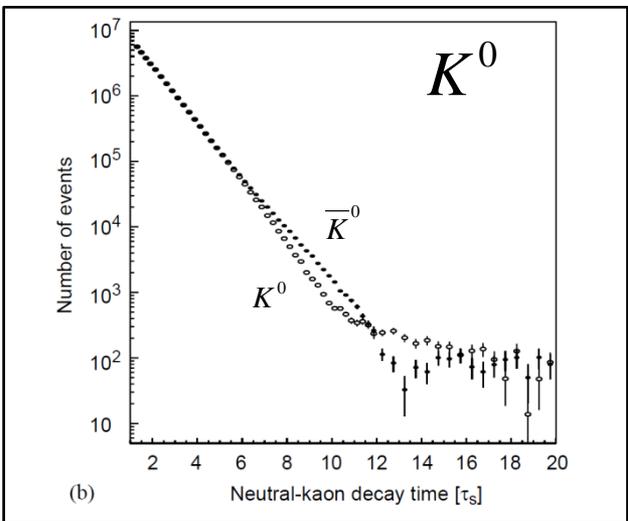
Guys...? Guys...??

# What does that have to do with heavy flavour physics?

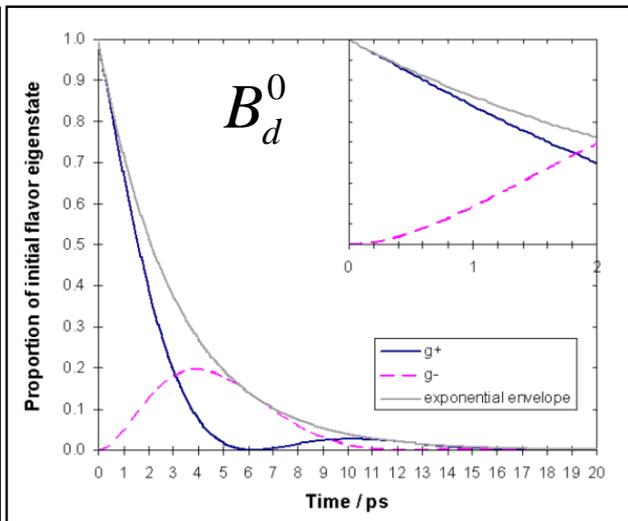


- There are in general two types of new physics searches

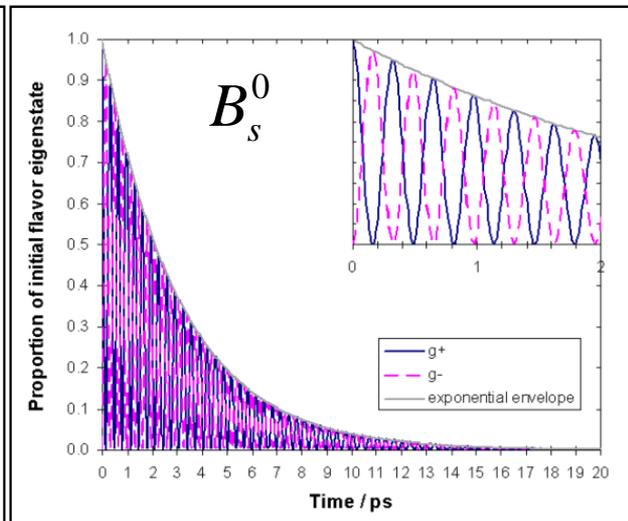




(CPLear)



(PDG)



(PDG)

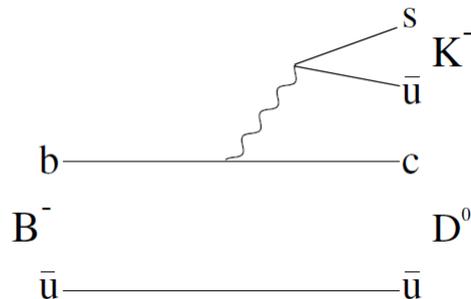
- “*mass-decay eigenstates are not the flavour eigenstates*”
  - Probably the weirdest phenomenon in physics!
  
- “*neither of those are the CP-eigenstates*”
  - CP-violation is very weird in itself
  - Observation of CPV in Kaons in 1964, before any predictions!

1. Where is the CP-violation we need?
  2. What is the flavour structure of new-physics?
- But first we ask ourselves:
- How can we best look for this new physics, and where?

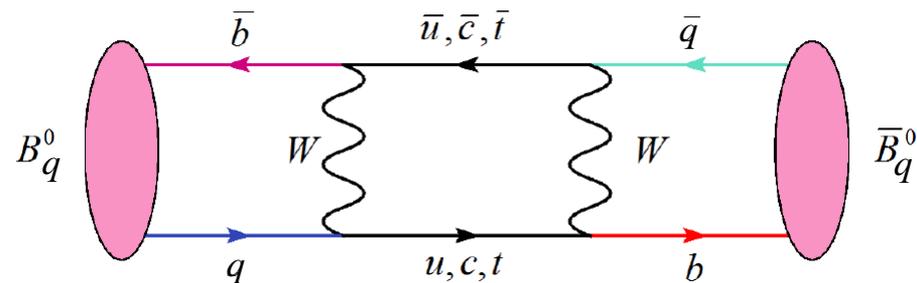


1. Find a place where new physics is unlikely
2. Precisely measure well-predicted observables
3. Find a place where new physics could enter
4. Precisely measure related observables

Unlikely: tree-level decays

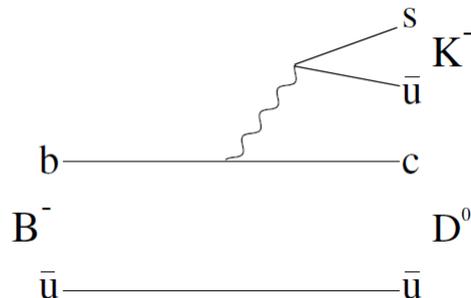


Likely: loops and penguins

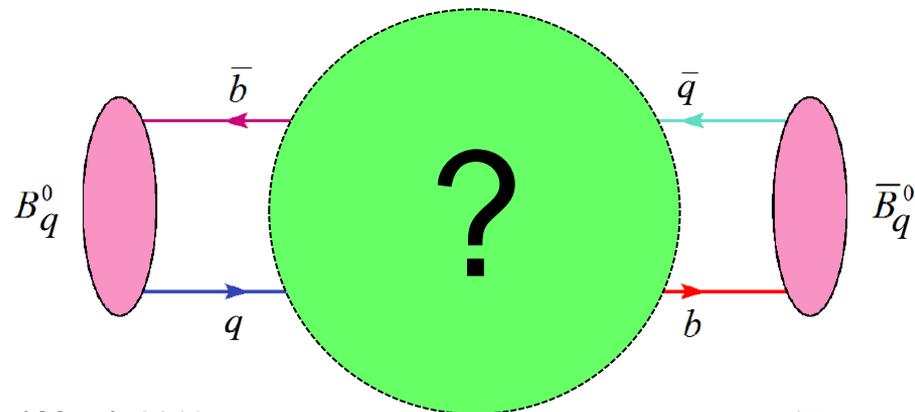


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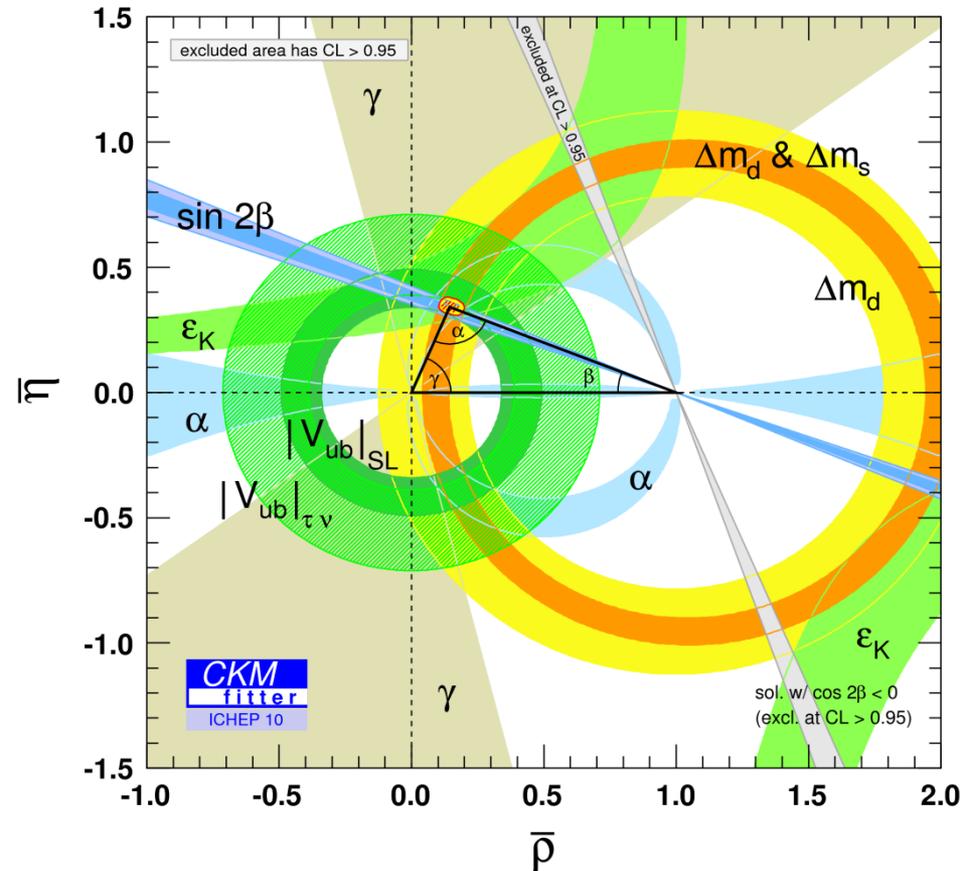


- SM has only one source of CPV, from the CKM, a phase

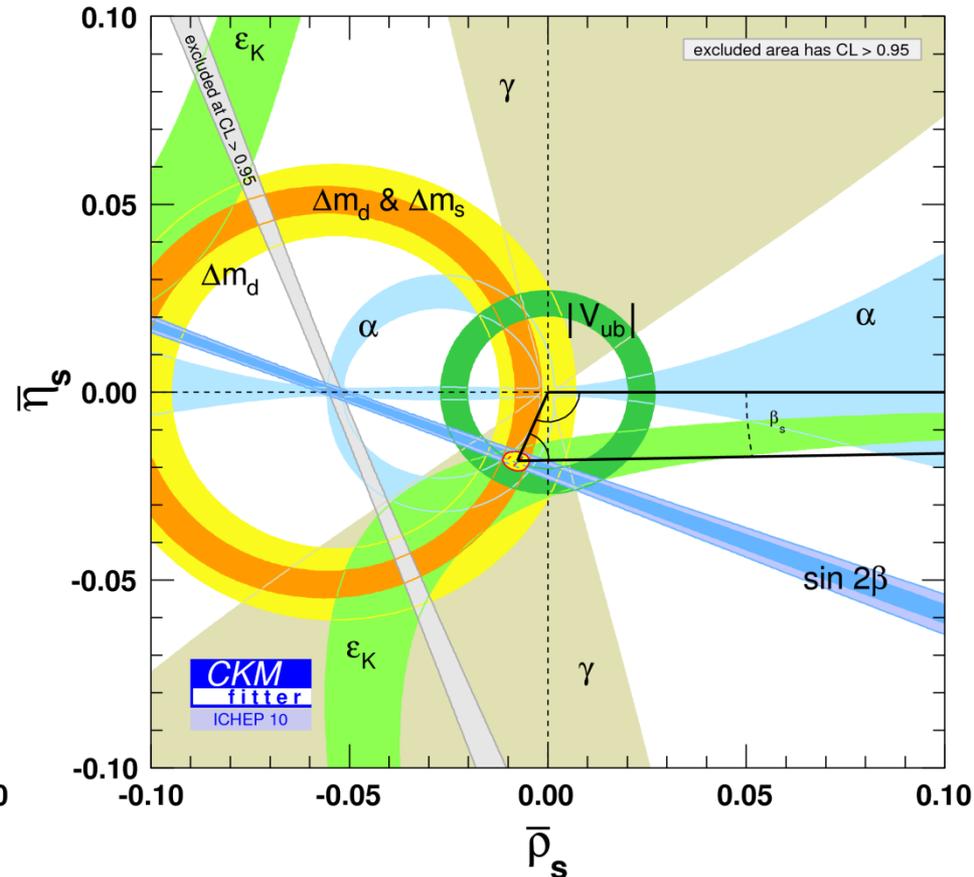
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- Observe this and any NP phase with interference:
  - Need observables with two competing amplitudes
- SM phase manifests most obviously in the  $b$ -quark system
- Measure in many different ways to constrain the same phase

- Plot everything together on a single graph
- Everything is consistent ... so far ...



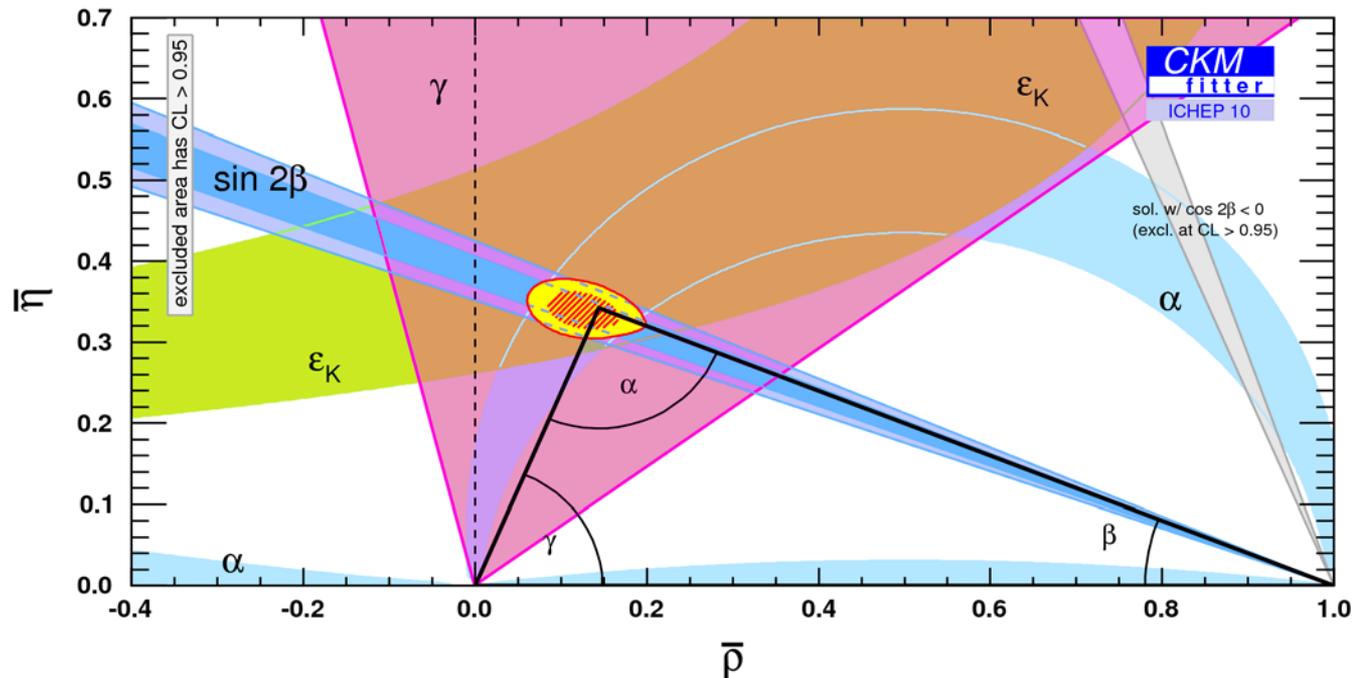
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Moriond QCD, 22nd March 2011

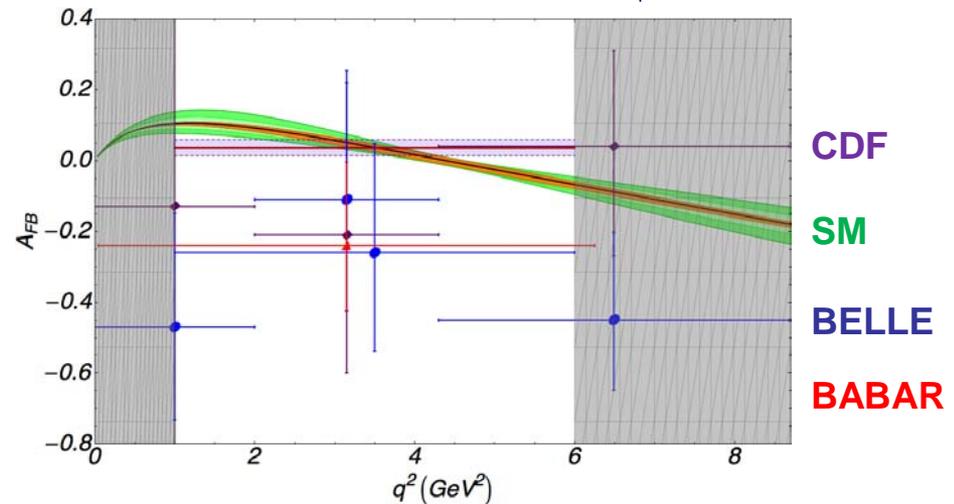
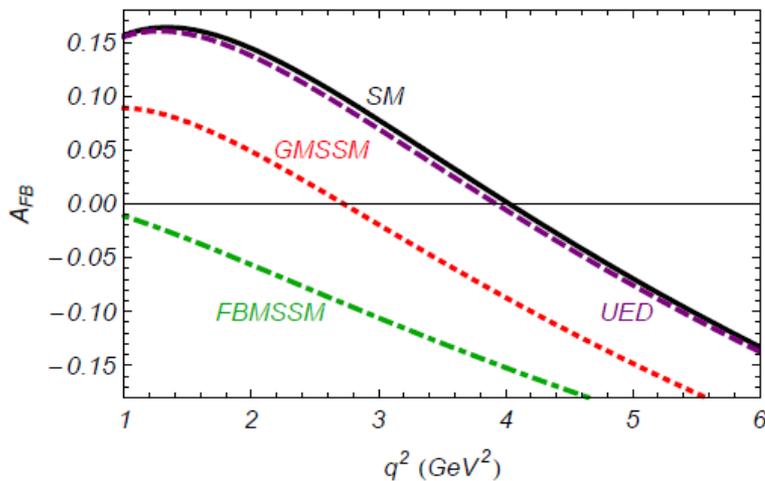
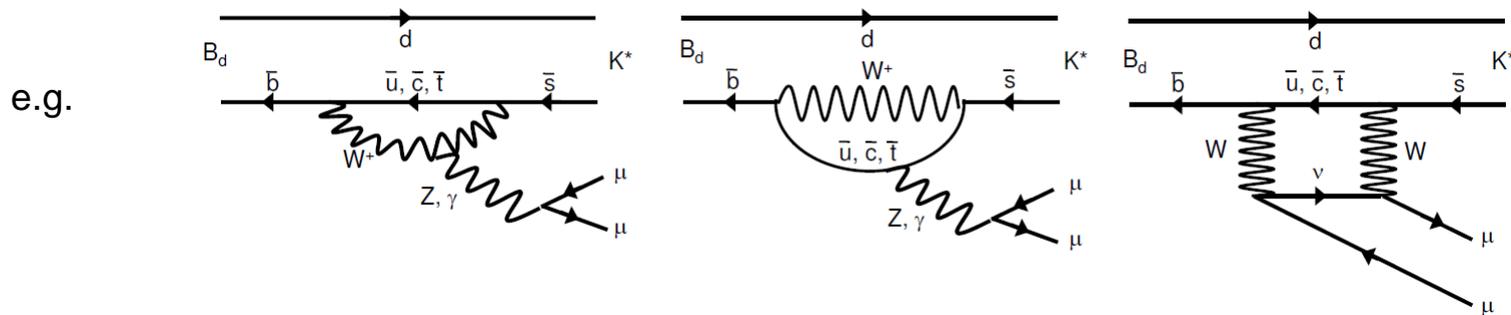
# *Hottest new physics searches*

- Check CP-violating observables
- Disagreement would point to CPV new physics

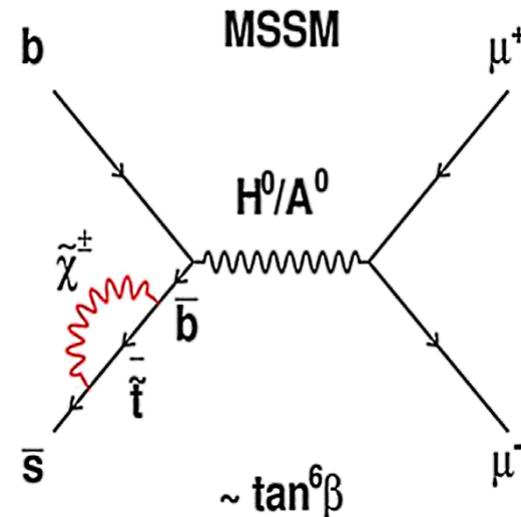
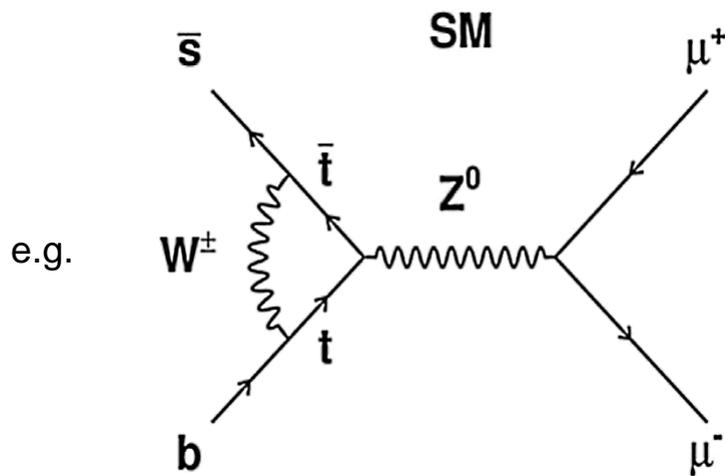


- No hints yet, but the angle  $\gamma$  is not well known

- $B_d \rightarrow K^* \mu \mu$  has both loops and penguins!
- Amongst many observables  $A_{fb}$  is sensitive to SUSY



- Very rare decays, where SM BR predictions are very good
- In the case of  $B_{s/d} \rightarrow \mu\mu$ , the rate is very sensitive to SUSY



$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)_{SM} = (3.2 \pm 0.2) \times 10^{-9}$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)_{SM} = (1.0 \pm 0.1) \times 10^{-10}$$

CDF Prelim:

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) < 43 \times 10^{-9}$$

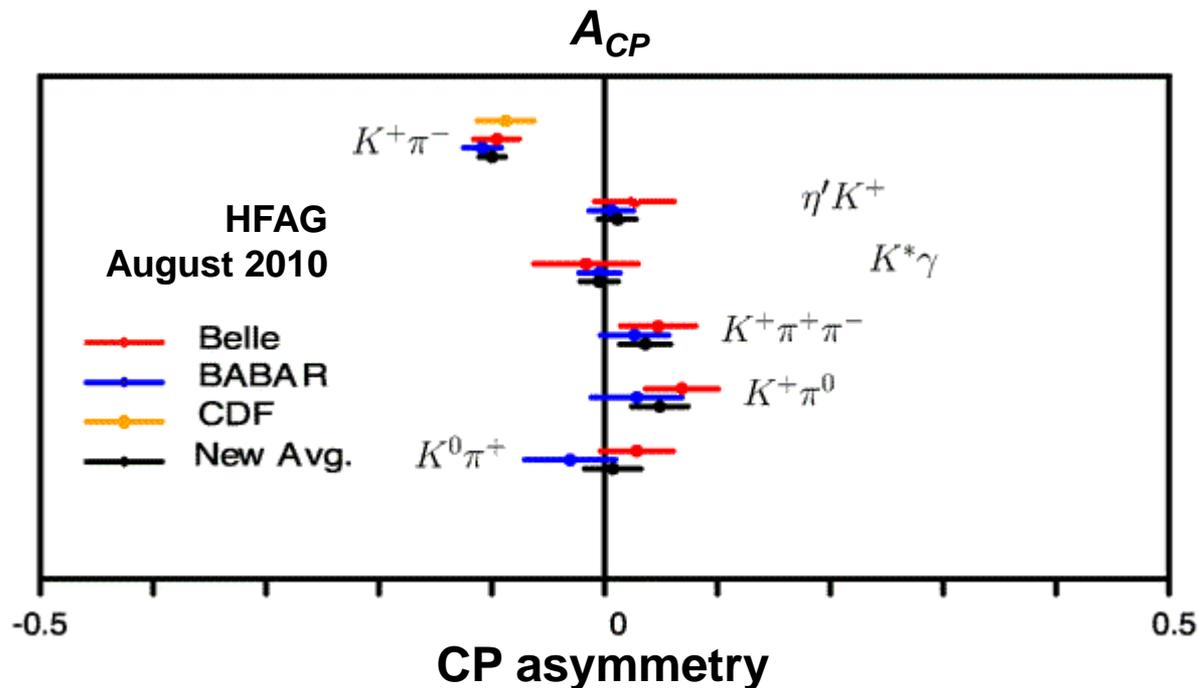
$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 7.6 \times 10^{-9}$$

LHCb:

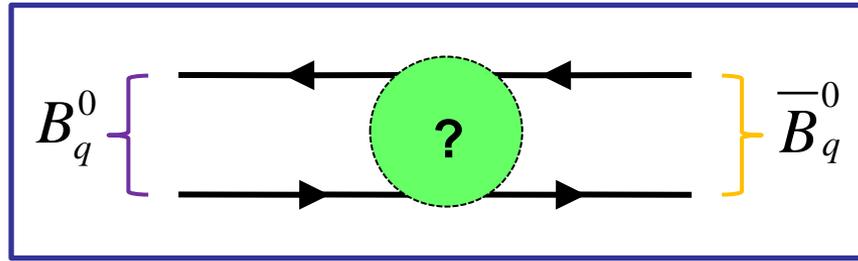
$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) < 5.6 \times 10^{-8}$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 1.5 \times 10^{-8}$$

- CP-asymmetry in decays (Direct CP-violation)
- Interesting hint: the  $B \rightarrow K\pi$  “puzzle”

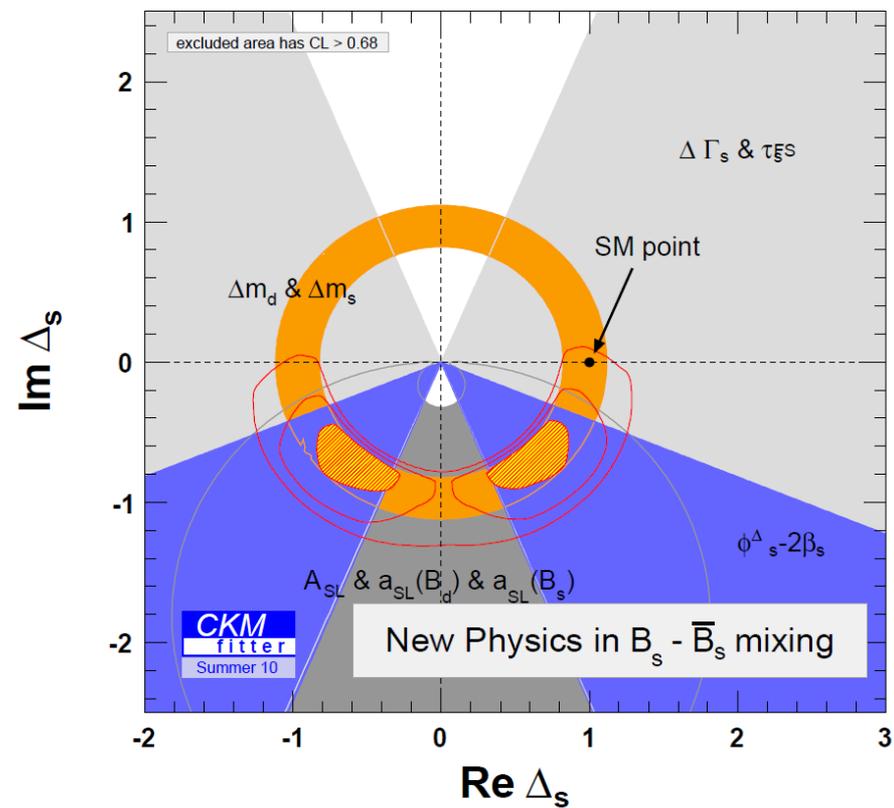
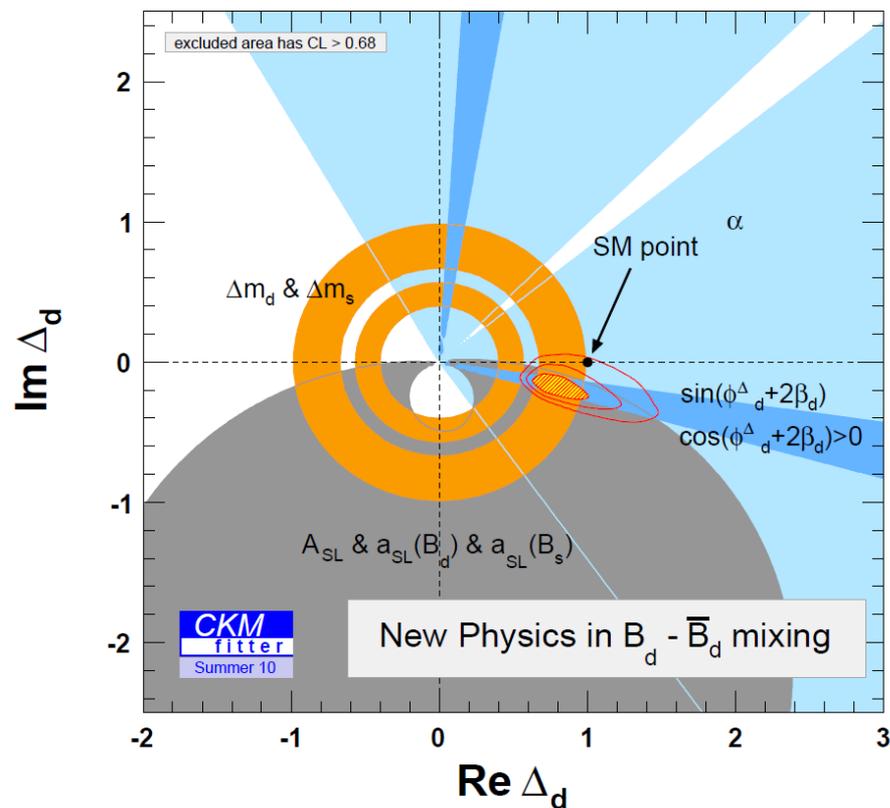


- Precision two-body  $B$ -decays will be very interesting



- Mixing can be modified in both magnitude *and* phase
- Define a complex number parameter  $\Delta_q$  for the new physics
- Just like we did with the CKM
  - Collect all the measurements together
  - Plot all at once in 2D (complex plane)

- SM is **disfavoured** by  $3.6\sigma$



- Owing a lot to the recent  $D\bar{D}$  measurement

# Flavour-specific asymmetry

... a smoking gun for new physics??

## Evidence for an anomalous like-sign dimuon charge asymmetry

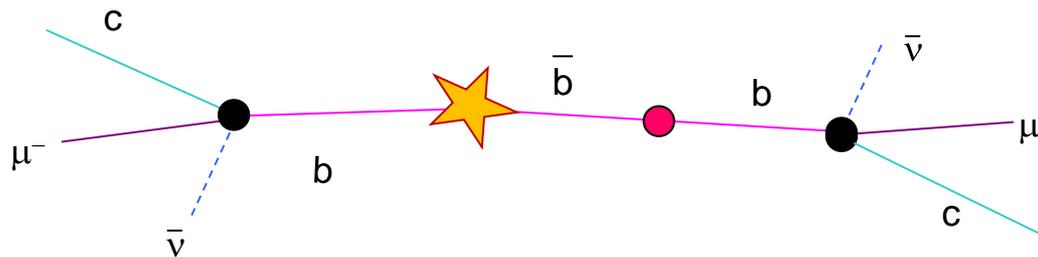
V.M. Abazov,<sup>36</sup> B. Abbott,<sup>74</sup> M. Abolins,<sup>63</sup> B.S. Acharya,<sup>29</sup> M. Adams,<sup>49</sup> T. Adams,<sup>47</sup> E. Aguilo,<sup>6</sup> G.D. Alexeev,<sup>36</sup>  
G. Alkhazov,<sup>40</sup> A. Alton<sup>a</sup>,<sup>62</sup> G. Alverson,<sup>61</sup> G.A. Alves,<sup>2</sup> L.S. Ancu,<sup>35</sup> M. Aoki,<sup>48</sup> Y. Arnoud,<sup>14</sup> M. Arov,<sup>58</sup>  
A. Askew,<sup>47</sup> B. Åsman,<sup>41</sup> O. Atramentov,<sup>66</sup> C. Avila,<sup>8</sup> J. BackusMayes,<sup>81</sup> F. Badaud,<sup>13</sup> L. Bagby,<sup>48</sup> B. Baldin,<sup>48</sup>  
D.V. Bandurin,<sup>47</sup> S. Banerjee,<sup>29</sup> E. Barberis,<sup>61</sup> A.-F. Barfuss,<sup>15</sup> P. Baringer,<sup>56</sup> J. Barreto,<sup>2</sup> J.F. Bartlett,<sup>48</sup>  
U. Bassler,<sup>18</sup> S. Beale,<sup>6</sup> A. Bean,<sup>56</sup> M. Begalli,<sup>3</sup> M. Begel,<sup>72</sup> C. Belanger-Champagne,<sup>41</sup> L. Bellantoni,<sup>48</sup>  
J.A. Benitez,<sup>63</sup> S.B. Beri,<sup>27</sup> G. Bernardi,<sup>17</sup> R. Bernhard,<sup>22</sup> I. Bertram,<sup>42</sup> M. Besançon,<sup>18</sup> R. Beuselinck,<sup>43</sup>

...

We measure the charge asymmetry  $A$  of like-sign dimuon events in  $6.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions recorded with the D0 detector at a center-of-mass energy  $\sqrt{s} = 1.96 \text{ TeV}$  at the Fermilab Tevatron collider. From  $A$ , we extract the like-sign dimuon charge asymmetry in semileptonic  $b$ -hadron decays:  $A_{\text{sl}}^b = -0.00957 \pm 0.00251 \text{ (stat)} \pm 0.00146 \text{ (syst)}$ . This result differs by 3.2 standard deviations from the standard model prediction  $A_{\text{sl}}^b(SM) = (-2.3_{-0.6}^{+0.5}) \times 10^{-4}$  and provides first evidence of anomalous CP-violation in the mixing of neutral  $B$  mesons.

PACS numbers: 13.25.Hw; 14.40.Nd

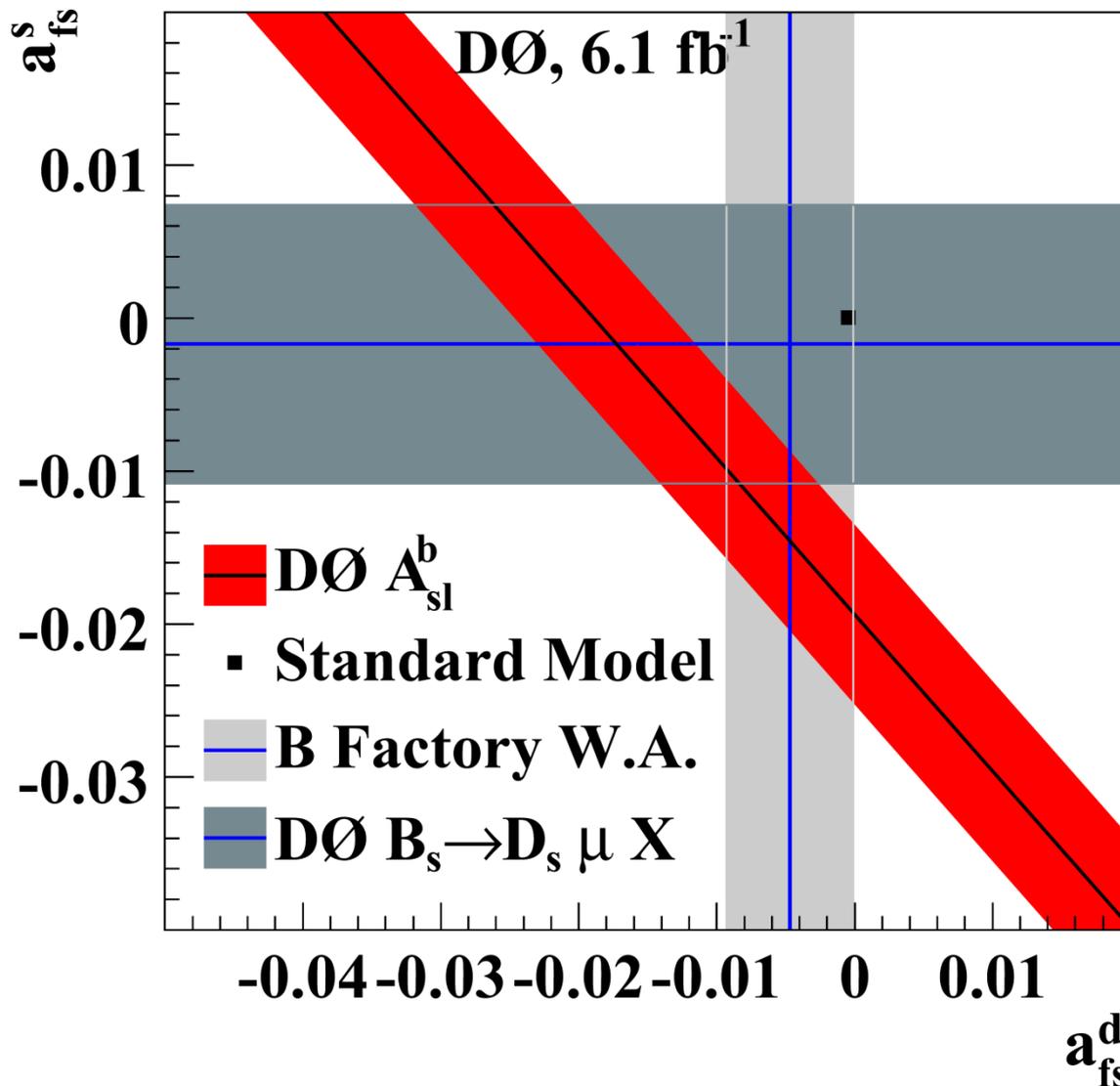
- Very difficult measurement
- Observe  $N(\mu^+ \mu^+) \neq N(\mu^- \mu^-)$
- Flavour-specific asymmetry from  $B^0$ -mixing in the SM:



CP asymmetry  
in mixing,  $a_{fs}$

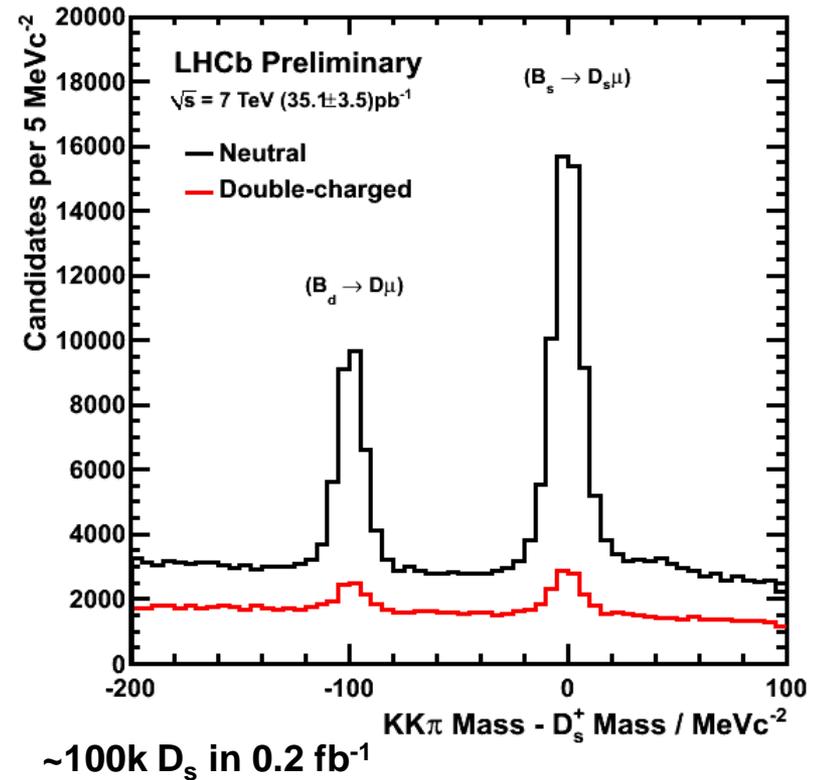
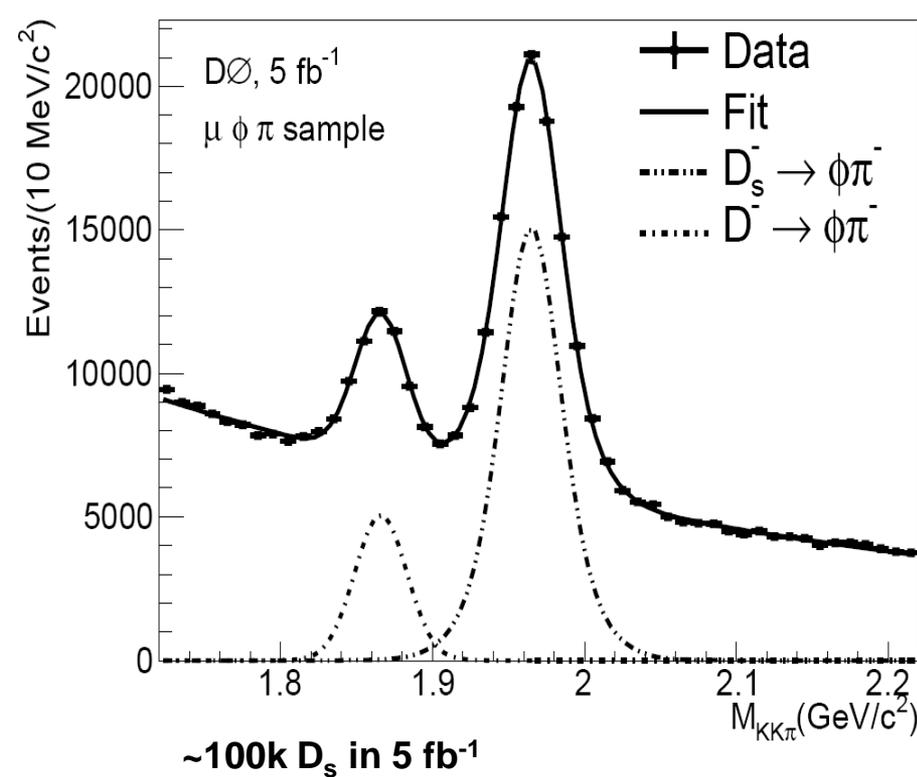
- In the standard model  $a_{fs}$  is almost negligible

$$A^b \approx \frac{a_{fs}^s + a_{fs}^d}{2} \quad SM = (-2.0 \pm 0.3) \times 10^{-4} \quad D\emptyset \approx (-1 \pm 0.3)\%$$



Situation could really be cleared up by LHCb

- LHCb is reconstructing both  $B_s^0 \rightarrow D_s^\mp \mu^\pm \nu_\mu$  and  $B_d^0 \rightarrow D_d^\mp \mu^\pm \nu_\mu$
- LHCb is catching up with DØ very quickly



- LHC is a pp-collider, not a p $\bar{p}$ -collider
- LHCb is in the forward region
  - Can't measure the same thing as DØ
  - Need a clever new method

NB: DØ  
(inclusive)

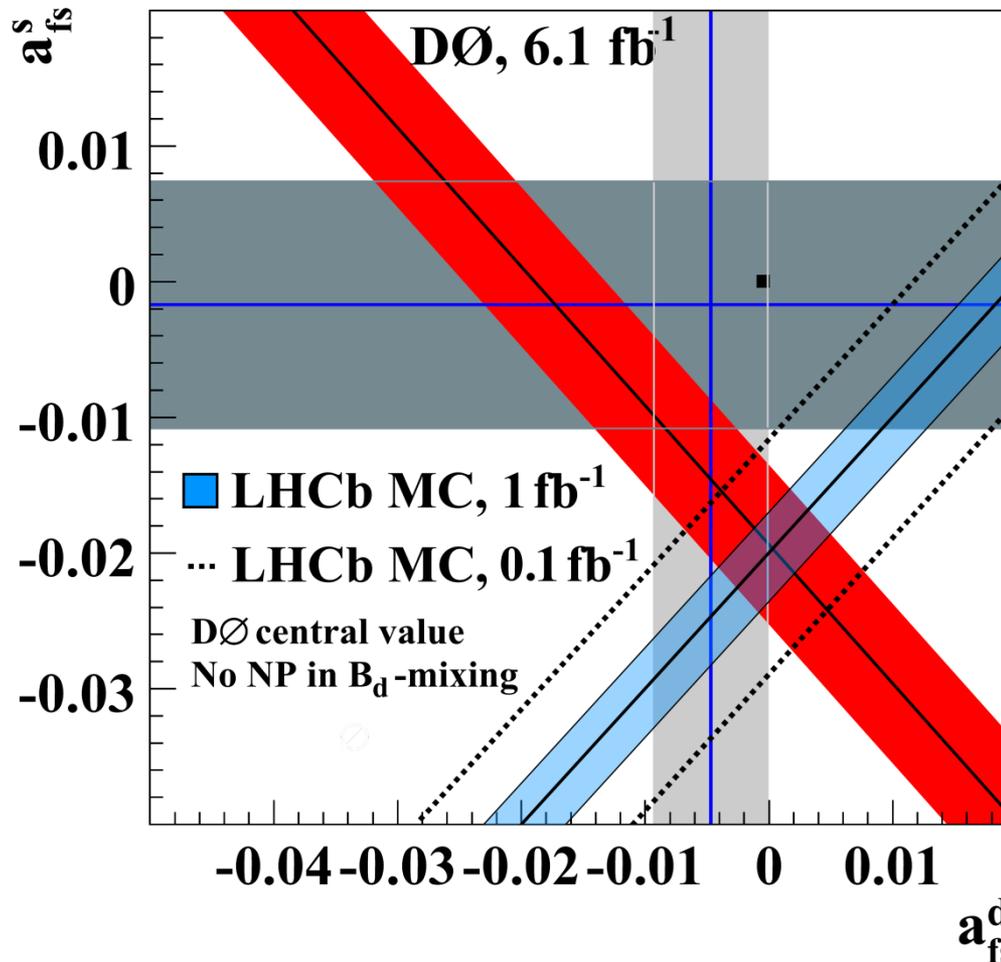
$$\left( A^b \sim \frac{a_{fs}^s + a_{fs}^d}{2} \sim -(2.0 \pm 0.3) \times 10^{-4} \right)$$

- Subtract two asymmetries to eliminate systematics

LHCb  
(subtraction)

$$\Delta A_{fs} = \frac{a_{fs}^s - a_{fs}^d}{2} = (2.1 \pm 0.3) \times 10^{-4}$$

- LHCb measurement cuts at right-angles to  $D\emptyset$



Only one example of the great physics on the way from LHCb

- ? Need new physics to explain the observed universe
  - ✓ LHC is a discovery machine
  - ✓ Precision measurements complement direct searches
- ✓ LHCb is the flavour experiment at the LHC
  - ✓  $B \rightarrow K\pi$ , CKM-angle  $\gamma$ ,  $B_{s/d} \rightarrow \mu\mu$ ,  $B_d \rightarrow K^* \mu\mu$ ,  $B_s \rightarrow J/\psi\Phi$  ...
- ? We've seen a hint of new physics already from  $D\emptyset$ 
  - ✓ LHCb will make an early complementary measurement
- This is only the start of the LHC era, so ....
  - Stay tuned for the latest experimental results!

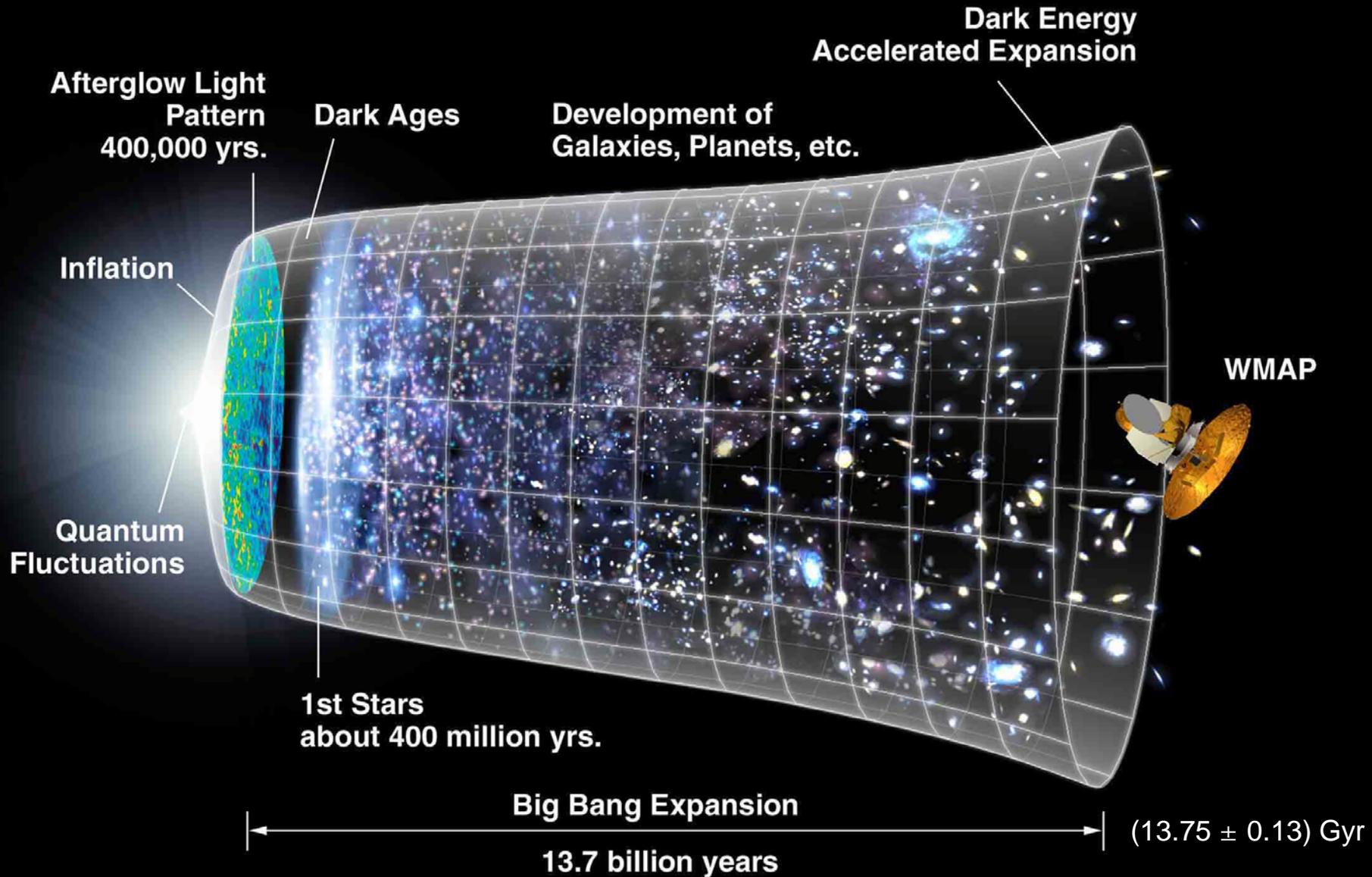
- Backups are often required

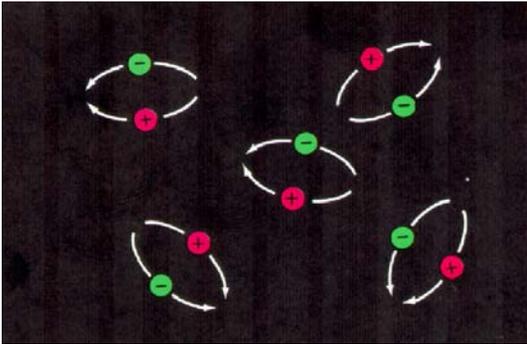


- Ulrich Kerzel for discussions on two-body  $B$ -decays
- Guy Wilkinson and Thomas Ruf for their great advice
- Johannes Albrecht for discussions on  $B_s \rightarrow \mu\mu$
- The CKM-fitter members of LHCb for updating the  $\beta_s$  plot, pointing out to me a long-standing physics goof in our TDR and other publications, and for putting up with my crazy questions about their fitting methods

- LHCb:
  - $B_s \rightarrow \mu\mu$  first result: <http://arxiv.org/abs/1103.2465>
  - Detector paper: J. of Instrumentation (No. 3 pp. S08005P)
  - “Roadmap” of physics analyses: arXiv:0912.4179
    - Chapter 2:  $\gamma$
    - Chapter 3:  $B \rightarrow K\pi$
    - Chapter 5:  $B_{s/d} \rightarrow \mu\mu$
    - Chapter 6:  $K^* \mu\mu$
  - $\Delta A_{fs}$  studies:
    - R.W. Lambert, CERN-THESIS-2009-001
    - N. Brook *et al.*, CERN-LHCb-2007-054
- CPLEar: Kaon mixing: [Physics Reports, Volume 374, Issue 3, Pages 165-270 \(January 2003\)](#)
- Experimental averages:
  - CKM fitter group : <http://ckmfitter.in2p3.fr/>
  - HFAG ( $B \rightarrow K\pi$ ): <http://www.slac.stanford.edu/xorg/hfag/rare/ichep10/acp/index.html>
- More on  $B \rightarrow K\pi$ 
  - Theory Status: S. Mishima from CKM 2010, arXiv:1101.1501
  - New Physics : S. Baek *et al.*, arXiv:hep-ph/0412086
- CDF  $B_{s/d} \rightarrow \mu\mu$  : CDF Public Note 9892 (preliminary)

# Further introduction





**Matter + Antimatter = photons**

**CP-violation, CPV**  
observable difference between  
matter and antimatter

## REALITY

$$\frac{n_{baryon}}{n_{\gamma}} = (5.5 \pm 0.5) \times 10^{-10}$$

Mass of entire solar system:  $2 \times 10^{30}$  kg

## SM (maximal CPV)

$$\frac{n_{baryon}}{n_{\gamma}} < 10^{-20}$$

Mass of largest asteroid, Ceres:  $10^{21}$  kg

Area ~ Kazakhstan: Population ~ one small dog

- CPV in the SM is ensconced in a single unitary matrix



## The Nobel Prize in Physics 2008

"for the discovery of the mechanism of spontaneous broken symmetry in subatomic physics"

"for the discovery of the origin of the broken symmetry which predicts the existence of at least three families of quarks in nature"



Photo: University of Chicago

**Yoichiro Nambu**

🏆 1/2 of the prize



Photo: KEK

**Makoto Kobayashi**

🏆 1/4 of the prize



Photo: Kyoto University

**Toshihide Maskawa**

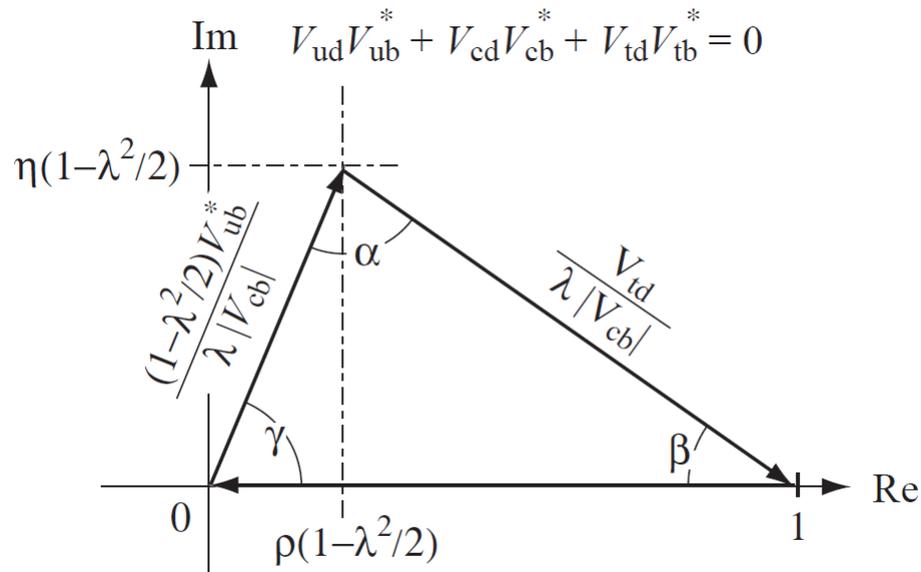
🏆 1/4 of the prize

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

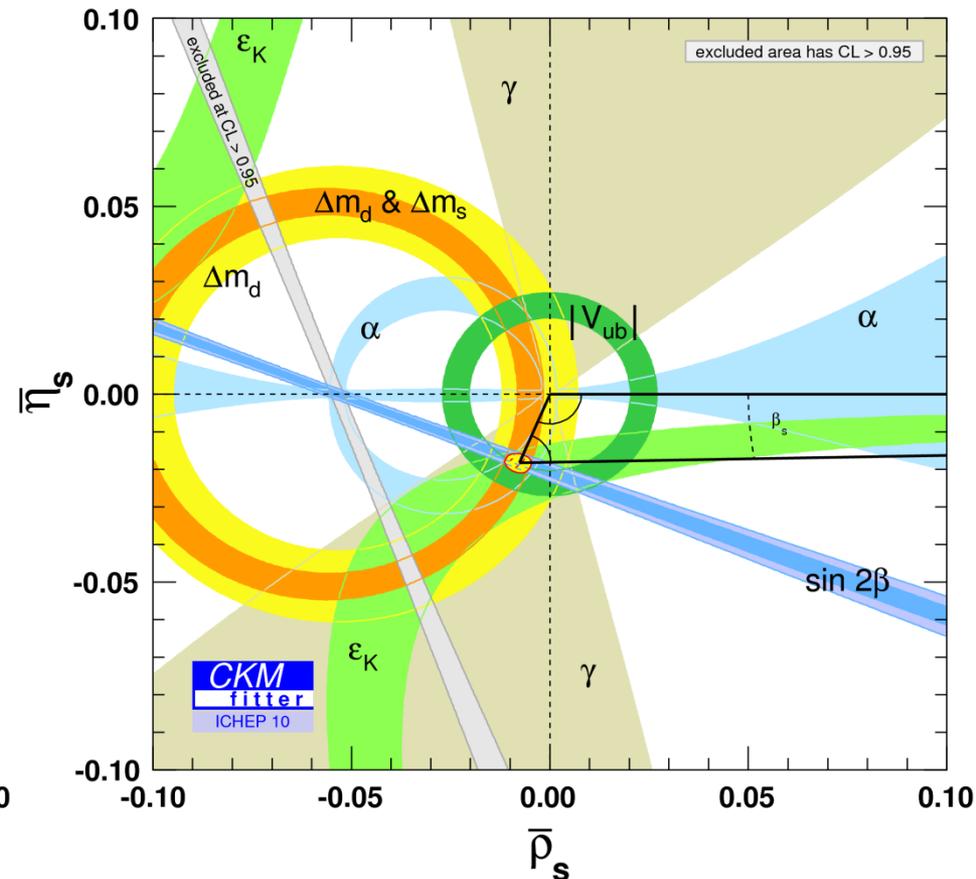
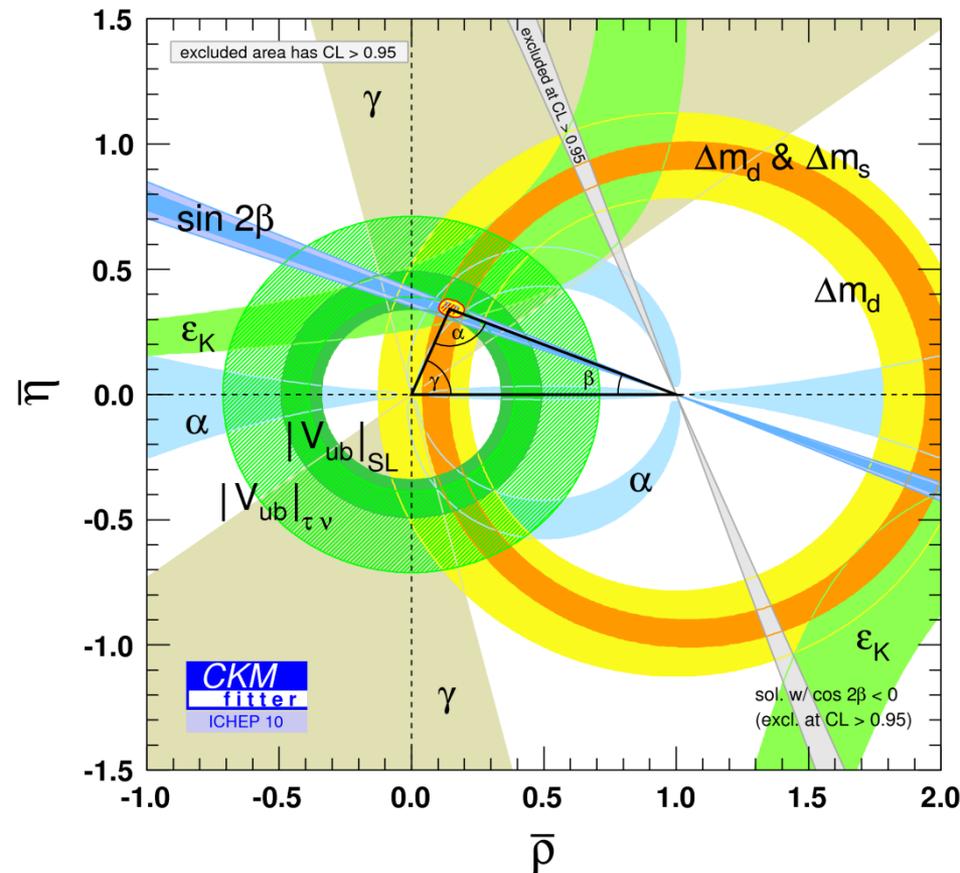
- The CKM matrix
- Three real parameters
- One complex phase violates CP

- The phase is most readily observed in the  $b$ -quark system

- Product of rows and columns are constrained by unitarity
- Of the nine relationships, six form a unitarity triangle
- The most well-known triangle is:



- Couplings, rates and mixings constrain magnitudes
- Asymmetries and mixings constrain phases

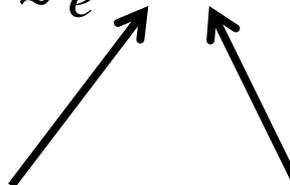


# Mixing observables

- The most basic hamiltonian of anything

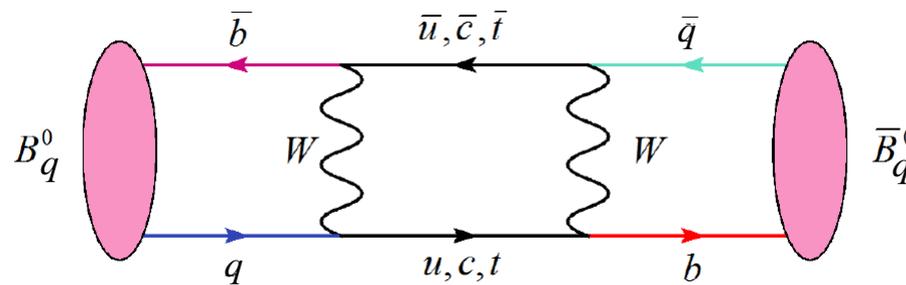
$$H = i \frac{d}{dt} |X\rangle = \left( M_X - \frac{i}{2} \Gamma_X \right) |X\rangle$$

- Because:

$$|X\rangle(t) \sim e^{-iHt} \sim e^{(-iM - \Gamma)t}$$


Wave-like propagation
Decay

- It's weird, it's confusing... it must be quantum mechanics
- In the  $b$ -system, for example, we have two coupled states



- Simplest one-line hamiltonian is now a matrix

$$i \frac{d}{dt} \begin{pmatrix} |B_q^0(t)\rangle \\ |\bar{B}_q^0(t)\rangle \end{pmatrix} = \left( \underline{\underline{M}}_q - \frac{i}{2} \underline{\underline{\Gamma}}_q \right) \begin{pmatrix} |B_q^0(t)\rangle \\ |\bar{B}_q^0(t)\rangle \end{pmatrix}$$

- Off-diagonal elements provide mixing and interference

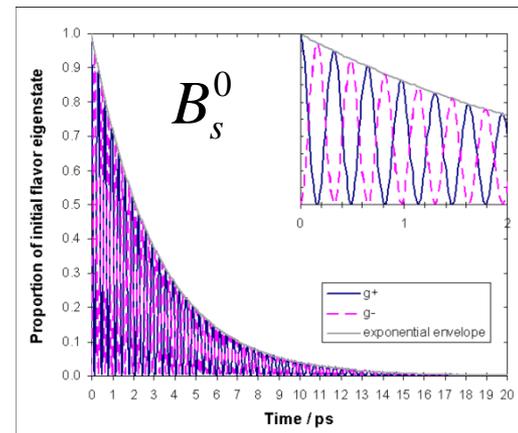
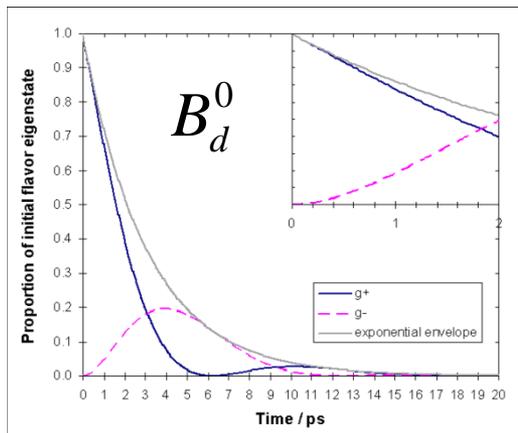
- So, it's not a diagonal matrix... OK
- let's diagonalize it to find:

$$|B_H\rangle(t) \sim e^{-iHt} \sim e^{(-iM_H - \Gamma_H)t}$$

These are the mass-decay-eigenstates

$$|B_L\rangle(t) \sim e^{-iHt} \sim e^{(-iM_L - \Gamma_L)t}$$

- Not the flavour states, a time-dependent mixture of them!



➤ Four simple observables:

1. Average width  $\bar{\Gamma}, \Gamma_{11} + \Gamma_{22}$

2. Average mass  $\bar{M}, M_{11} + M_{22}$

3. Width Difference  $\Delta\Gamma_q = (\Gamma_H^q - \Gamma_L^q) = 2|\Gamma_{12}^q| \arg\left\{\frac{\Gamma_{12}^q}{M_{12}^q}\right\}$

4. Mass Difference  $\Delta m_q = (M_H^q - M_L^q) = 2|M_{12}^q|$

➤ And we also have a **phase**, which violates CP:

$$\phi_q = \arg\left\{-\frac{M_{12}^q}{\Gamma_{12}^q}\right\} \quad \text{and/or} \quad \alpha_{fs}^q = \text{Im}\left\{\frac{\Gamma_{12}^q}{M_{12}^q}\right\}$$

➤ All very predictable observables in the SM

# Flavour-specific asymmetry

... a smoking gun for new physics??

1.  $p\bar{p}$ -interactions within a symmetric experiment
2. Correct all experimental biases (magnets, mis-id ...)

3. Observe  $N(\mu^+ \mu^+) \neq N(\mu^- \mu^-)$

4. In the SM, the favoured way to make charge asymmetry is if:

$$b\bar{b} \longrightarrow \mu^+ \mu^+ \neq b\bar{b} \longrightarrow \mu^- \mu^-$$

5. Which comes from  $B^0$ -mixing:

$$b\bar{b} \Rightarrow \bar{B}^0 B^0 \sim \bar{B}^0 \bar{B}^0 \rightarrow \mu^+ \mu^+ X \neq b\bar{b} \Rightarrow \bar{B}^0 B^0 \sim B^0 B^0 \rightarrow \mu^- \mu^- X$$

➤ In the standard model it is almost negligible

$$A^b \approx \frac{a_{fs}^s + a_{fs}^d}{2} \quad SM = (-2.0 \pm 0.3) \times 10^{-4} \quad D\emptyset \approx (-1 \pm 0.3)\%$$

- $a_{fs}$  is very sensitive to new physics (NP) even if:
  - ❖ Tree-level processes are SM-dominated
  - ❖ SM flavour structure
  - ❖ Unitary CKM
- With very weird scenarios (like *leptoquarks*)
  - Probe NP mixing, interference and/or decays
- Usual formula is modified:

$$a^{SM} \approx \text{Im} \left\{ \frac{\Gamma_{12}^{SM}}{M_{12}^{SM}} \right\}$$

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- With very weird scenarios (like *leptoquarks*)
  - Probe NP mixing, interference and/or decays
- If we allow a single NP phase in the mixing  $\Theta$

$$a^{NP} \approx \text{Im} \left\{ \frac{\Gamma_{12}^{SM}}{M_{12}^{SM}} \right\} \cos \Theta - \text{Re} \left\{ \frac{\Gamma_{12}^{SM}}{M_{12}^{SM}} \right\} \sin \Theta$$

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  - (first part is just the SM value)

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  - ❖ Tree-level processes are SM-dominated
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- With very weird scenarios (like *leptoquarks*)
  - Probe NP mixing, interference and/or decays
- If we allow a single NP phase in the mixing  $\Theta$ 
  - (first part is just the SM value)

$$a^{NP} \approx 2.1 \times 10^{-5} \cos \Theta + 4.0 \times 10^{-3} \sin \Theta$$

- Up to **200-times** the SM!!! [[[ ... still... < DØ measurement ]]]

# Flavour-specific asymmetry

At LHCb

- At the LHC we have extra complications in the measurement
- Polluting asymmetries, which are all much larger than  $a_{fs}$ 
  - Production asymmetry  $\delta_p \sim (10^{-2})$
  - Detector asymmetry  $\delta_c \sim (10^{-2})$
  - Background asymmetry  $\delta_b \sim (10^{-3})$
- Use a, time-dependent, untagged, simultaneous fit to  $B_s + B_d$
- Subtract two asymmetries to eliminate detector component

$$\Delta A_{fs} = \frac{a_{fs}^s - a_{fs}^d}{2} = (2.1 \pm 0.3) \times 10^{-4}$$

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NB:  $D^0$

$$\left( A^b \sim \frac{a_{fs}^s + a_{fs}^d}{2} \sim -(2.0 \pm 0.3) \times 10^{-4} \right)$$

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

$$A_{fs}^q(t) = \frac{a_{fs}^q}{2}$$



$10^{-3} \rightarrow 10^{-5}$

$$-\left( \frac{a_{fs}^q}{2} \right) \frac{\cos(\Delta m_q t)}{\cosh(\Delta \Gamma_q t / 2)}$$

Very Complicated

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

$$A_{fs}^q(t) = \frac{a_{fs}^q}{2} \left[ -\frac{\delta_c^q}{2} - \left( \frac{a_{fs}^q}{2} + \frac{\delta_p^q}{2} \right) \frac{\cos(\Delta m_q t)}{\cosh(\Delta \Gamma_q t / 2)} + \frac{\delta_b^q}{2} \left( \frac{B}{S} \right)^q \right]$$

$10^{-3} \rightarrow 10^{-5}$        $10^{-2}$        $10^{-2}$        $10^{-3}$

➤ Polluting asymmetries are much larger than  $a_{fs}$

- Detector asymmetry  $\delta_c$        $\sim(10^{-2})$
- Production asymmetry  $\delta_p$        $\sim(10^{-2})$
- Background asymmetry  $\delta_b$        $\sim(10^{-3})$

$$\delta_c = \frac{\varepsilon(\bar{f}_i)}{\varepsilon(f_i)} - 1$$

$$\delta_p = \frac{N(\bar{I}_0)}{N(I_0)} - 1$$

$$\delta_b = \frac{\bar{B}/\bar{S}}{B/S} - 1$$

- We measure time-dependent decay rates:

$$\Gamma(f) = N e^{-\Gamma t} (1 + A_c) \left[ (1 + A_{fs}) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (2A_p - A_{fs}) \cos(\Delta m t) \right]$$

$$\rightarrow N e^{-\Gamma t} \left[ (1 + A_c + A_{fs}) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (2A_p - A_{fs} + 2A_p A_c) \cos(\Delta m t) \right]$$

$$\Gamma(\bar{f}) \rightarrow N e^{-\Gamma t} \left[ (1 - A_c - A_{fs}) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (A_{fs} - 2A_p + 2A_p A_c) \cos(\Delta m t) \right]$$

- $A_c$ ,  $A_p$  and  $A_{fs}$  are correlated and cannot be separately fitted
- First, reparameterise

- Just to make it easier to see what we're doing...

$$\Gamma(f) = Ne^{-\Gamma t} \left[ (1 + x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 + x_3) \cos(\Delta mt) \right]$$

$$\Gamma(\bar{f}) = Ne^{-\Gamma t} \left[ (1 - x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 - x_3) \cos(\Delta mt) \right]$$

where:  $x_1 = A_c + a_{fs}$        $x_2 = 2A_c A_p$        $x_3 = 2A_p - a_{fs}$

- production asymmetry is an *initial state asymmetry*
- Changes the mixing amplitude, does not change the physics
- Fit for  $x_1$  independently, which now only has detector asym

- Take  $B_s/B_d$  with the same final states ( $f=KK\pi\mu$ )

$$\Gamma(f) = Ne^{-\Gamma t} \left[ (1 + x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 + x_3) \cos(\Delta mt) \right]$$

$$\Gamma(\bar{f}) = Ne^{-\Gamma t} \left[ (1 - x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 - x_3) \cos(\Delta mt) \right]$$

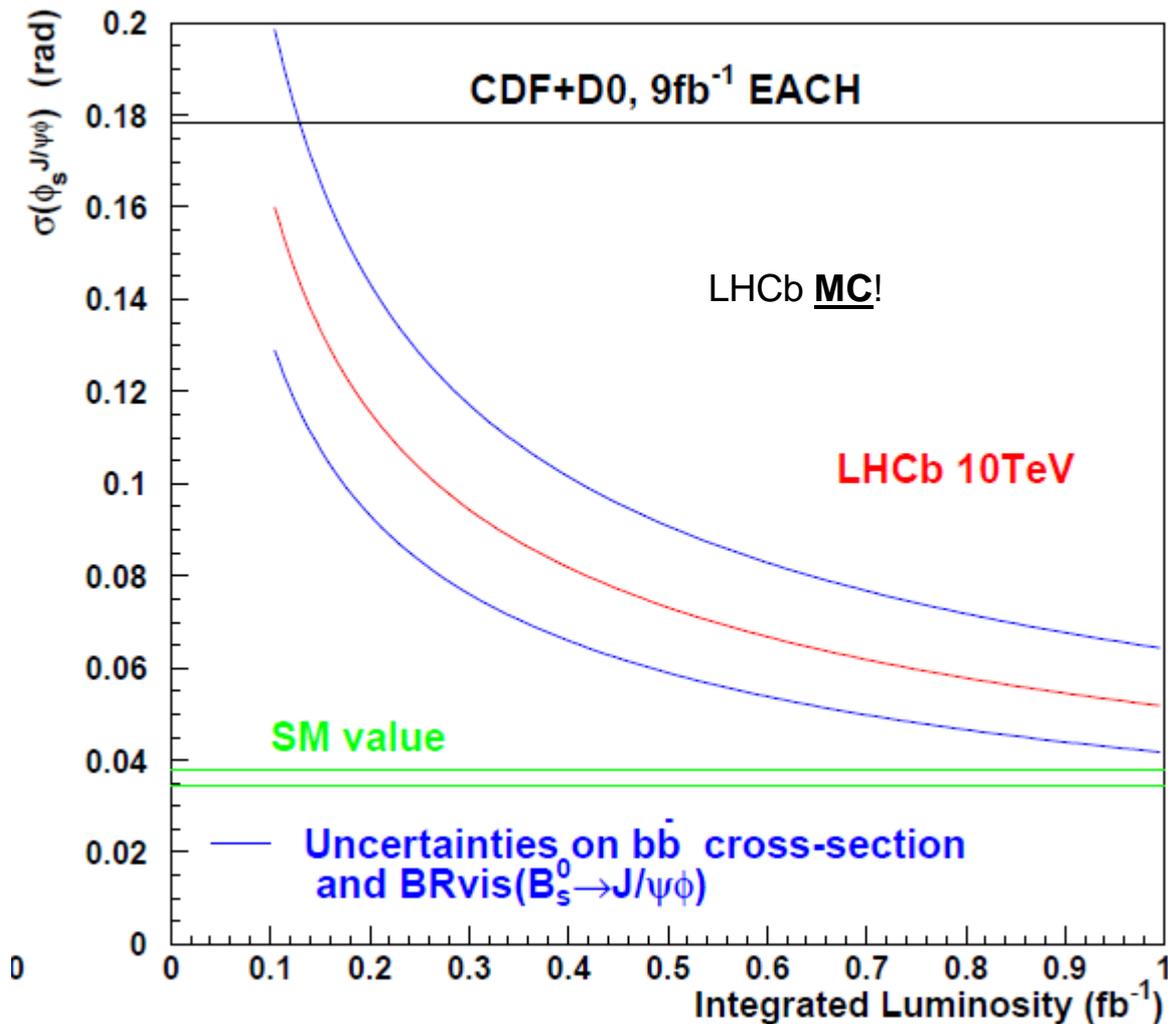
where:  $x_1 = A_c + a_{fs}$        $x_2 = 2A_c A_p$        $x_3 = 2A_p - a_{fs}$

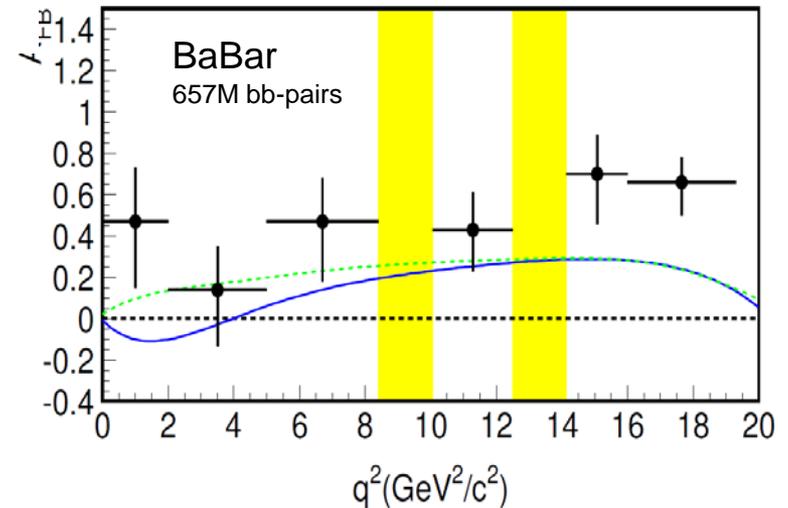
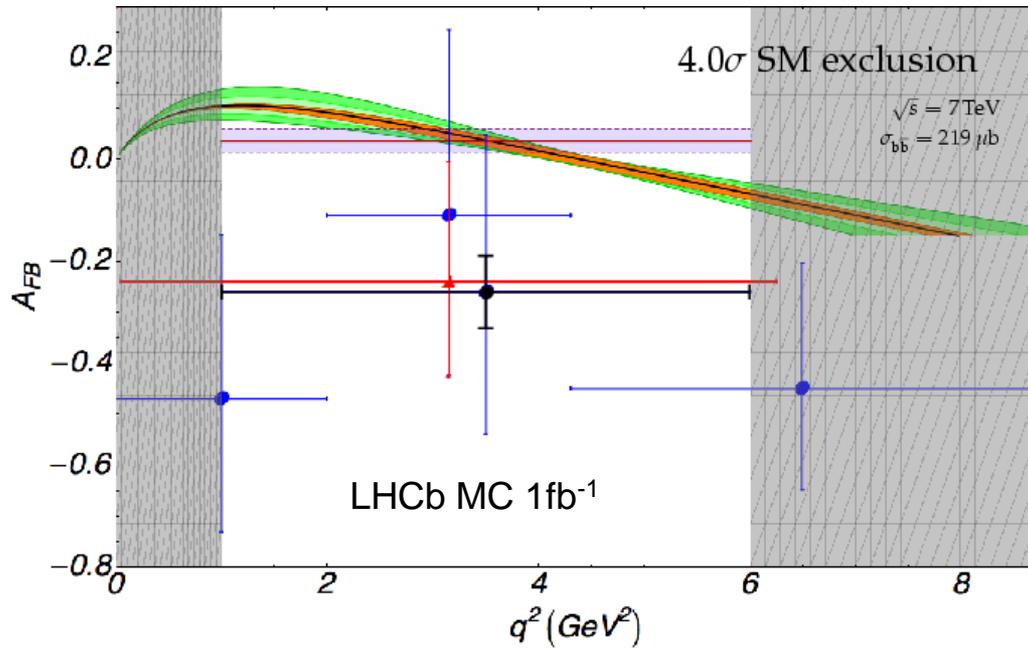
- All **production asymmetry** is in  $x_2/x_3$ , just throw it away
- Measure the **difference** between  $B_s$  and  $B_d$

$$\Delta A_{fs}^{s,d} = \frac{x_1^s - x_1^d}{2} = \frac{a_{fs}^s - a_{fs}^d}{2} \qquad SM = \left( +2.5^{+0.5}_{-0.6} \right) \times 10^{-4}$$

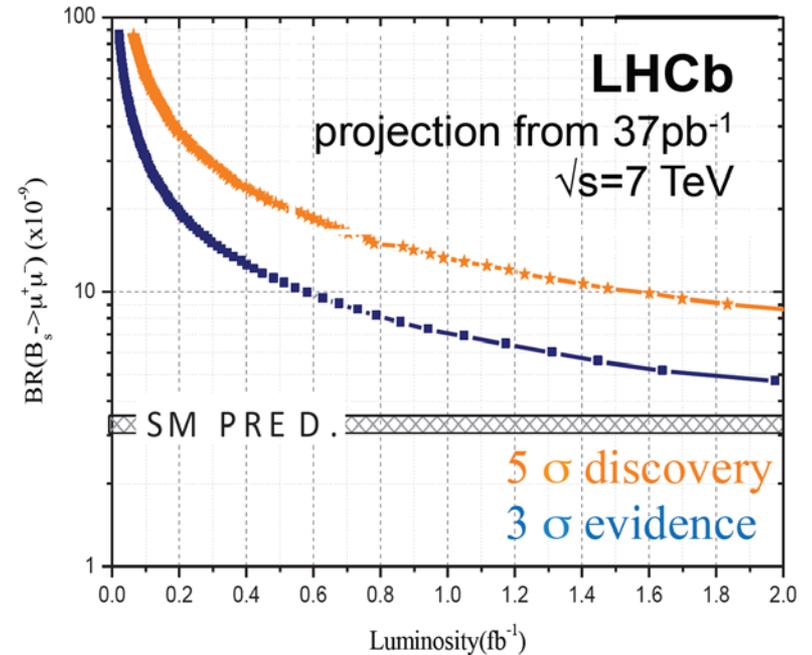
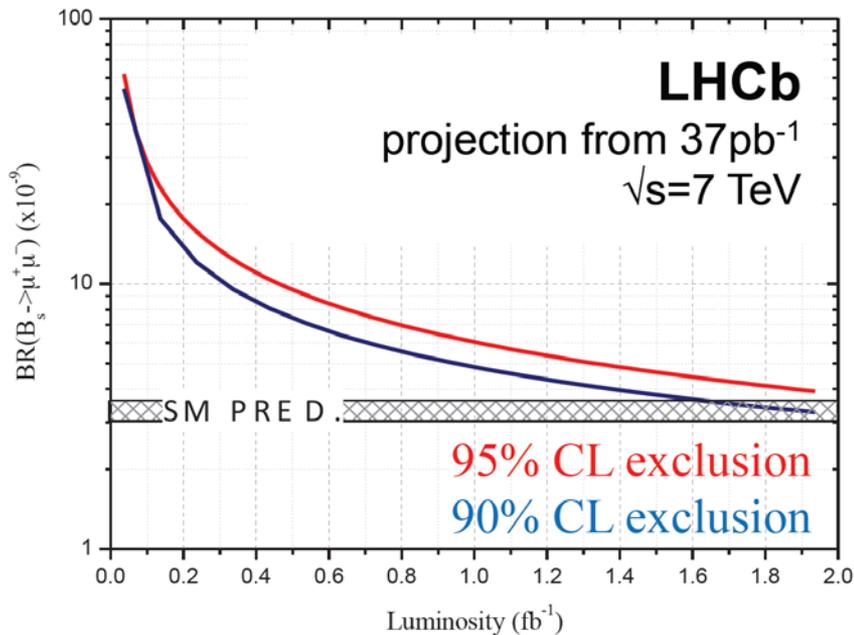
- MC sensitivities, Real data yields and systematics
  - $0.1 \text{ fb}^{-1}$   $\sigma \sim 5 \times 10^{-3}$  ... First result (2011)
  - $1.0 \text{ fb}^{-1}$   $\sigma \sim 2 \times 10^{-3}$  ...  $5\sigma$  observation? (2012/2013)

# LHCb projections



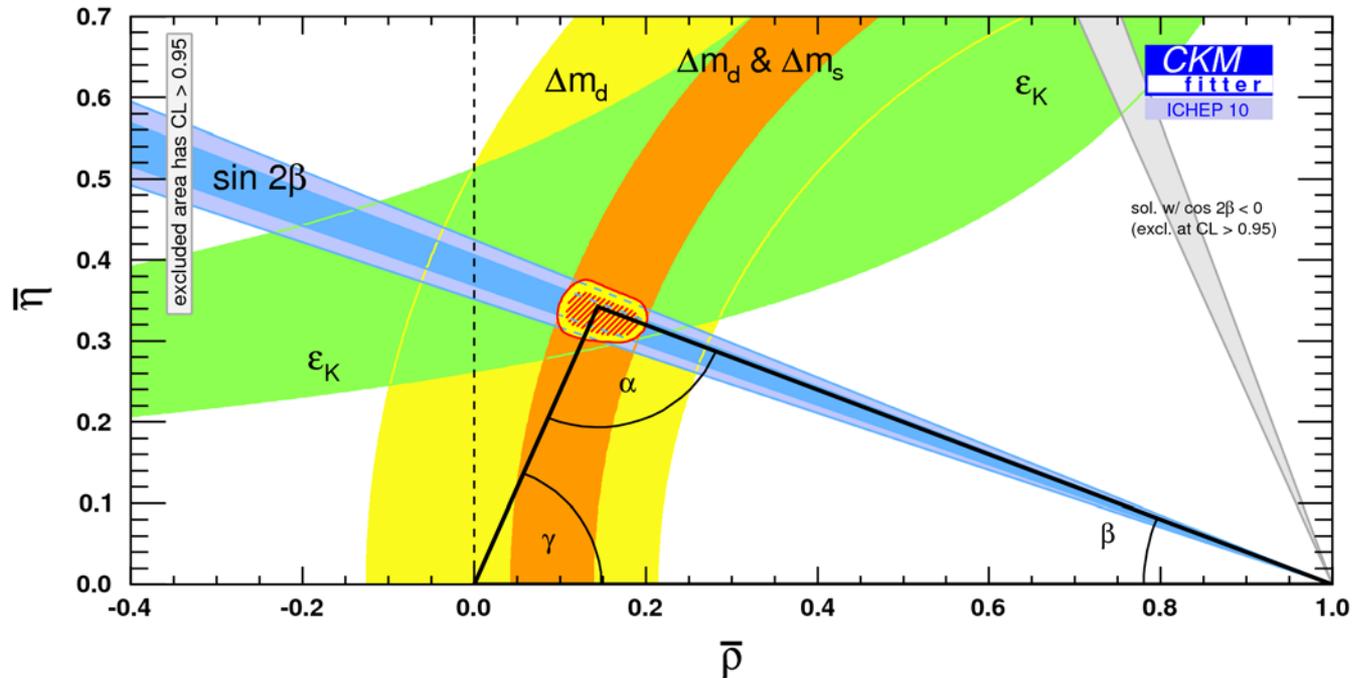


➤ LHCb will exclude most SUSY models this year!



# Misc

- Check loop-level observables
- Would need a very accurate determination of  $d_{\text{md}}/d_{\text{ms}}$



➤  $B_s^0 \rightarrow J/\psi \Phi$

- Directly Measure  $\sin \phi_s$
- $\sigma(\phi_s) = 0.05^c$  in  $1 \text{ fb}^{-1}$

➤  $a_{fs}^s$

- Effectively Measures

$$\text{Im} \left\{ \frac{\Gamma_{12}}{M_{12}} \right\} \cos \Theta - \text{Re} \left\{ \frac{\Gamma_{12}}{M_{12}} \right\} \sin \Theta$$

- $\sigma(\Theta) = 0.5^c$  in  $1 \text{ fb}^{-1}$

➤ But they constrain NP differently

- Effective power enhanced
- NB physical limit of  $a_{fs}$  is at  $4 \times 10^{-3} < \text{current } D\emptyset \text{ result!}$

